

## Extreme intensity pulses in semiconductor lasers with optical injection or feedback

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UNIVERSITAT POLITÈCNICA  
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*Campus d'Excel·lència Internacional*

Workshop on Nonlinear Optics and  
Nanophotonics  
Sao Paulo, Brazil, December 2013





- José A. Reinoso (Universidad Nacional de Educación a Distancia, UNED, Spain)



- Jordi Zamora Munt (IFISC, Mallorca, Spain)
- B. Garbin, M. Feyereisen, S. Barland, M. Giudici (INLN, Nice, France)
- Jorge Tredicce (INLN, now at Universite de la Nouvelle Calédonie)
- Jose Rios Leite (Universidade Federal de Pernambuco, Recife, Brasil)

# Where are we?

## UPC Campus Terrassa



1. Barcelona
2. Castelldefels
3. Igualada
4. Manresa
5. Mataró
6. Sant Cugat del Vallès
7. Terrassa
8. Vilanova i la Geltrú



## Gaia research Building

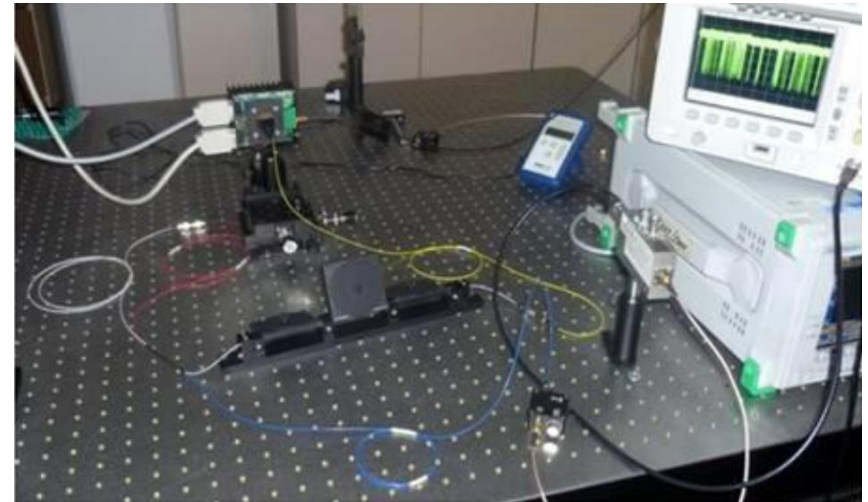
Viernes, 25 de septiembre de 2009 *Diari de Terrassa*



El edificio Gaia centraliza grupos científicos consolidados y emergentes.



# Nonlinear optics and laser dynamics research labs



- Introduction to extreme events & optical rogue waves.
- Semiconductor laser with **optical injection**: experimental observations & numerical results.
- Semiconductor laser with **optical feedback**: numerical results.
- Summary and conclusions.



RWs are rare, ultra-high waves that fall outside (and far from) the main part of long-tailed probability distributions.

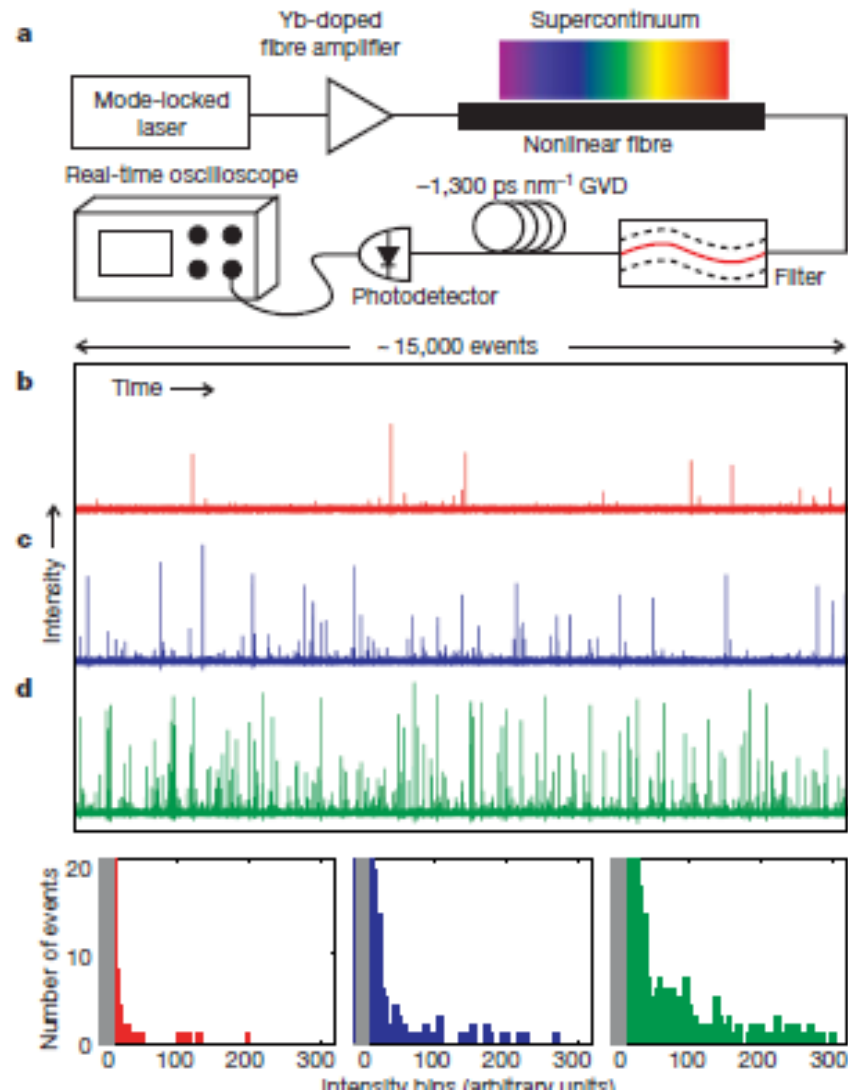


The Great Wave of Kanagawa, Katsushika Hokusai. Source: Wikipedia

- A RW is usually defined as a wave that is two times the significant wave height of the area. The significant wave height is the average of the highest one-third of waves that occur over a given period.
- Serious problem for the design of off-shore platforms.



# Optical RWs: first observation



D. R. Solli et al, Nature 450, 1054 (2007)





## A recent review of the state-of-the-art

IOP PUBLISHING

JOURNAL OF OPTICS

J. Opt. **15** (2013) 060201 (9pp)

[doi:10.1088/2040-8978/15/6/060201](https://doi.org/10.1088/2040-8978/15/6/060201)

### EDITORIAL

# Recent progress in investigating optical rogue waves

N Akhmediev<sup>1</sup>,  
J M Dudley<sup>2</sup>, D R Solli<sup>3,4</sup>  
and S K Turitsyn<sup>5</sup>

The science of rogue waves in optics is now over five years old, and it has emerged as an area of broad interest to researchers across the physical sciences [1]. This area of study was initiated by the pioneering measurement of

# RWs in optically injected semiconductor lasers

- We have recently shown experimentally and numerically that continuous-wave optically injected semiconductor lasers can display huge intensity pulses that we identified as deterministic rogue waves.
- These pulses can be predicted with a certain anticipation time.
- They are generated by an external crisis-like process.
- Noise can either enhance or diminish their probability of occurrence.

C. Bonatto et al, *Deterministic optical rogue waves*, PRL 107, 053901 (2011).

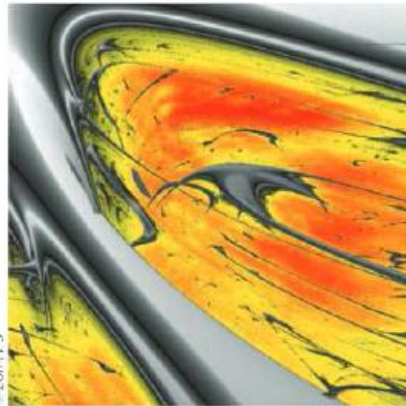
J. Zamora-Munt et al, *Rogue waves in optically injected lasers: origin, predictability and suppression*, PRA 87, 035802 (2013).

# Highlighted in OPN and Nature Photonics

## ROGUE WAVES

### Surely deterministic

Phys. Rev. Lett. **107**, 053901 (2001)



The physical mechanism behind the formation of rogue waves — waves with seemingly spontaneous large amplitudes — has long been an interesting research topic for oceanologists and physicists, including those working in photonics. An important question is whether rogue waves can be described by a fully deterministic process with noise as a driving force. Cristian Bonatto and co-workers from Spain, France and Brazil recently carried out an investigation into the generation of rogue waves using a semiconductor laser that received optical injection from a continuous-wave master laser. The researchers not only showed that sporadic high-amplitude pulses can be observed with such a simple and inexpensive laser set-up, but also concluded that the rogue waves they observed are generated from deterministic nonlinearities. Their conclusion was based on good qualitative agreement between experimental results and simulated results from a simple, deterministic noise-free rate equation model.

RIV

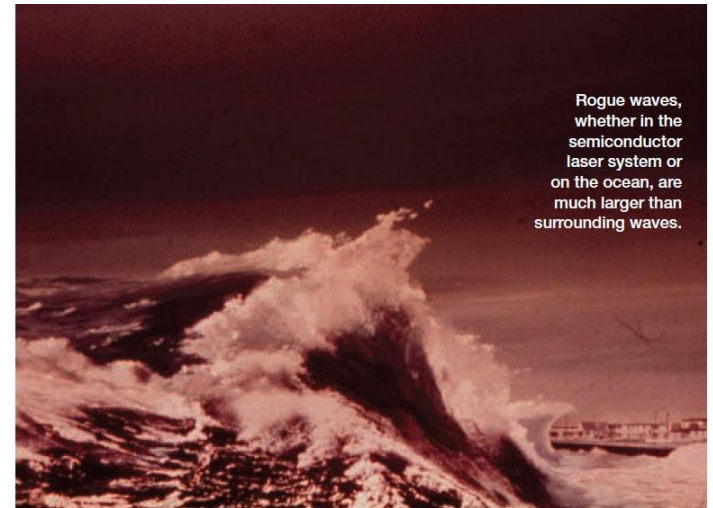
## SCATTERINGS | NEWS

### Deciphering Rogue Waves

Rare pulses with giant intensities—the optical equivalent of rogue ocean waves—have been shown to occur in common laser systems. A team of researchers from Spain, France and Brazil found a way to generate rogue waves and developed a model for understanding them as a result of a deterministic nonlinear process (Phys. Rev. Lett. **107**, 053901). Extremely high waves have been a subject of interest over the past decade in oceanography as well as in other fields (including optics), but we still don't fully understand what triggers them and how they develop.

Rogue waves on the ocean are typically twice the height of surrounding waves and have steep sides, like “a wall of water.” They have high amplitude, with a fast rise and fast fall. In the laser system demonstrated by the researchers, a rogue wave has an intensity so high that—according to Gaussian statistics—it should be vanishingly improbable. Such waves are unusual, but occur more often than Gaussian statistics can explain. Rogue waves also can be destructive: paper coauthor Jorge Tredicce at the Université de Nice Sophia Antipolis (France) says, “in mode-locked lasers, those extreme pulses may damage the optics and the crystal ... it is the death of the laser!”

Light from a continuous wave master laser is injected into a stabilized vertical



Rogue waves, whether in the semiconductor laser system or on the ocean, are much larger than surrounding waves.

Wikimedia Commons

cavity surface-emitting laser (VCSEL) with stabilized pump current and temperature. The VCSEL emitted at 980 nm in a single transverse mode.

Researchers detuned the injection laser from the VCSEL and found that the slave laser output falls into four regions—one of which is stable-locked behavior. As the VCSEL current is increased, the output becomes more and more chaotic. Near the border of the mode-locked region, the researchers found a series of small pulses interrupted by occasional extremely large pulses. Coauthor Cristina Masoller at the Universitat Politècnica de Catalunya (Spain) explains, “we identify two types of optical chaos: one in which rogue waves are rare but they certainly appear and one in which practically they do not exist.” There appear to be areas where rogue waves don't occur even if the behavior is chaotic.

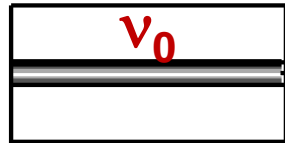
The experiments were inspired by Tredicce's theoretical paper that suggested the existence of huge intensity pulses in this laser system. The researchers found that a simple noise-free rate equation model produced results that agree with the experiments. This allowed them to interpret the sporadic high amplitude pulses as the result of a deterministic nonlinear process.

The group now has a simple system that allows them to experiment with optical rogue waves, as well as a model for describing them. Because rogue waves occur in other systems, ranging from ocean surface to acoustic waves to economics, their work could have implications far beyond the realm of optics. Next, they want to find a mechanism that creates rogue waves in their system, as well as whether they can increase or decrease their likelihood.

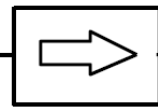
—Yvonne Carls-Powell

# Optically injected semiconductor lasers

Master Laser

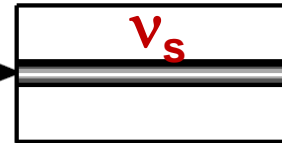


Tunable SCL



Isolator

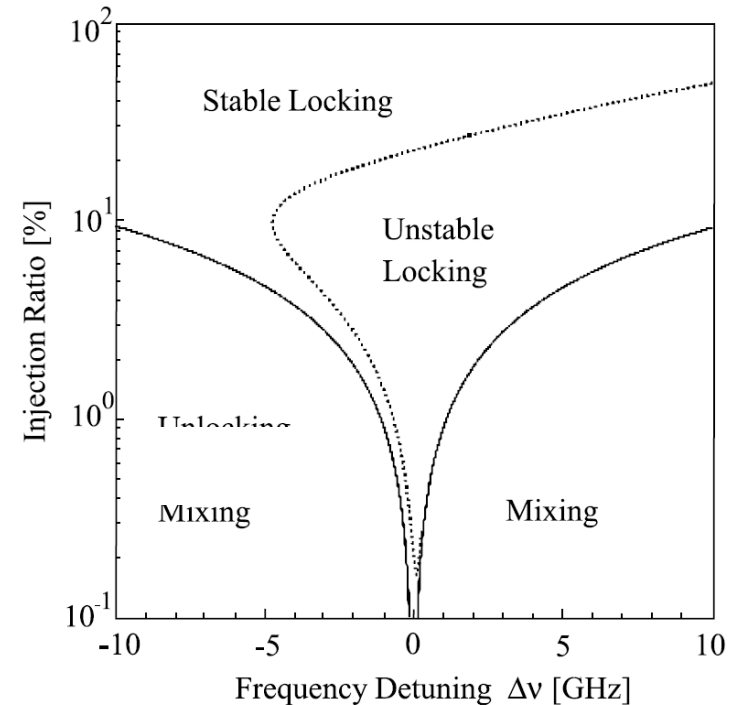
Slave Laser



VCSEL 980 nm

Detection system  
(photo detector,  
oscilloscope,  
spectrum  
analyzer)

- Parameters:
  - Injection ratio
  - Frequency detuning  $\Delta\nu = \nu_s - \nu_0$
- Dynamical regimes:
  - Injection locking (cw output)
  - Period-one oscillation
  - Period-two oscillation
  - Chaos



# Labyrinth bifurcations in optically injected diode lasers

V. Kovanis<sup>1</sup>, A. Gavrielides<sup>2</sup>, and J.A.C. Gallas<sup>3,4,5,a</sup>

<sup>1</sup> Air Force Research Laboratory, 2241 Avionics Circle, Wright-Patterson AFB, Dayton OH 45433, USA

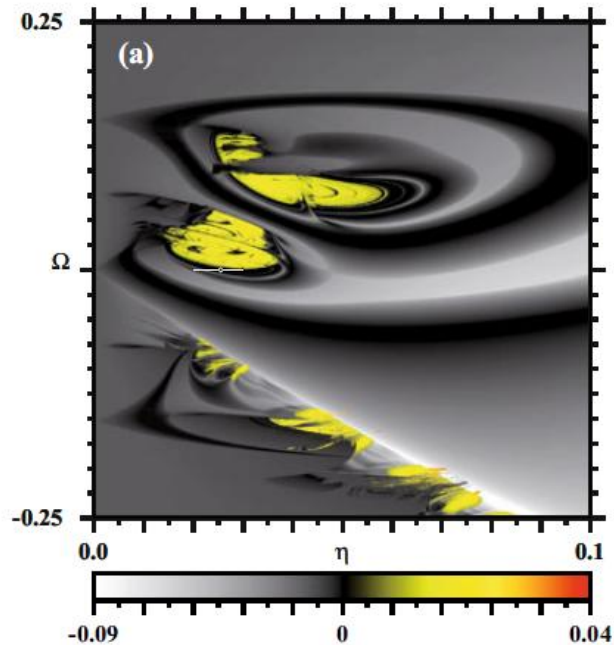
<sup>2</sup> USAF, Research Laboratory, High Power Solid State Lasers Branch, Kirtland AFB, NM 87117, USA

<sup>3</sup> TecEdge, Wright Brothers Institute, 5100 Springfield Street, Dayton OH 45431, USA

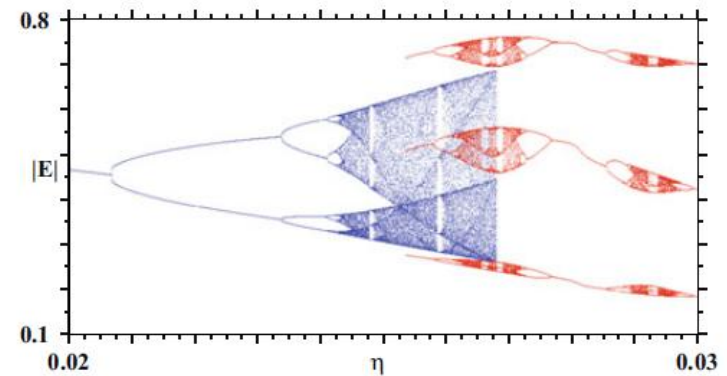
<sup>4</sup> Departamento de Física, Universidade Federal da Paraíba, 58051-970 João Pessoa, Brazil

<sup>5</sup> Instituto de Física, Universidade Federal do Rio Grande do Sul, 91501-970 Porto Alegre, Brazil

## Lyapunov diagram



## Bifurcation diagram

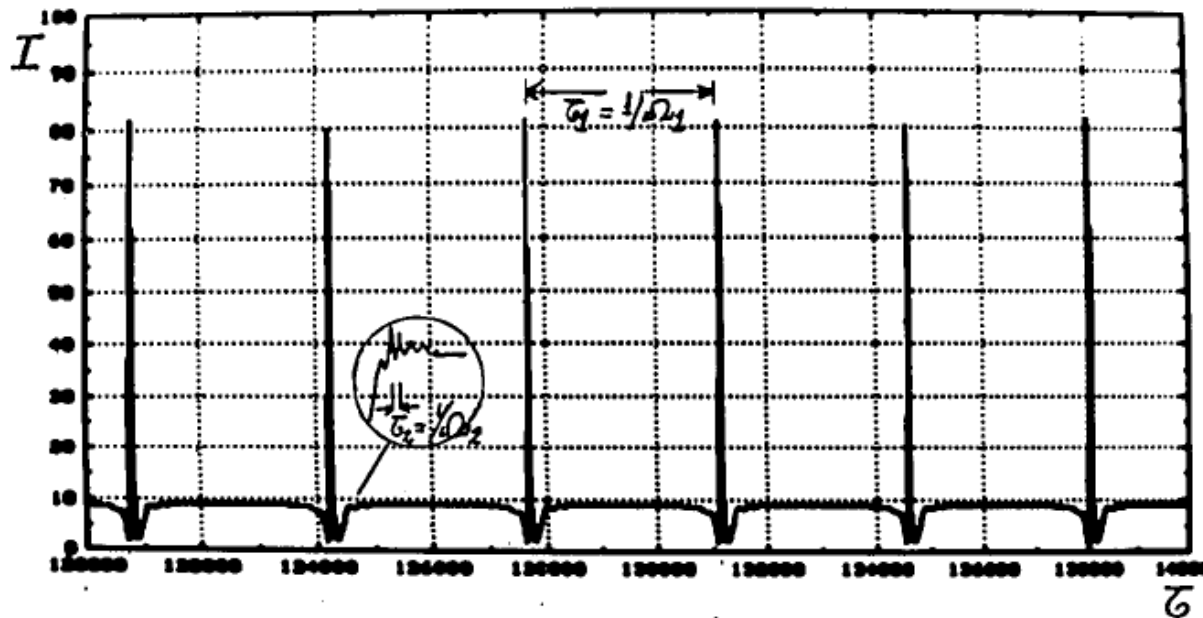




# Instabilities in lasers with an injected signal

J. R. Tredicce, F. T. Arecchi, G. L. Lippi, and G. P. Puccioni

178 J. Opt. Soc. Am. B/Vol. 2, No. 1/January 1985

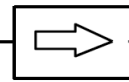


# Experimental setup and control parameters

Master Laser



$\nu_0$   
(fixed)



Isolator

Slave Laser



$\nu_s$   
(variable)

- Photo detector: 9.5GHz
- Oscilloscope: 6 GHz, 20GS/s
- Spectrum analyzer: 72GHz

The **frequency detuning** between the lasers,  $\Delta\nu = \nu_s - \nu_0$ , is controlled by the slave laser pump current,  $I$

When  $I$  increases:

- **Joule heating**
- **the temperature modifies the cavity refractive index**
- **decreases the cavity resonance frequency**

$$\nu_s = g(\text{Temp}) = f(I)$$

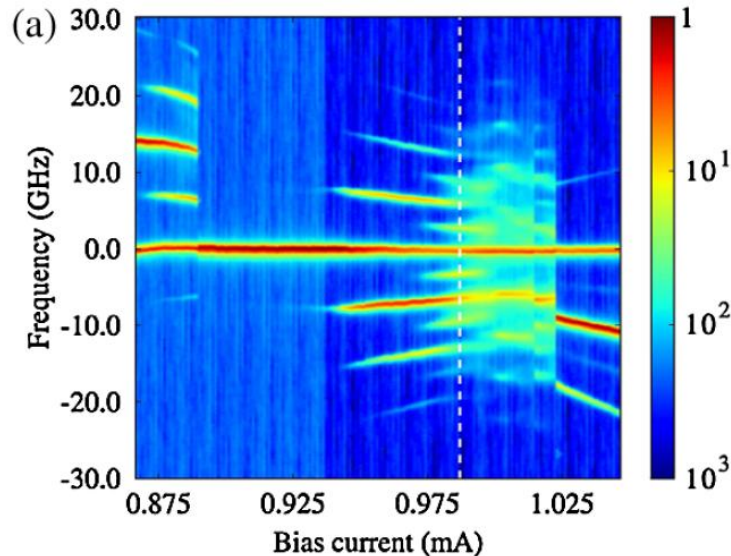
(f approximately linear)

We varied the slave laser pump current and detected the output of the laser:

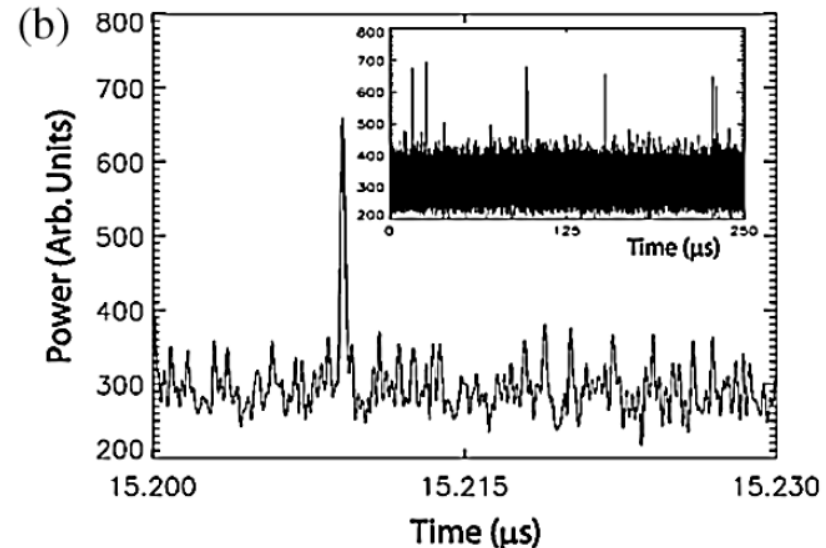
- **Intensity time - series (with a 6 GHz oscilloscope)**
- **Intensity Fourier spectrum (spectrum analyzer)**

# Experimental observations

## Fourier spectrum of the laser intensity



## Time series of the laser intensity



Five regions as  $I$  increases:

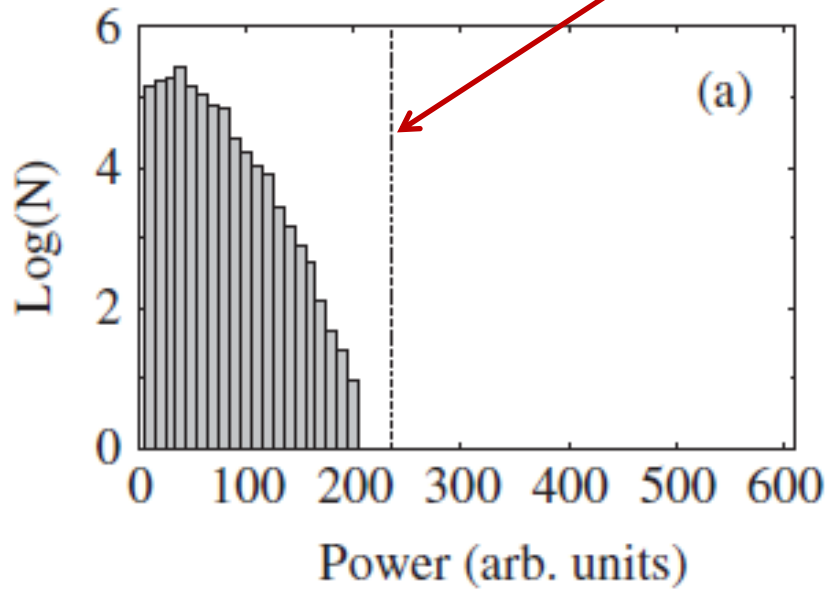
- Beating (independent lasers)
- Period 2 of the beat note
- Stable locking
- Periodic & chaotic oscillations
- Beating (independent lasers again)

(In the chaotic region,  
 $I = 0.976$  mA,  $\Delta\nu = -1.34$  GHz)

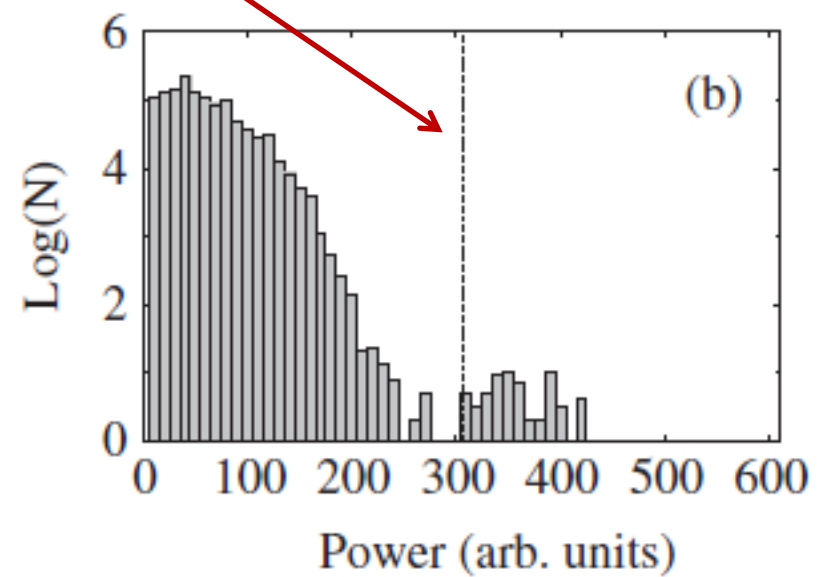
C. Bonatto et al, PRL 107, 053901 (2011)

# Histograms of pulse amplitude

Border = mean value +  $8 \sigma$



$I = 0.972 \text{ mA}$



$I = 0.976 \text{ mA}$

# Governing equations

- The complex optical field, **E** (photon number  $\propto |\mathbf{E}|^2$  )
- The carrier density, **N**

$$\frac{dE}{dt} = \frac{1}{2\tau_p} (1 + i\alpha)(N - 1)E + \underbrace{i\Delta\omega + \sqrt{P_{inj}}}_{\text{optical injection}} + \underbrace{\sqrt{2\beta_{sp} / \tau_N} \xi(t)}_{\text{spontaneous emission noise}}$$

$$\underbrace{\frac{dN}{dt} = \frac{1}{\tau_N} (\mu - N - N|E|^2)}_{\text{Solitary laser parameters: } \alpha \tau_p \tau_N \mu}$$

$\eta$ : injection strength  
 $\Delta\omega$ : frequency detuning

Typical parameter values:

$$\alpha = 3, \tau_p = 1 \text{ ps}, \tau_N = 1 \text{ ns}$$

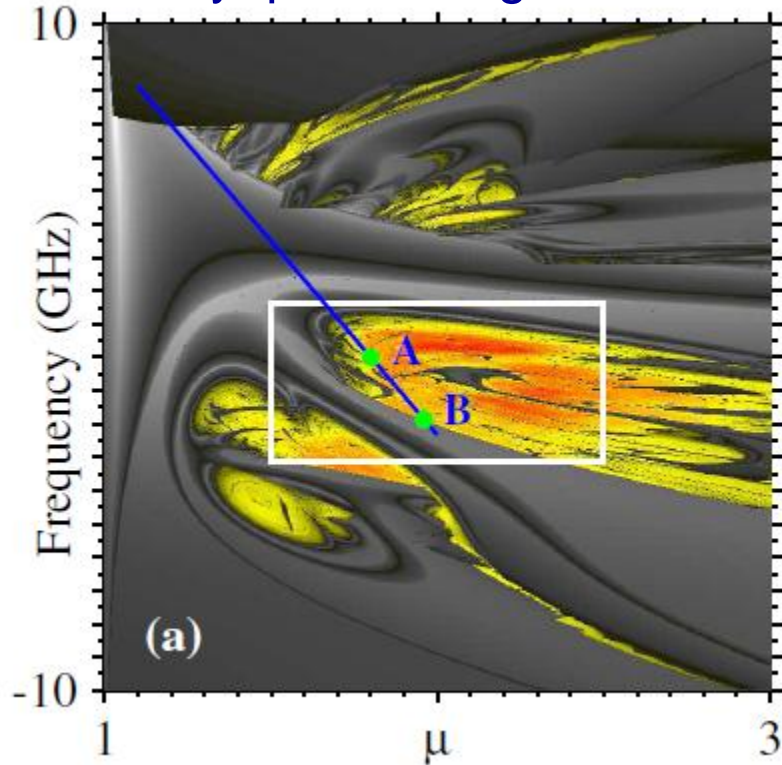
$\mu$ : normalized pump current parameter



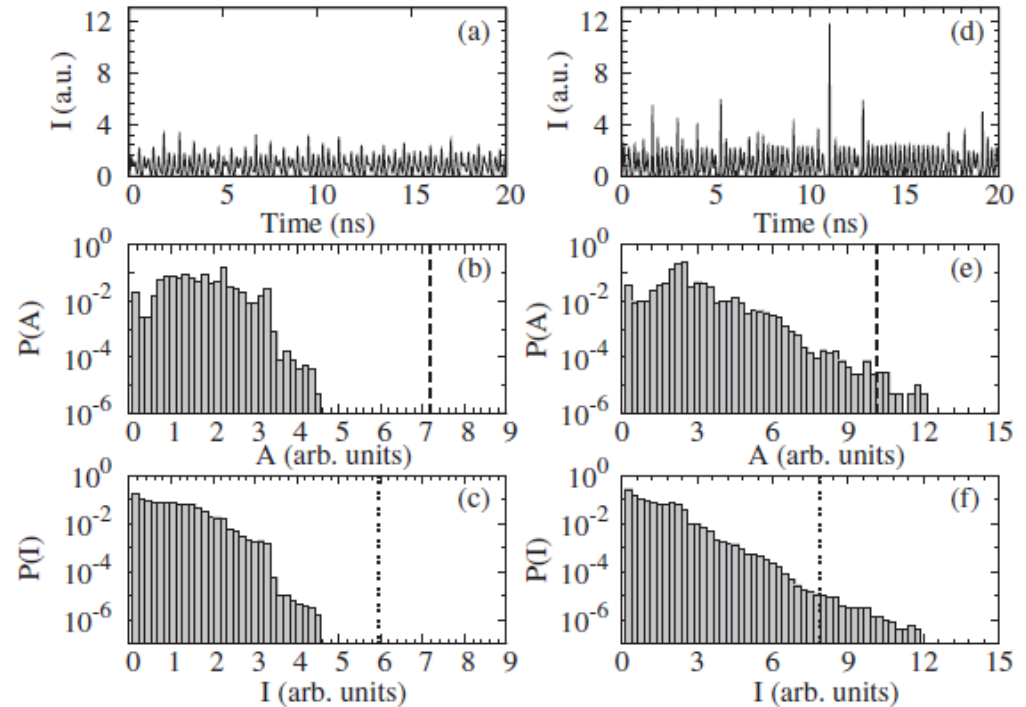
# Deterministic simulations

( $\beta_{sp}=0$ )

## Lyapunov diagram



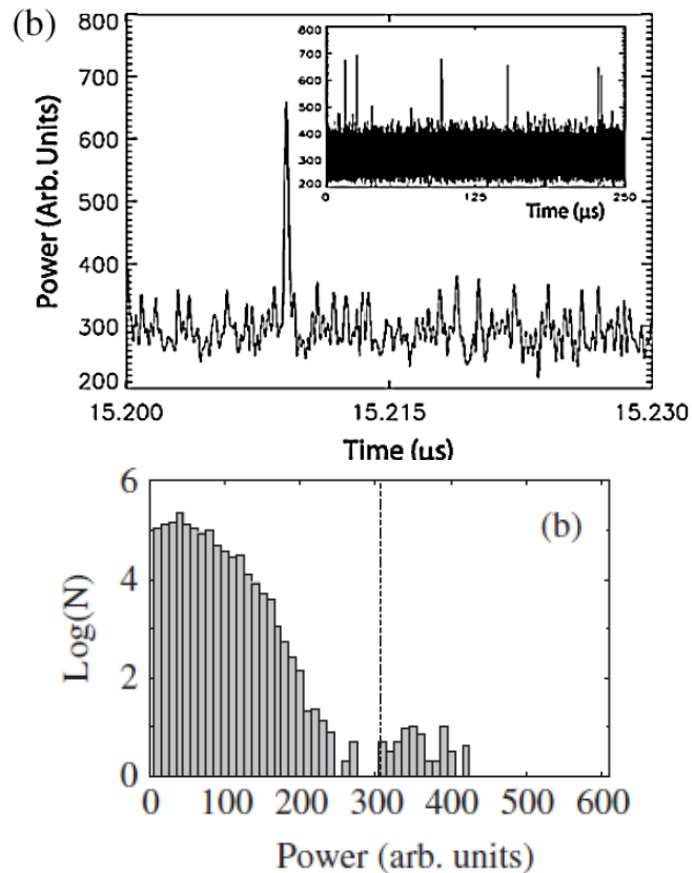
Slave laser pump current



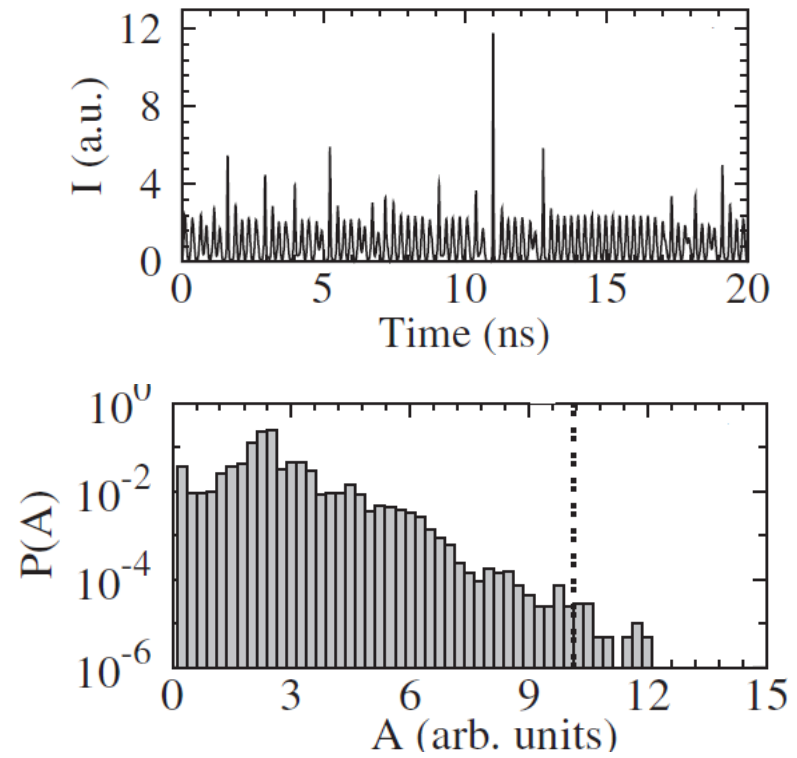
Point A: No RWs

Point B: RWs

## Experiments

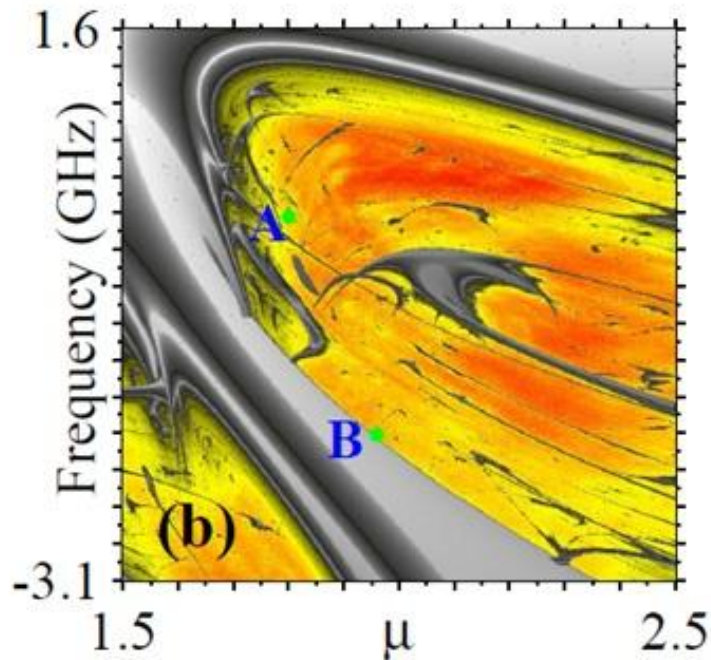


## Simulations

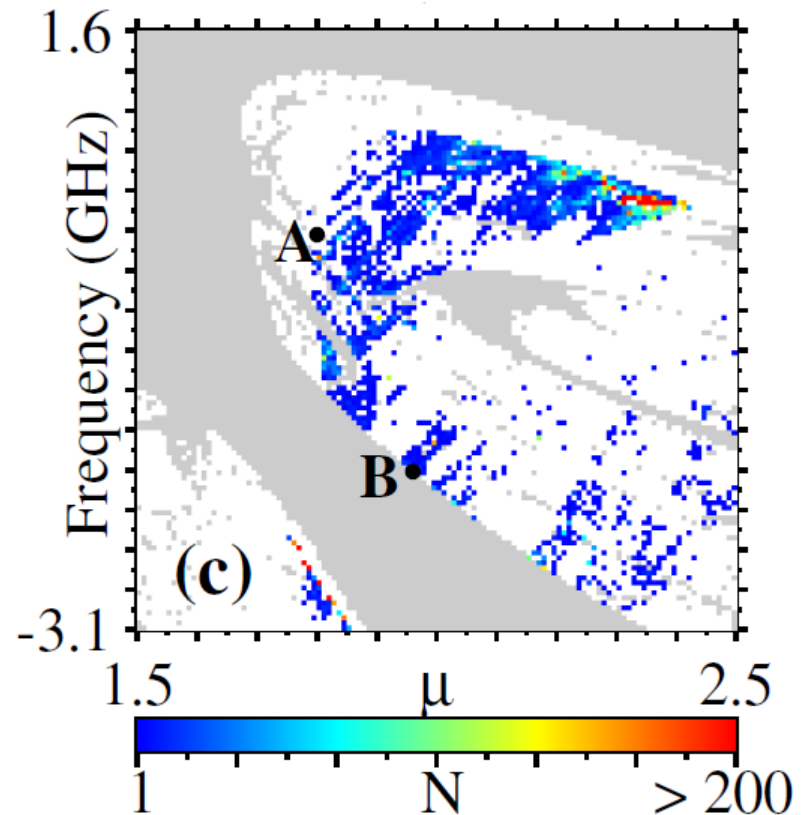


# Rogue waves in the parameter space (pump current, frequency detuning)

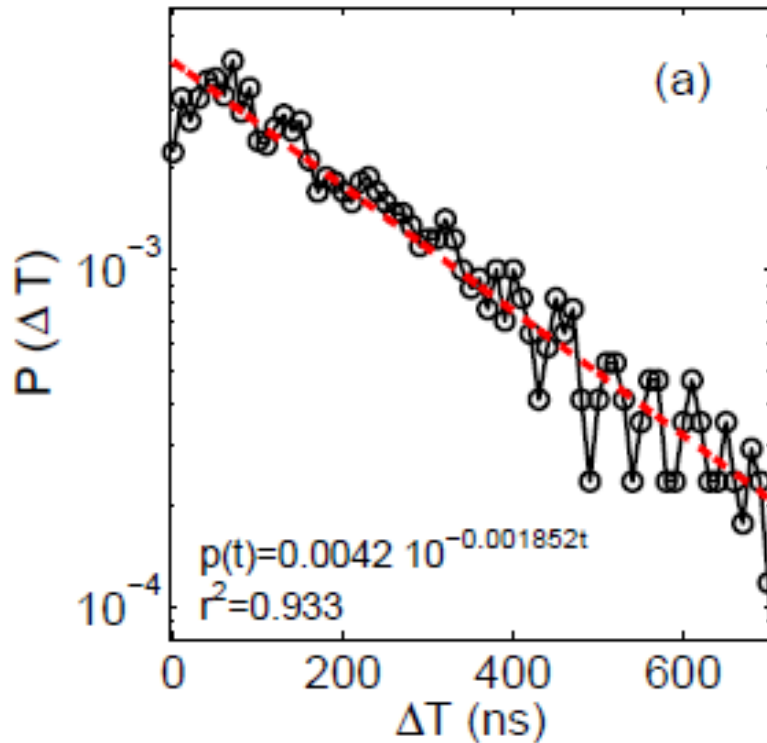
Lyapunov diagram



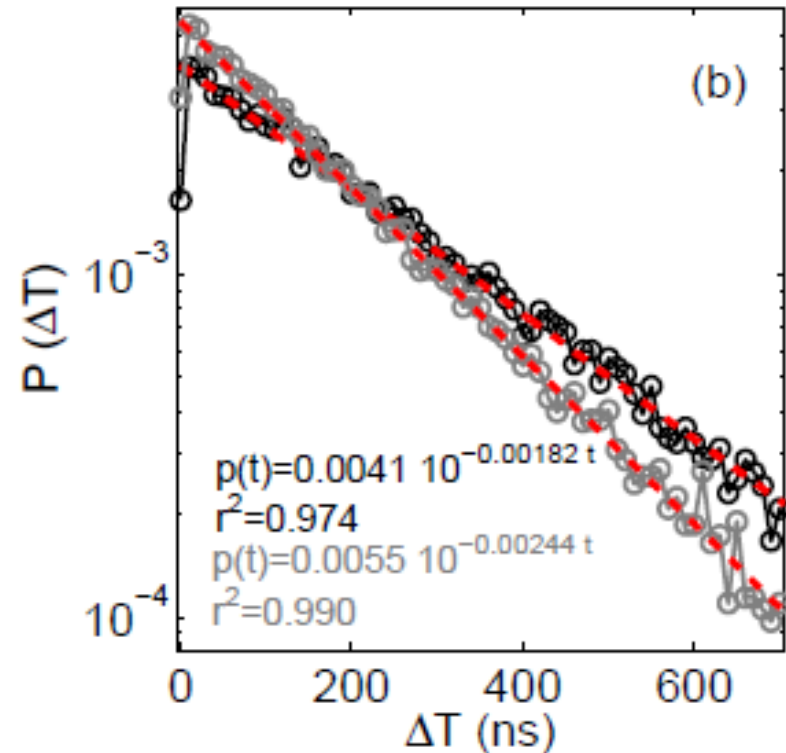
Number of RWs



# Statistics of the RW waiting time

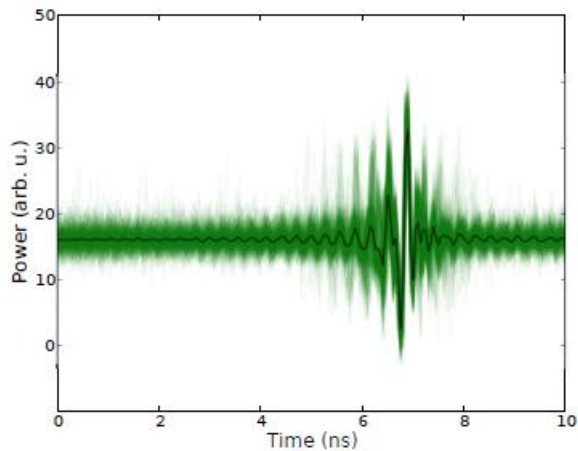


Experimental data

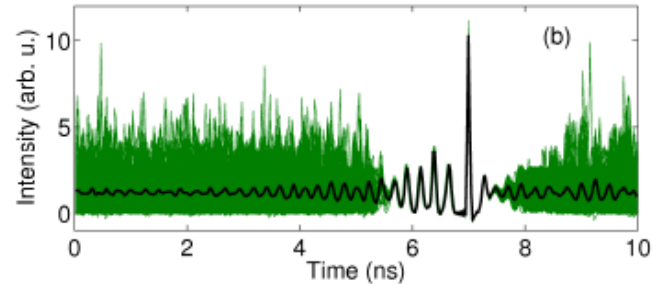


Deterministic & stochastic  
simulations ( $\beta_{sp} = 10^{-4}$ )

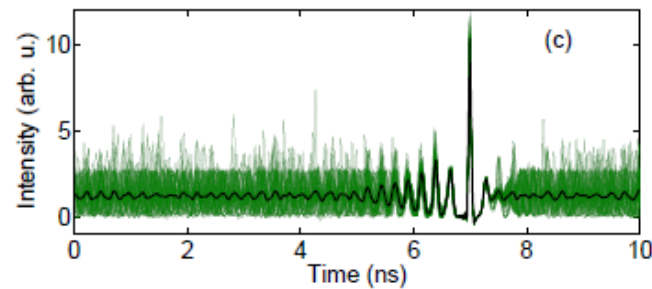
## Measured intensity time traces (500 RWs)



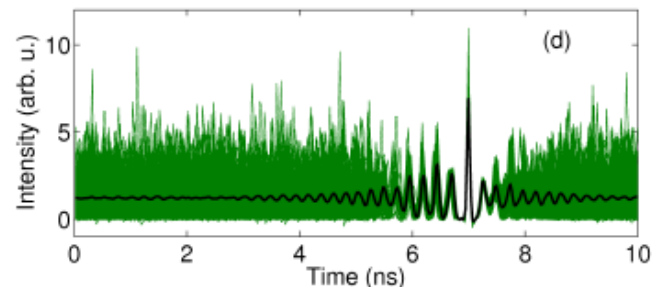
## Simulated time traces



$\langle H \rangle + 8\sigma$   
 $\beta_{sp} = 0$   
459 RWs



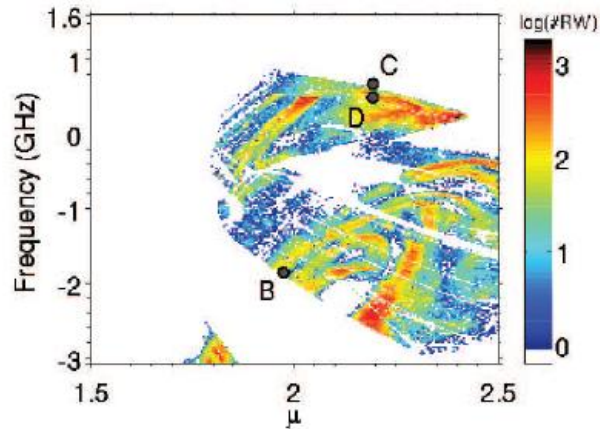
$\langle H \rangle + 8\sigma$   
 $\beta_{sp} = 10^{-2}$   
53 RWs



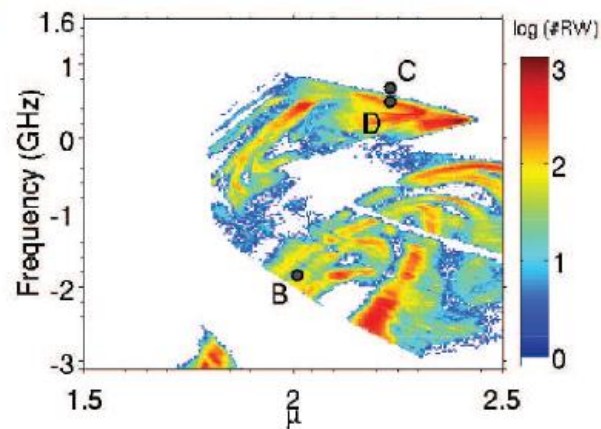
$\langle H \rangle + 4\sigma$



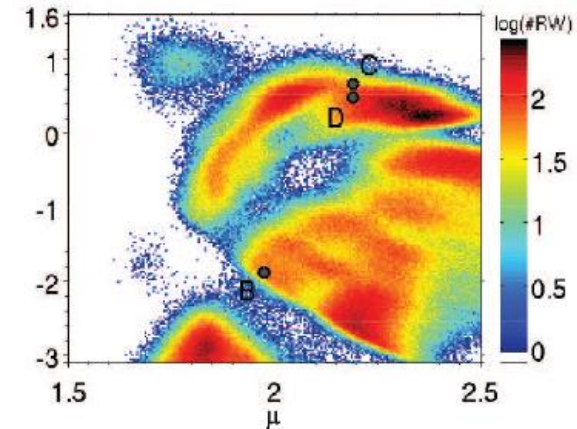
# Influence of spontaneous emission noise



$$\beta_{sp} = 0$$



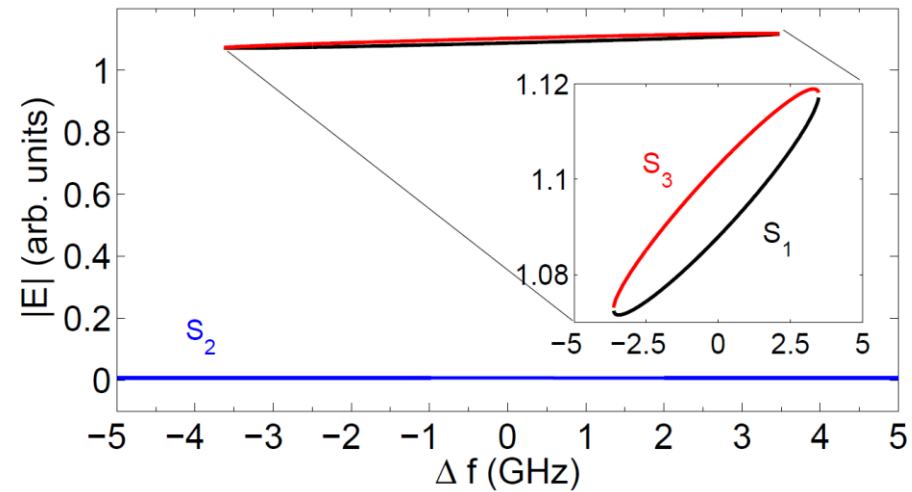
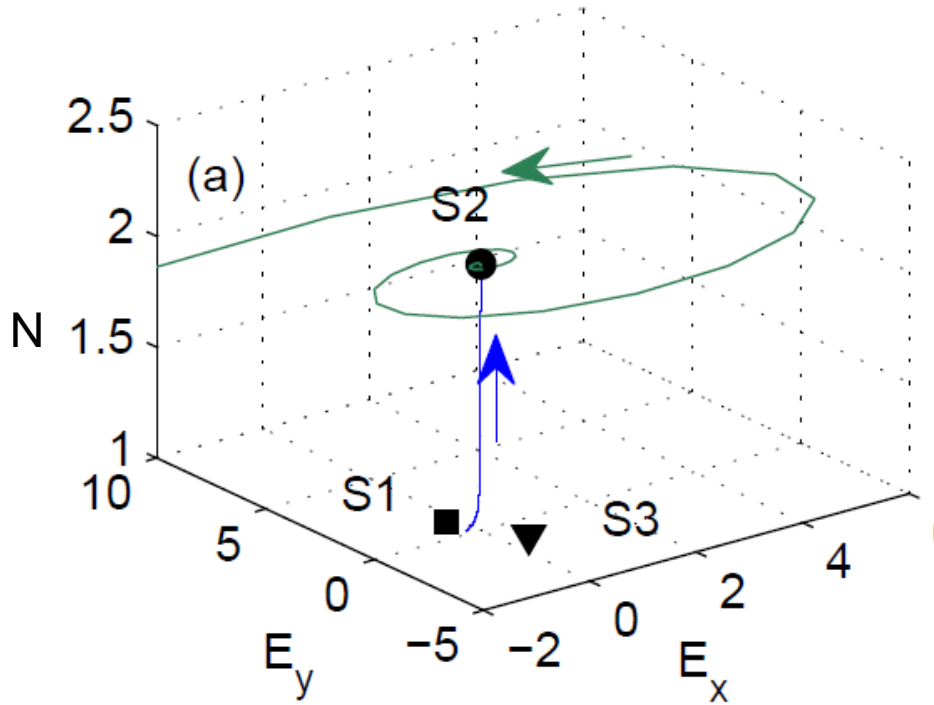
$$\beta_{sp} = 10^{-4}$$

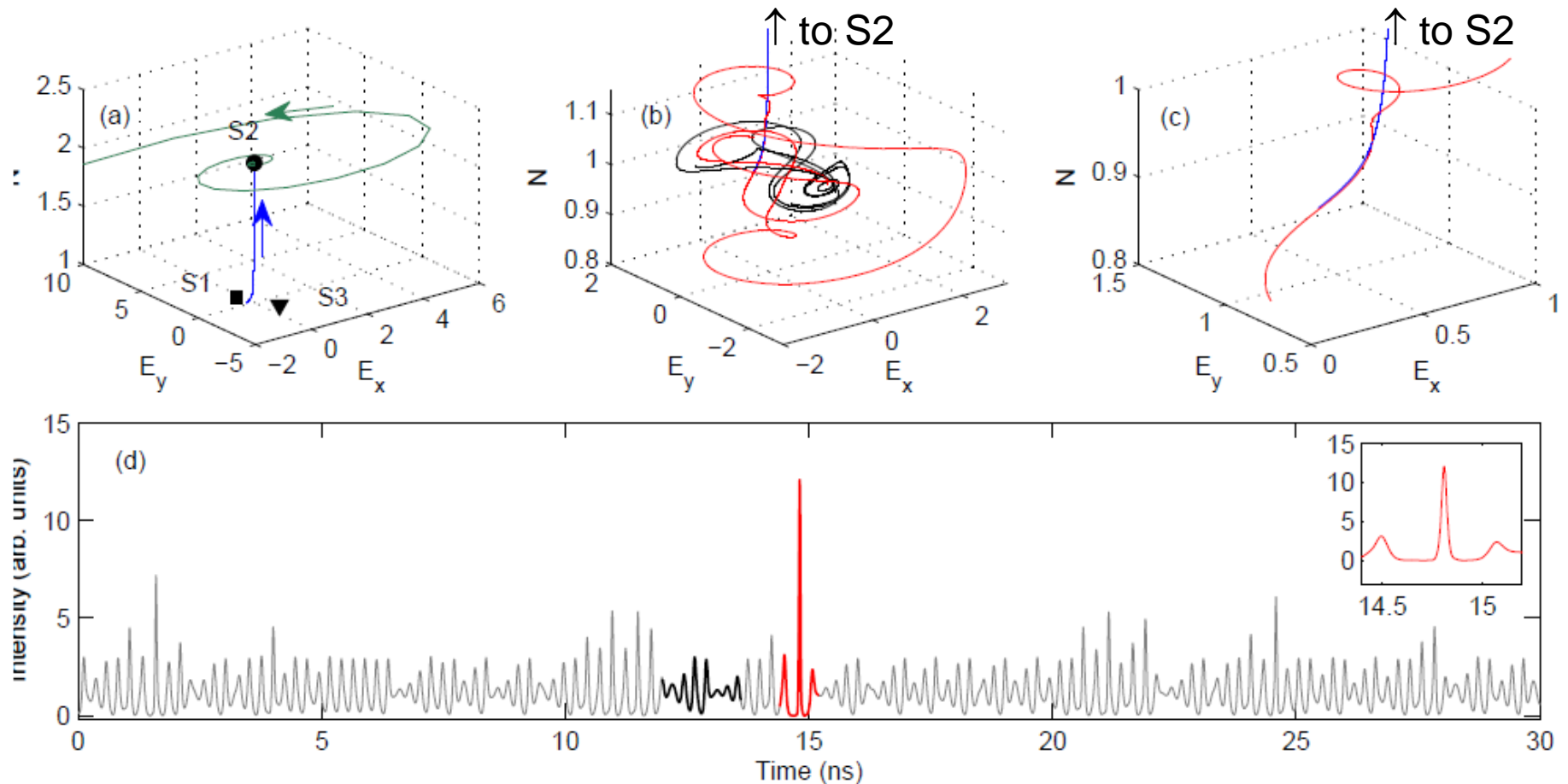


$$\beta_{sp} = 10^{-2}$$

# What triggers a RW pulse?

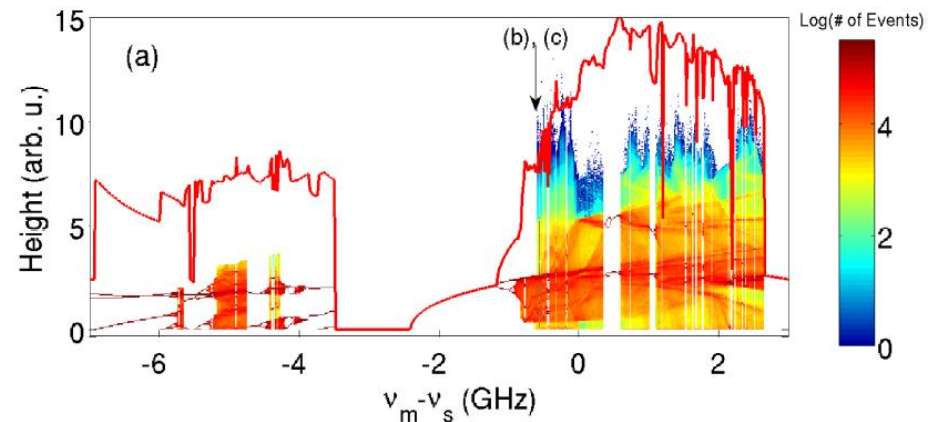
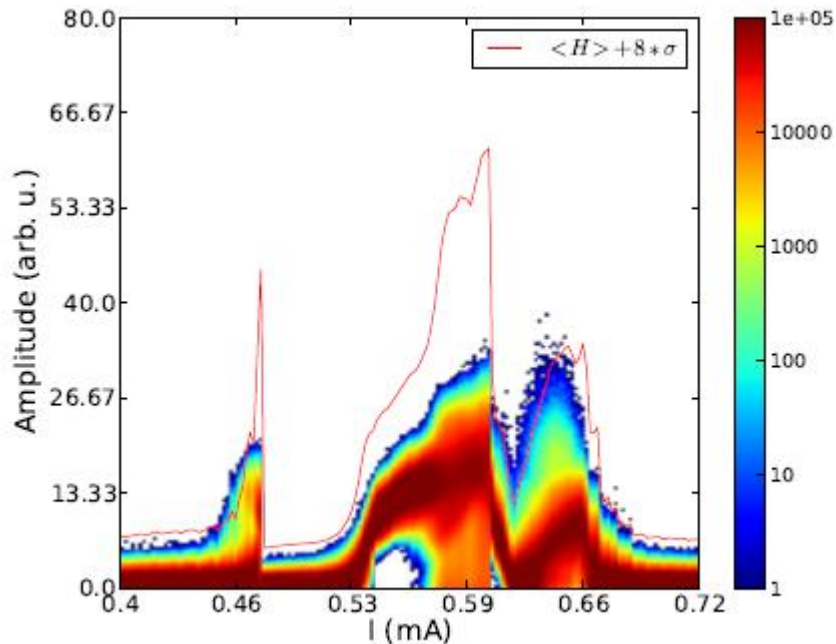
## Fixed points in the phase space





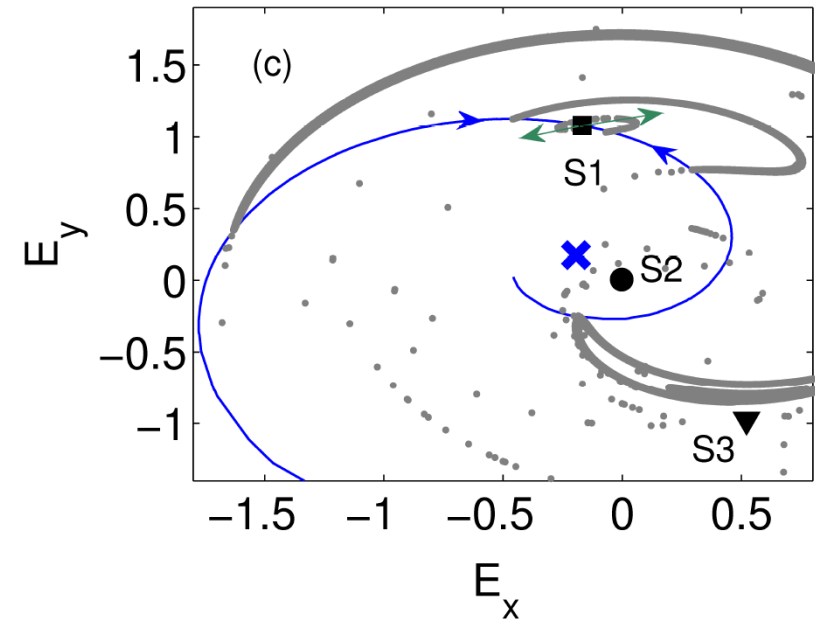
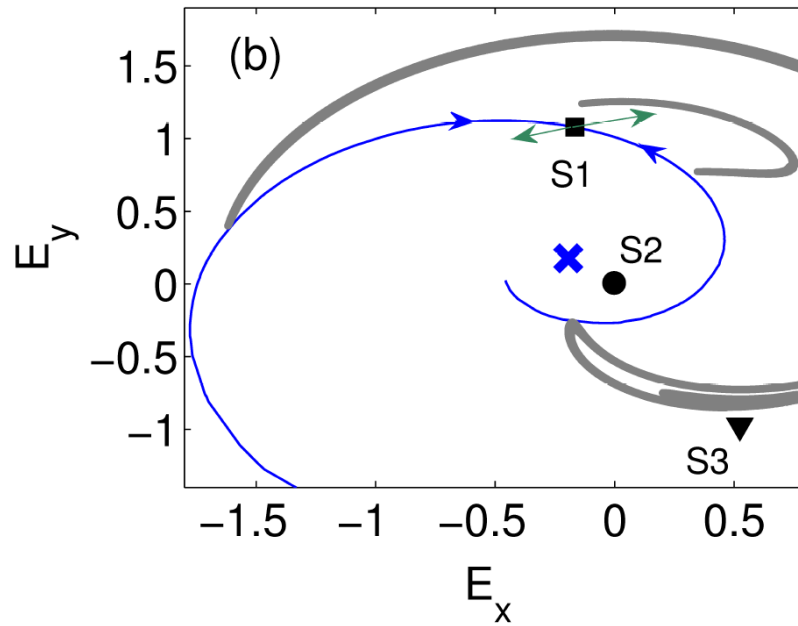
A RW is triggered whenever the trajectory closely approaches **the unstable manifold of S2**.

# chaos with RWs and chaos without them



Experimental data: amplitude of  
the intensity pulses

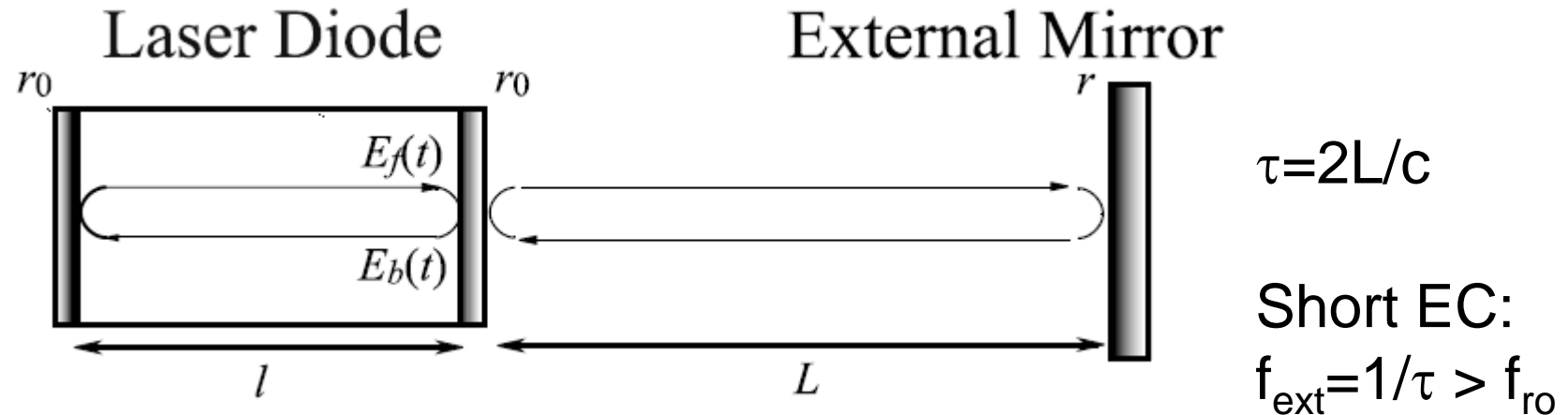
# Why chaos with RWs and chaos without them?



An external crises-like process enables access to the phase space region where the stable manifold of  $S2$  ( $x$ ) is.



# Dynamics with optical feedback from a short external cavity



$$dE/ds = (1 + i\alpha)NE(s) + \eta e^{-i\omega\theta} E(s - \theta) + \beta\xi,$$

$$TdN/ds = J - N - (1 + 2N)|E(s)|^2.$$

Typical parameter values:

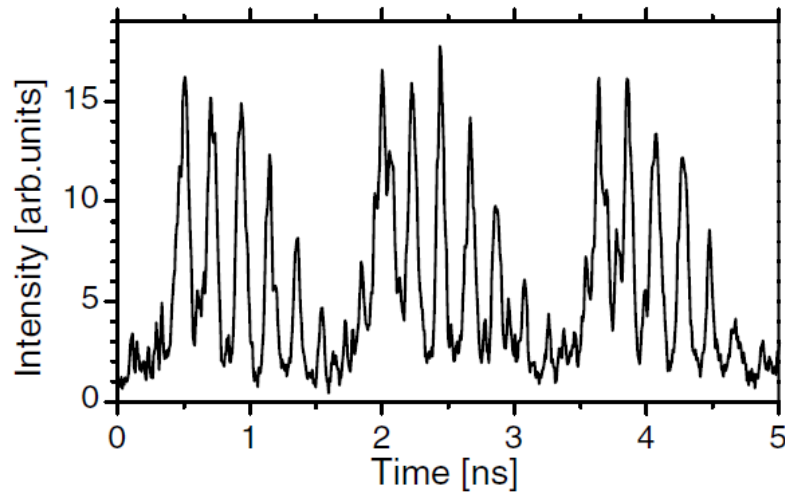
**$\alpha = 5$ ,  $T = 1710$ ,  $\theta = 70$ ,  $J = 1.155$**

$$s = t/\tau_p \quad \theta = \tau/\tau_p$$

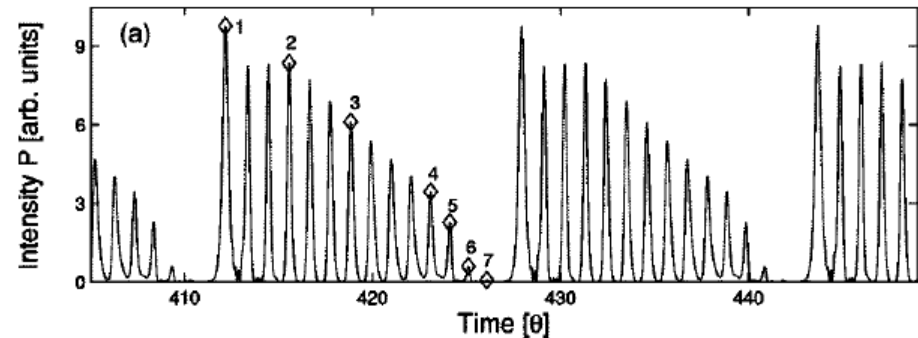
$$T_{RO} = \pi\sqrt{2T/J} = 171$$

# Regular Pulse Packages (RPPs)

## ■ Experiments



## ■ Simulations

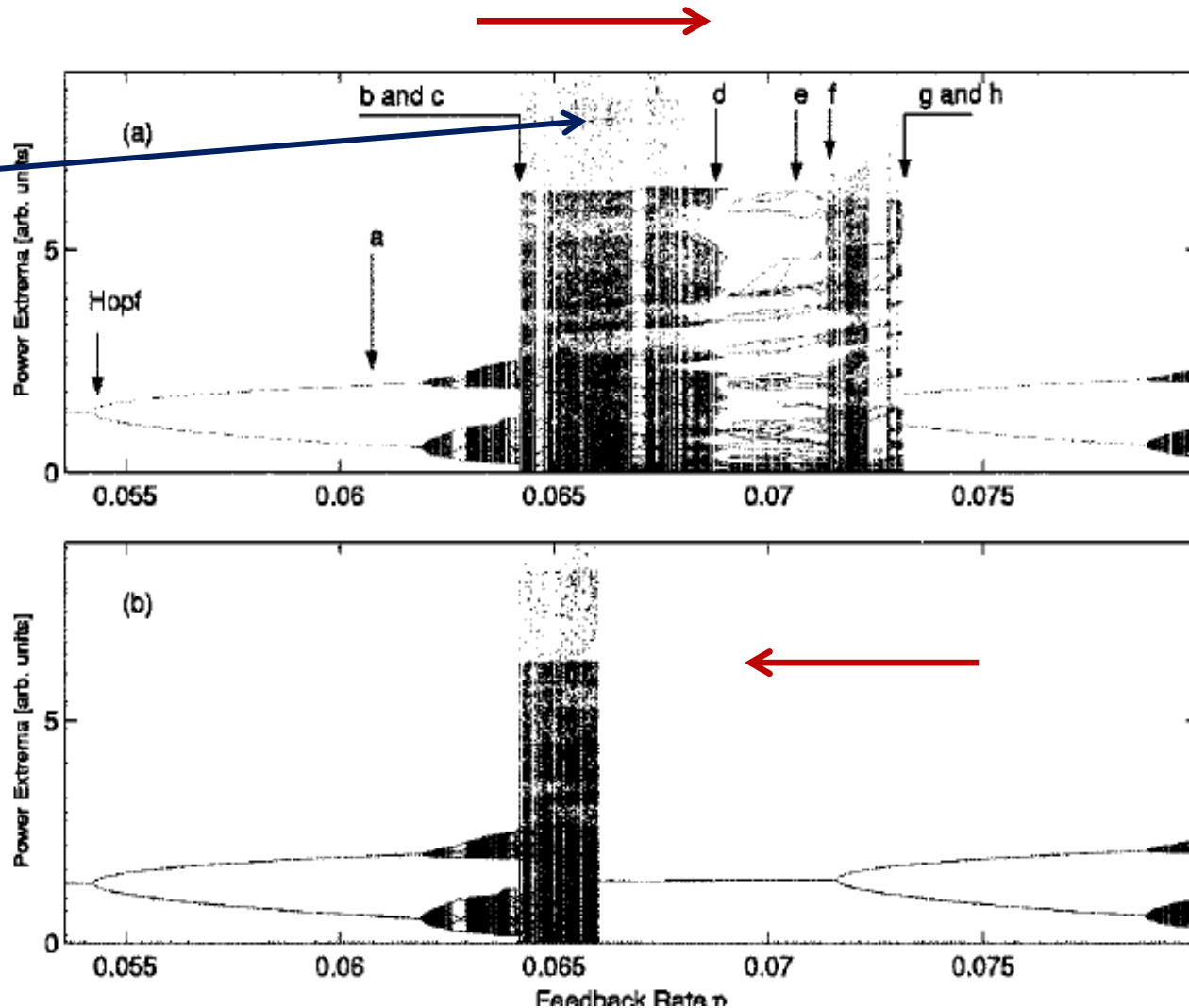


T. Heil et al, PRL 87, 243901 (2001)

A. Tabaka, et al. PRE 70, 036211 (2004)

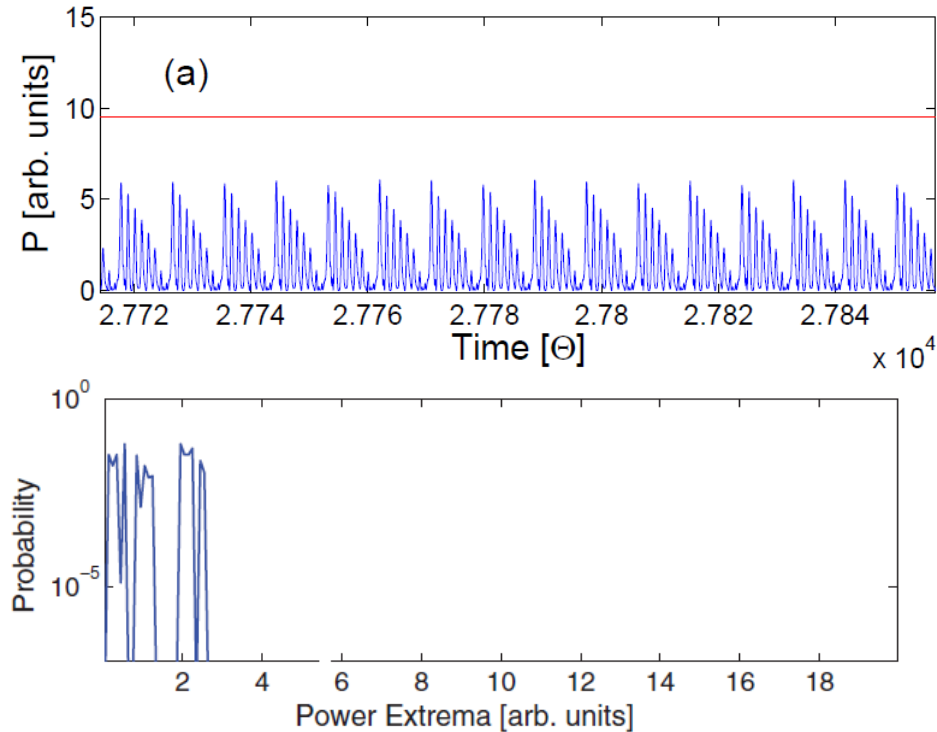
# Numerical bifurcation diagram

EPs?

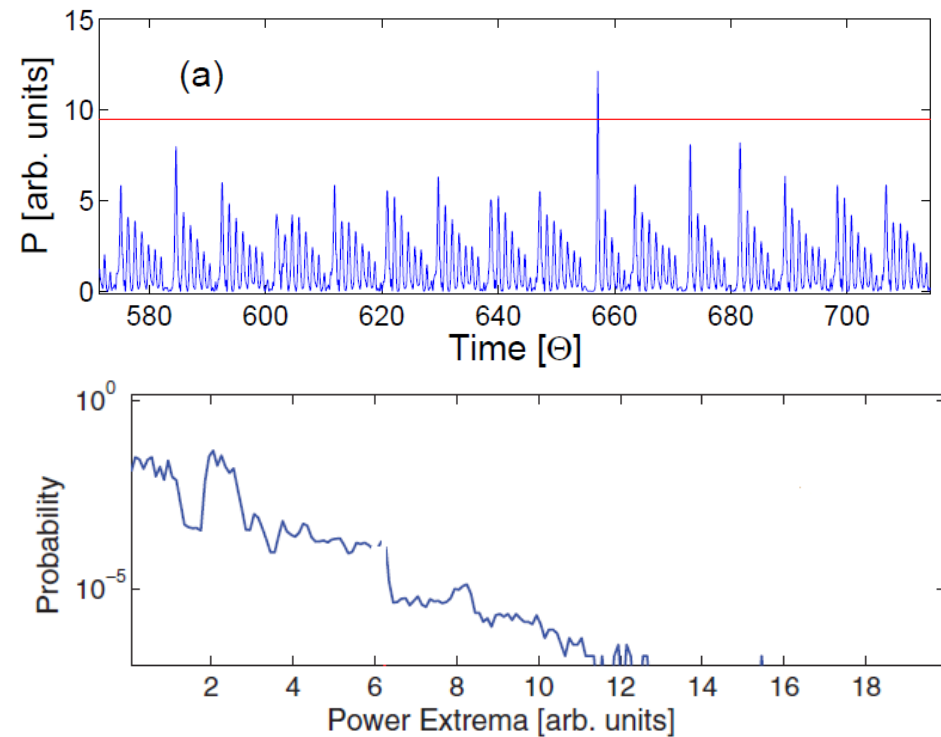


A. Tabaka, et al. PRE 70, 036211 (2004)

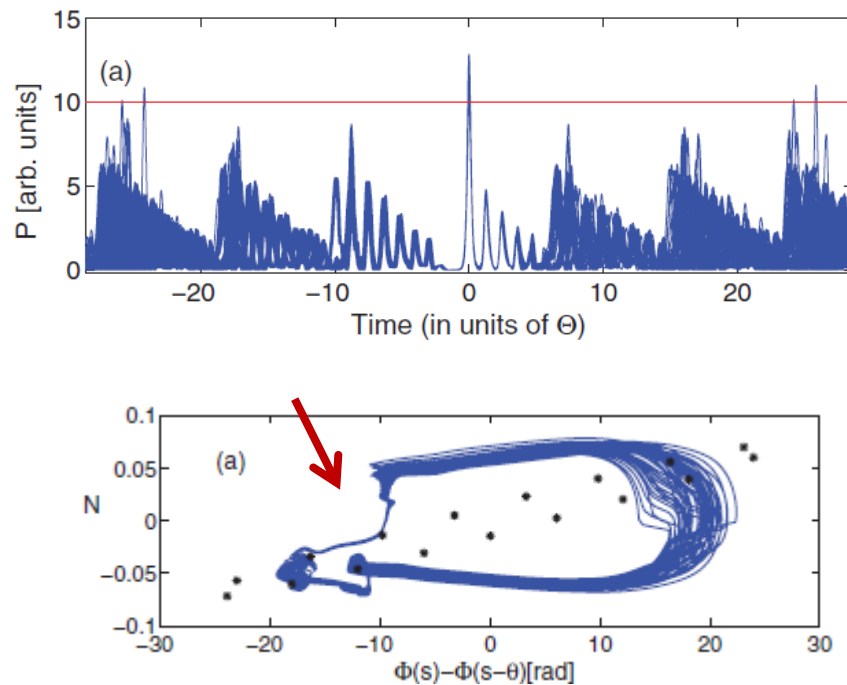
## Regular pulse packages



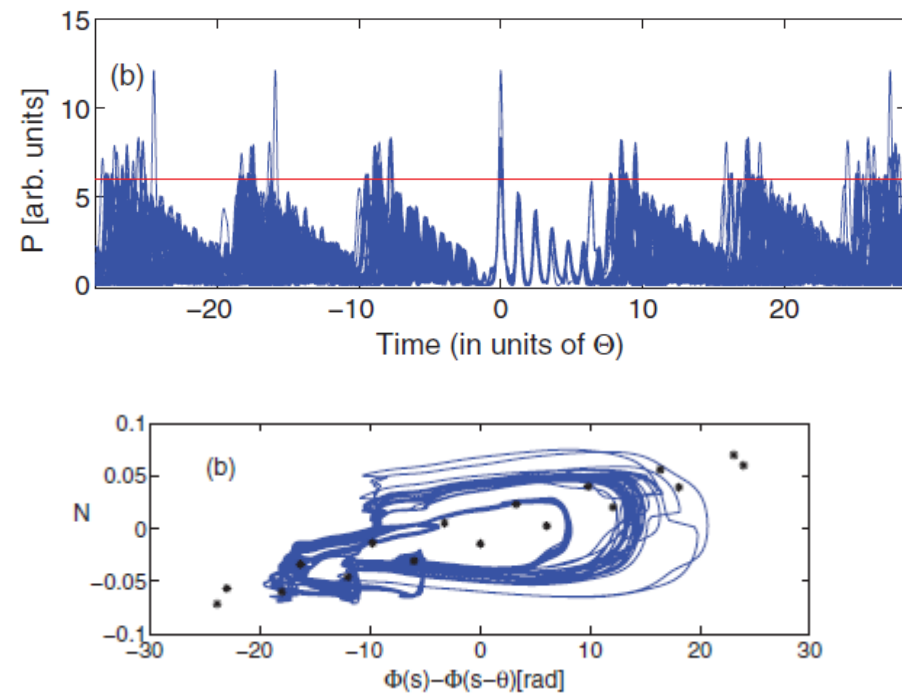
## Extreme pulses



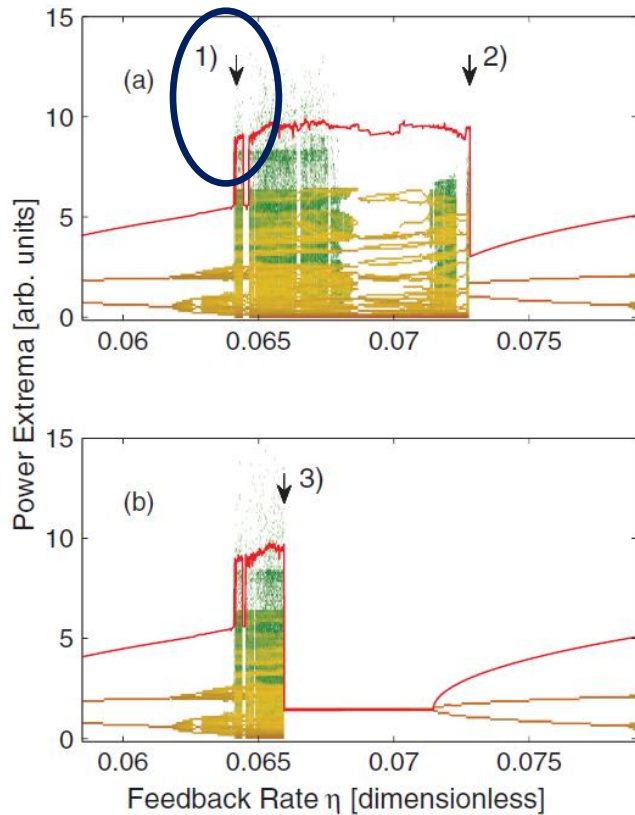
## Using a high threshold



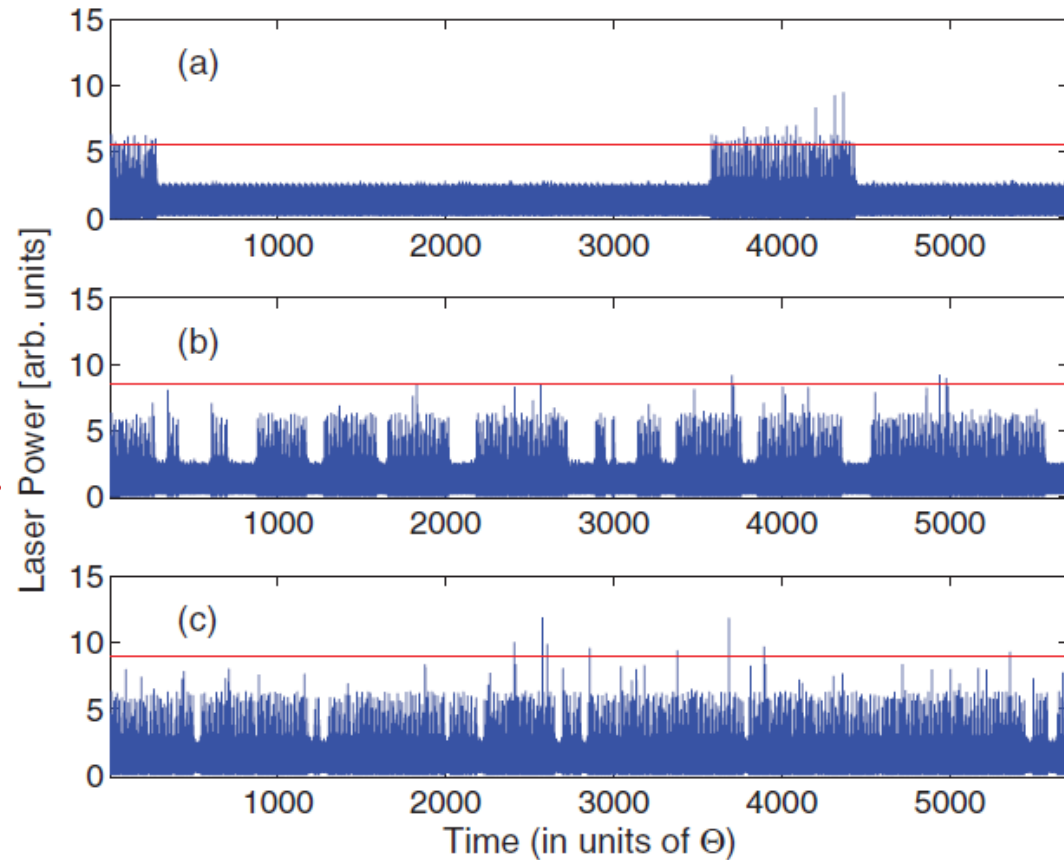
## Using a lower threshold



# Deterministic intermittency

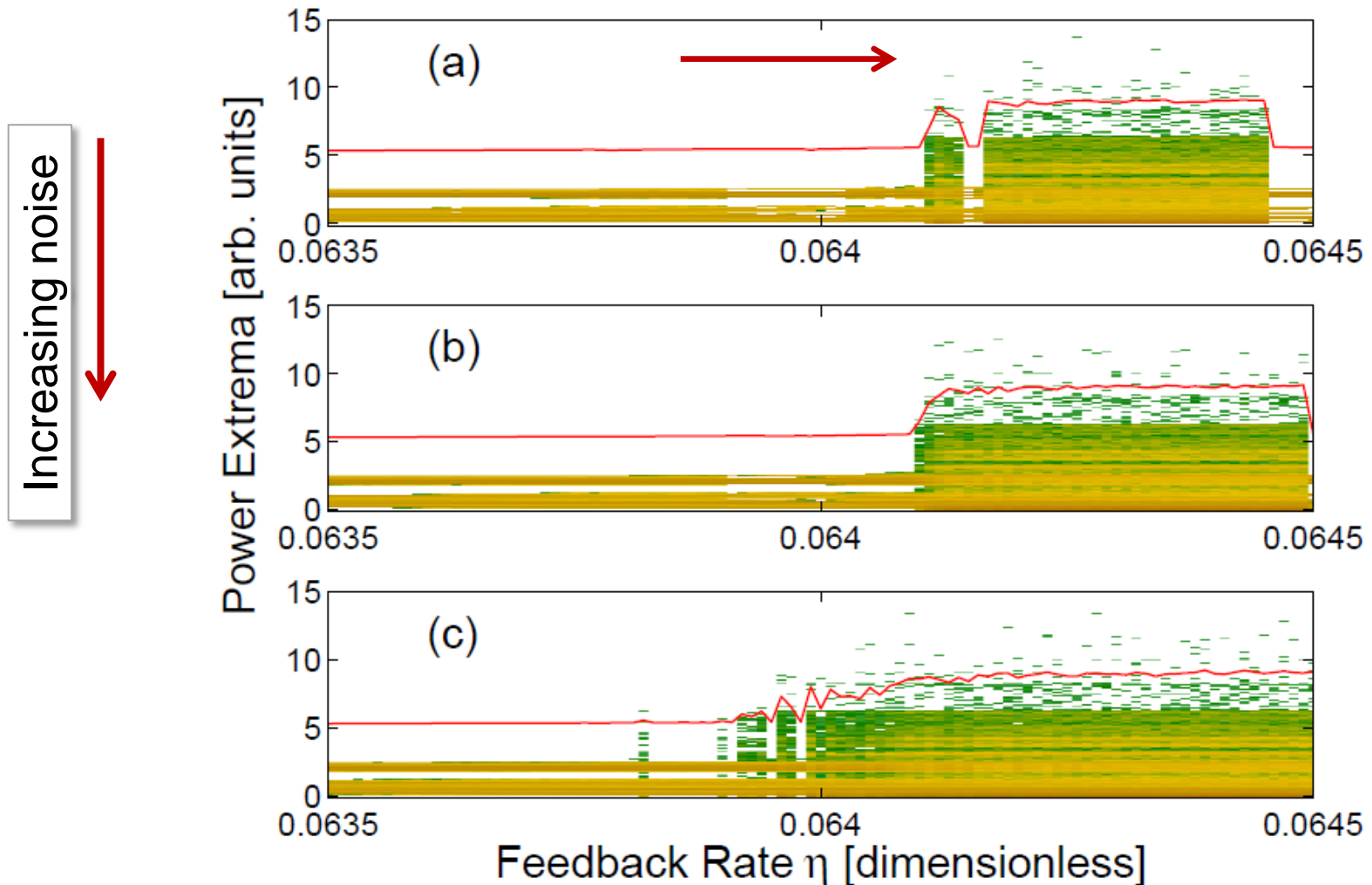


Increasing  $\eta$

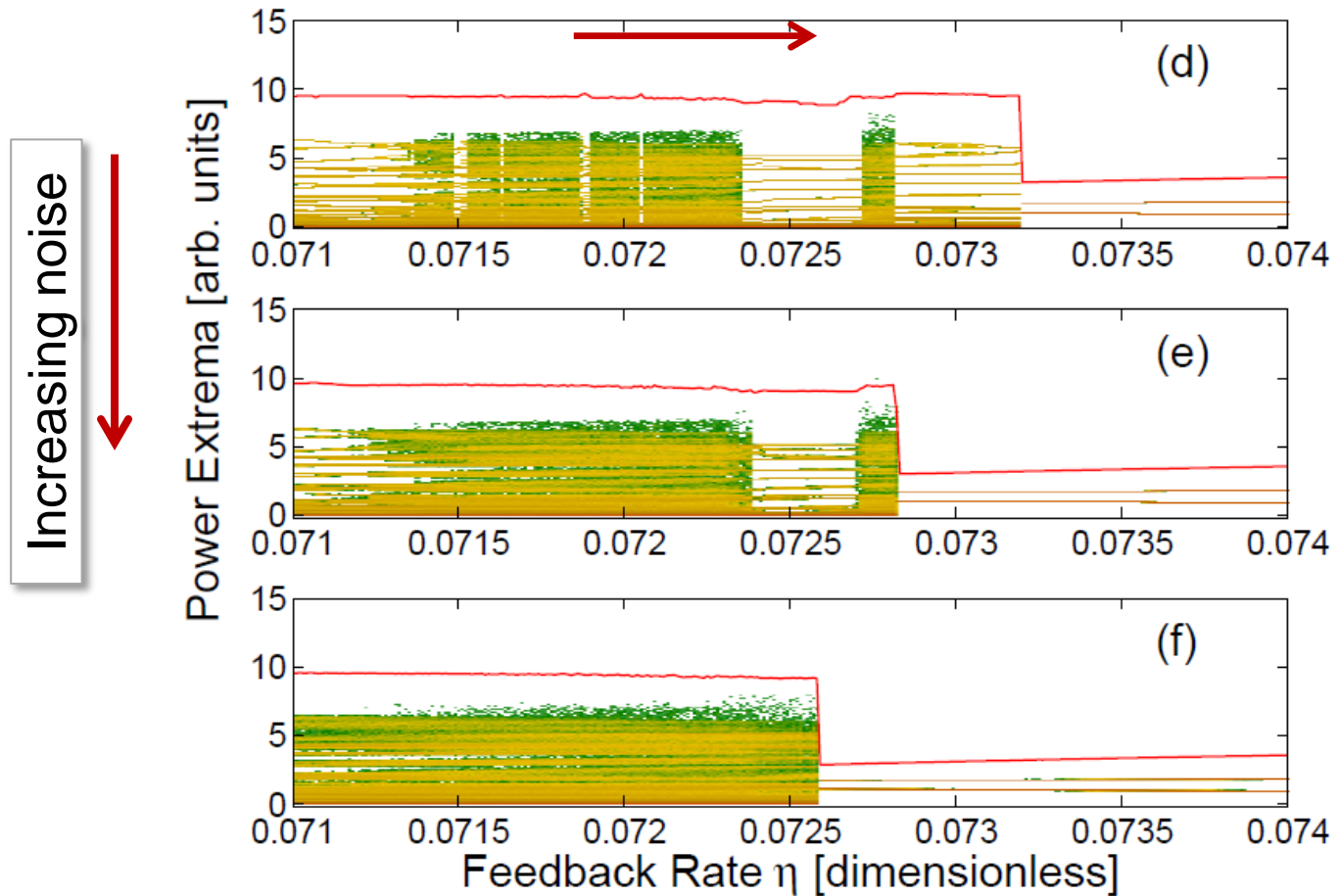




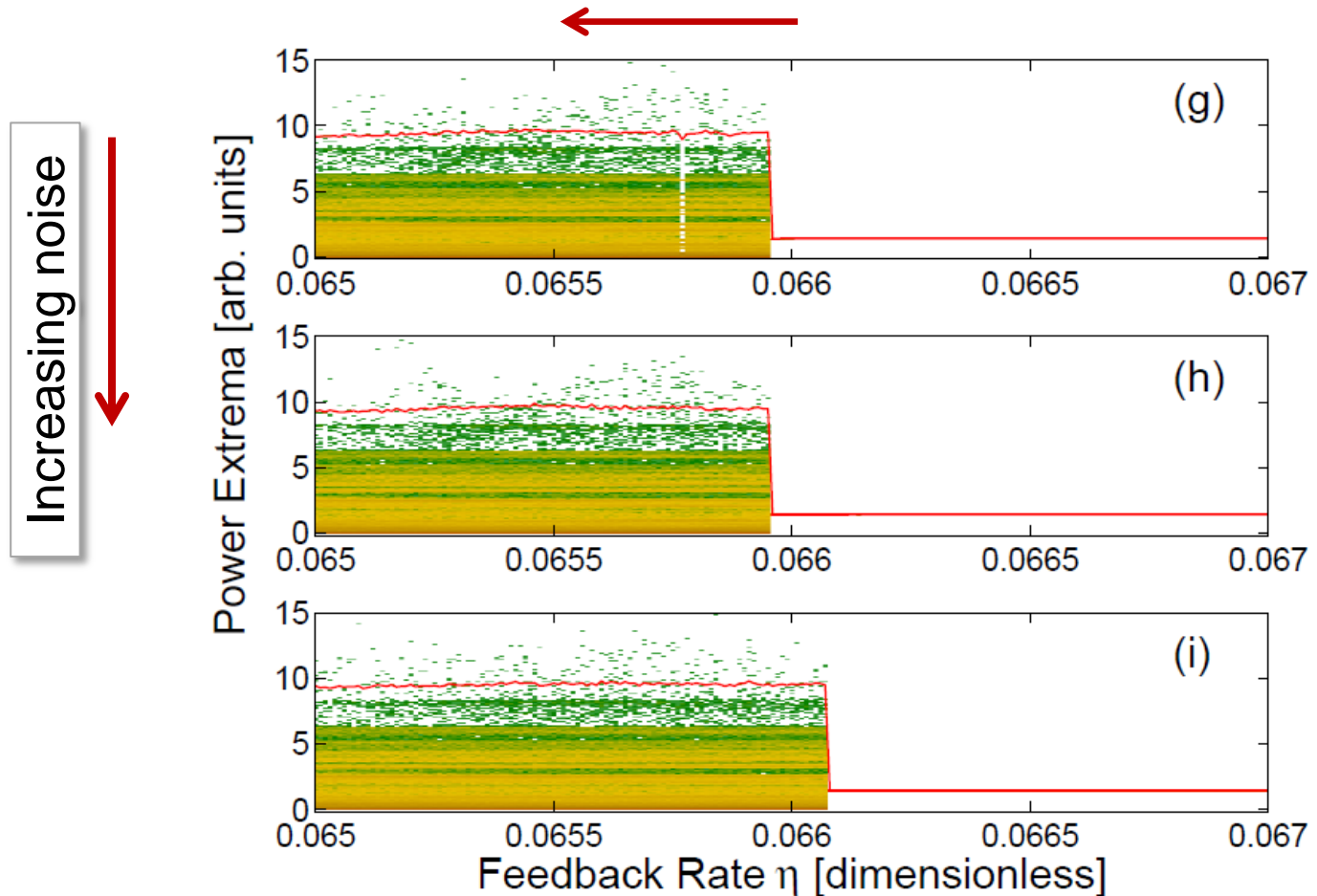
## ■ Transition 1: noise induced EPs



- Transition 2: noise advances the switching



- Transition 3: noise advances the switching



## Optically injected semiconductor lasers:

- Intensity pulses characterized by long-tailed histograms; giant rare pulses interpreted as Rogue Waves.
- Different types of chaos identified: without and with rogue waves.
- **Origin** of RWs: deterministic. An external crises-like process enables access to the region in phase space where RWs can be triggered.
- **Predictability**: in our system RWs can be predicted with some anticipation.
- **Control**: noise strongly affects their probability of occurrence.

## External-cavity lasers:

- Similar results, intermittency is the route to extreme pulses.

THANK YOU FOR YOUR ATTENTION !

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<http://www.fisica.edu.uy/~cris/>



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