



Spikes and extreme pulses in the dynamics of semiconductor lasers

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Experimental Techniques in Nonlinear Dynamics
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UNIVERSITAT POLITÈCNICA
DE CATALUNYA
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Campus d'Excel·lència Internacional



International
Year of Light
2015





Where are we?

1. Barcelona
2. Castelldefels
3. Igualada
4. Manresa
5. Mataró
6. Sant Cugat del Vallès
7. Terrassa
8. Vilanova i la Geltrú



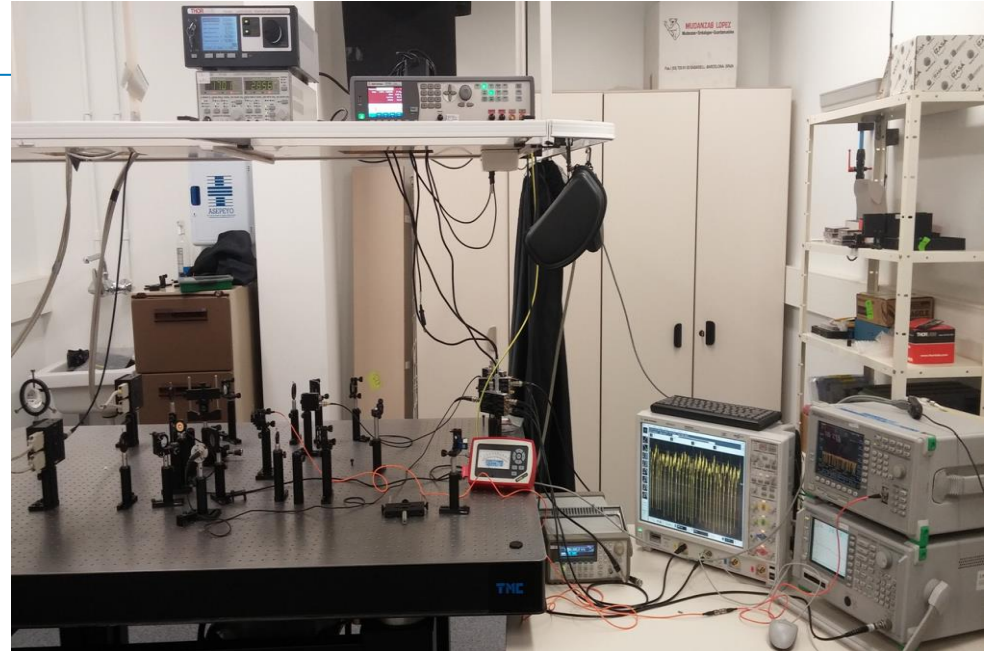
Viernes, 25 de septiembre de 2009 *Diari de Terrassa*



El edificio Gala centraliza grupos científicos consolidados y emergentes.

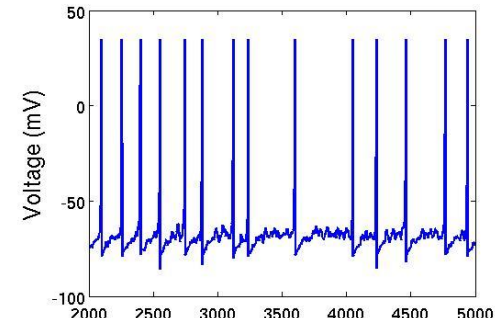
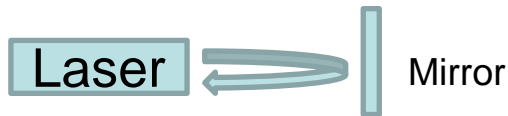


People and lab



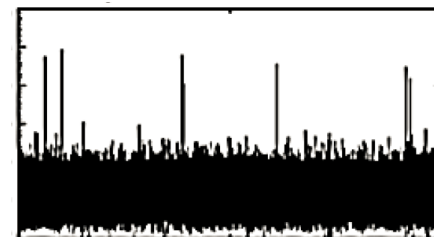
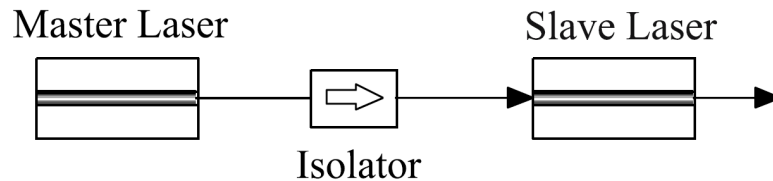
12 senior researchers,
2 postdoctoral researchers
10 phd students

■ Optical feedback



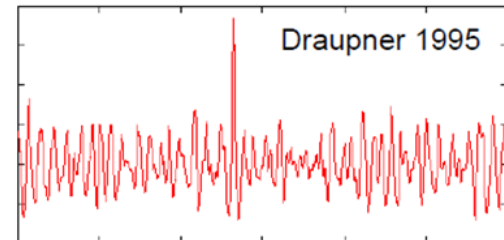
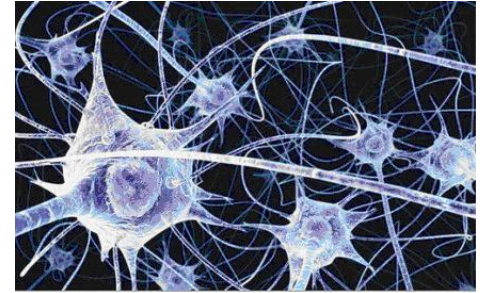
- Are optical spikes similar to neuronal spikes ?

■ Optical injection



- Extreme pulses
 - Can they be predicted?
 - Can they be controlled?

- Optical Neurons open the possibility of developing novel photonic neuro-inspired information processing systems: ultrafast (**ms vs ns- μ s**).
- Extreme events occur in many complex systems
 - “big data” approach: the laser setup allows for recording large amounts of data,
 - Useful for testing methods of prediction (“early warning signals”) and control.



- Introduction
 - Nonlinear dynamics of semiconductor lasers
 - Method of time series analysis

- Part 1: optical neurons

- Part 2: extreme pulses

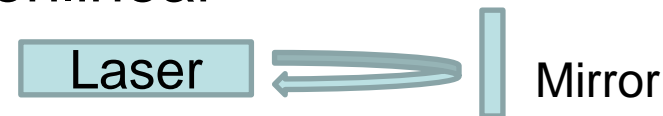
- Widely used in:

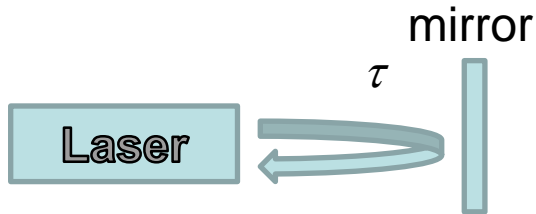
- Communications
- Data storage (CDs, DVDs ...)
- Barcode scanners, laser printers, computer mice
- Life sciences (imaging, sensing ...)
- Etc.



- Optical feedback can improve the emission characteristics but it can also induce nonlinear dynamics:

- Irregular power dropouts
- Chaotic emission





$|E|^2 \sim$ photon number (output intensity)

$N \sim$ number of carriers (electron-holes)

$$\frac{dE}{dt} = \frac{1}{2\tau_p} (1 + i\alpha)(G - 1)E + \underbrace{\eta E(t - \tau)e^{-i\omega_0\tau}}_{\text{feedback}} + \underbrace{\sqrt{\beta_{sp}}\xi}_{\text{noise}}$$

$$\frac{dN}{dt} = \frac{1}{\tau_N} (\underbrace{\mu}_{\text{pump current}} - N - G|E|^2)$$

feedback noise

η = feedback strength

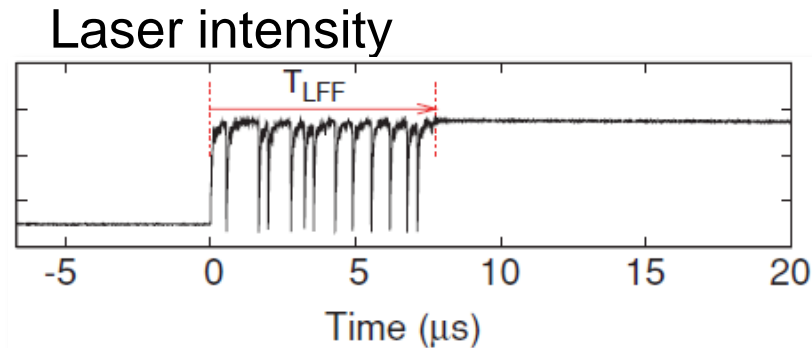
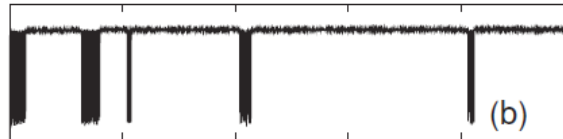
τ = feedback delay time

μ = pump current

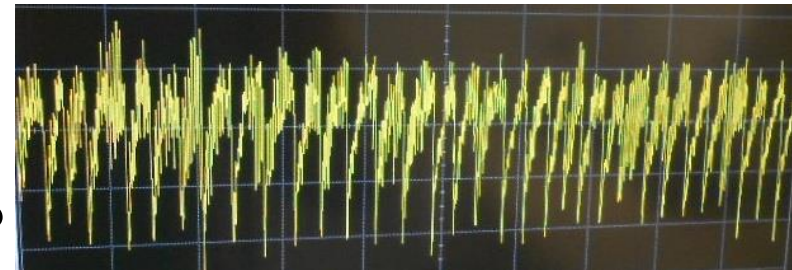
(control parameter)

Gain: $G = N / (1 + \varepsilon|E|^2)$

- In **deterministic** simulations: the spikes are **transient**.
- But in **stochastic** simulations: **bursts** of spikes.



- In the experiments:
which spikes are **noise-induced**
and which ones are **deterministic**?
- Can we infer signatures of determinism?
- Is there any **information** in the spike sequence?



How to extract information from optical spikes?

- Problem: we can measure only one variable (the laser output intensity).
- Also a problem: the detection system (photodiode, oscilloscope) has a finite *bandwidth* that gives limited temporal resolution.

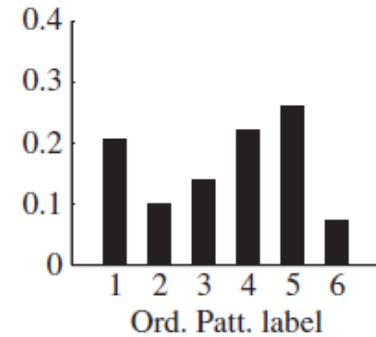
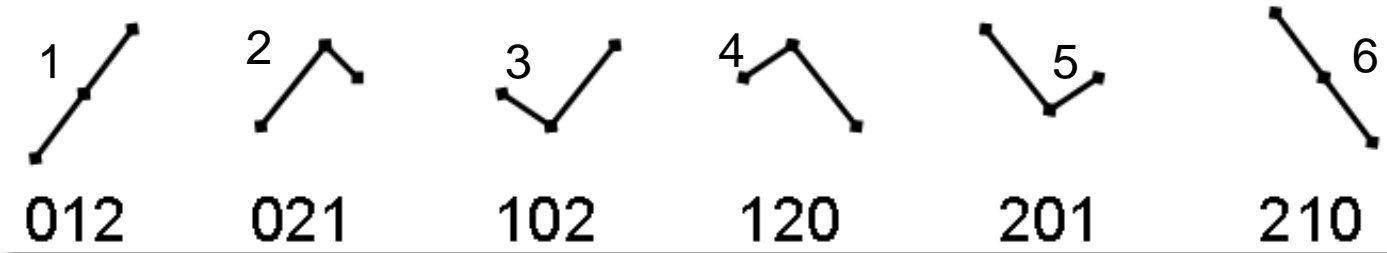


- Solution: **event-level description**. We analyze the sequence of **inter-spike-intervals (ISIs)**:

$$\Delta T_i = t_{i+1} - t_i$$

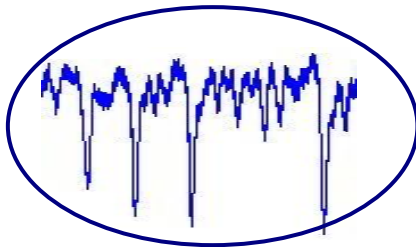
- Method of analysis: symbolic ordinal analysis

- $X = \{\dots \Delta T_i, \Delta T_{i+1}, \Delta T_{i+2}, \dots\}$ (**inter-spike-intervals**)

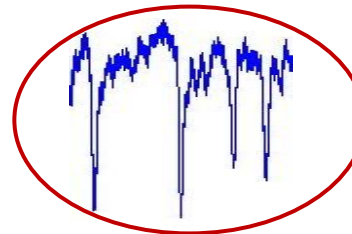


Example: (5, 1, 7) gives “102” because $1 < 5 < 7$

012



210

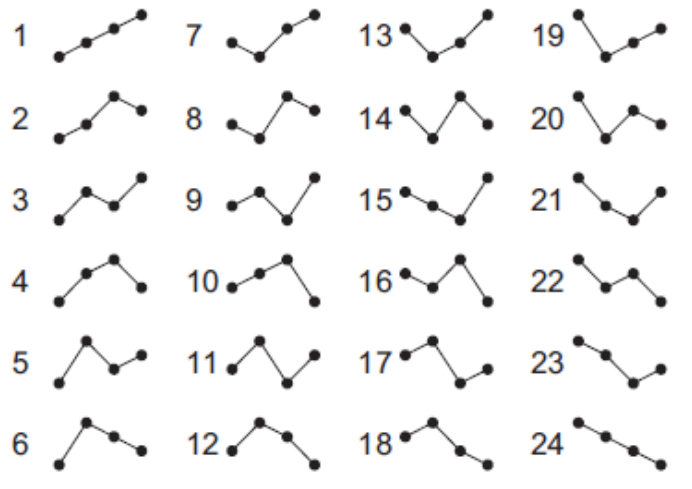


Random spikes
⇒ all patterns
are equally
probable

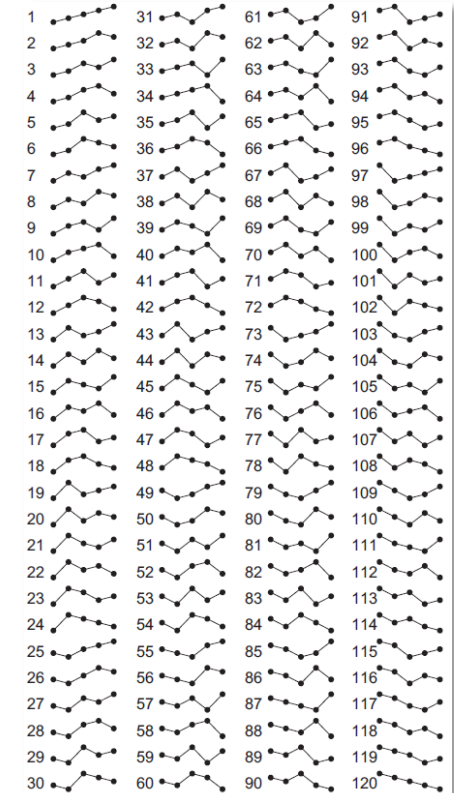
- Advantage: the OP **probabilities** uncover spike correlations.
- Drawback: we lose information (5,1,100) also gives “102”.

Number of possible ordinal patterns: D!

D=4



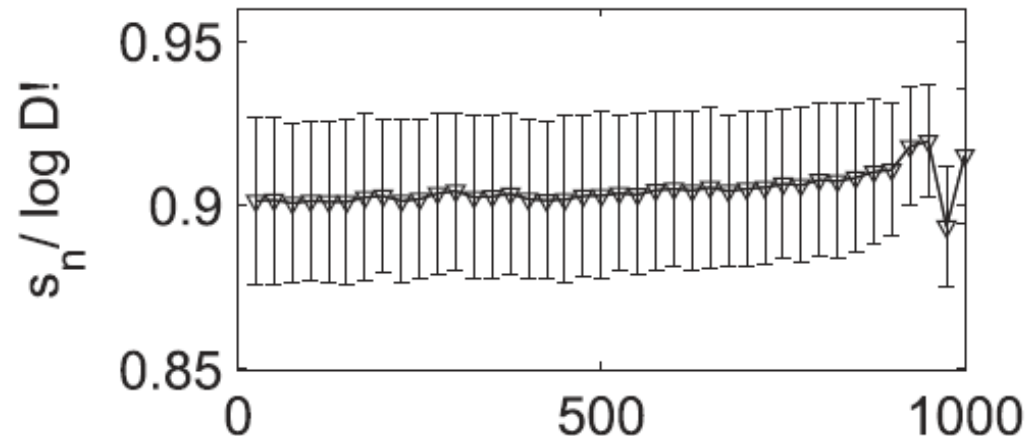
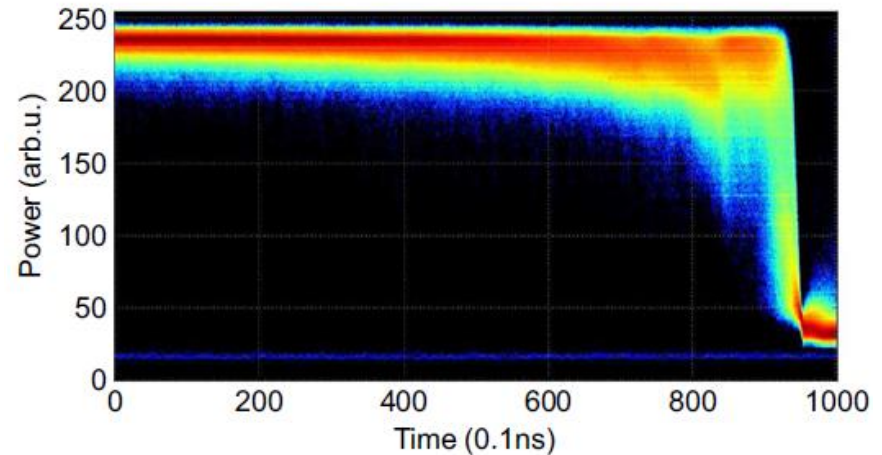
D=5



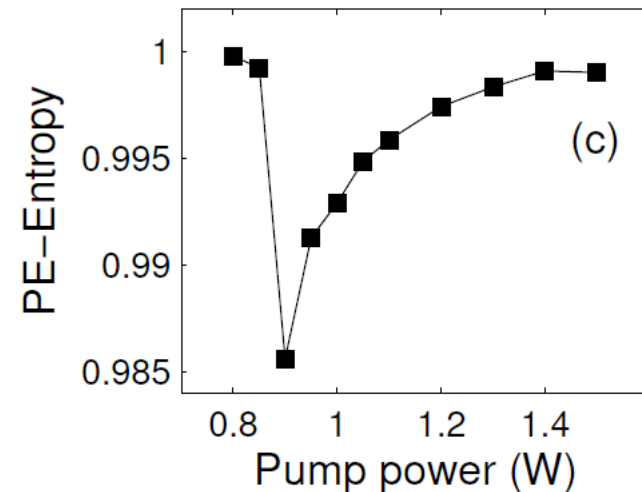
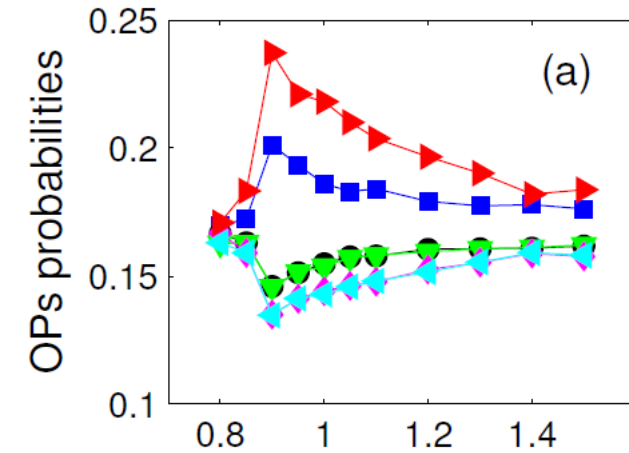
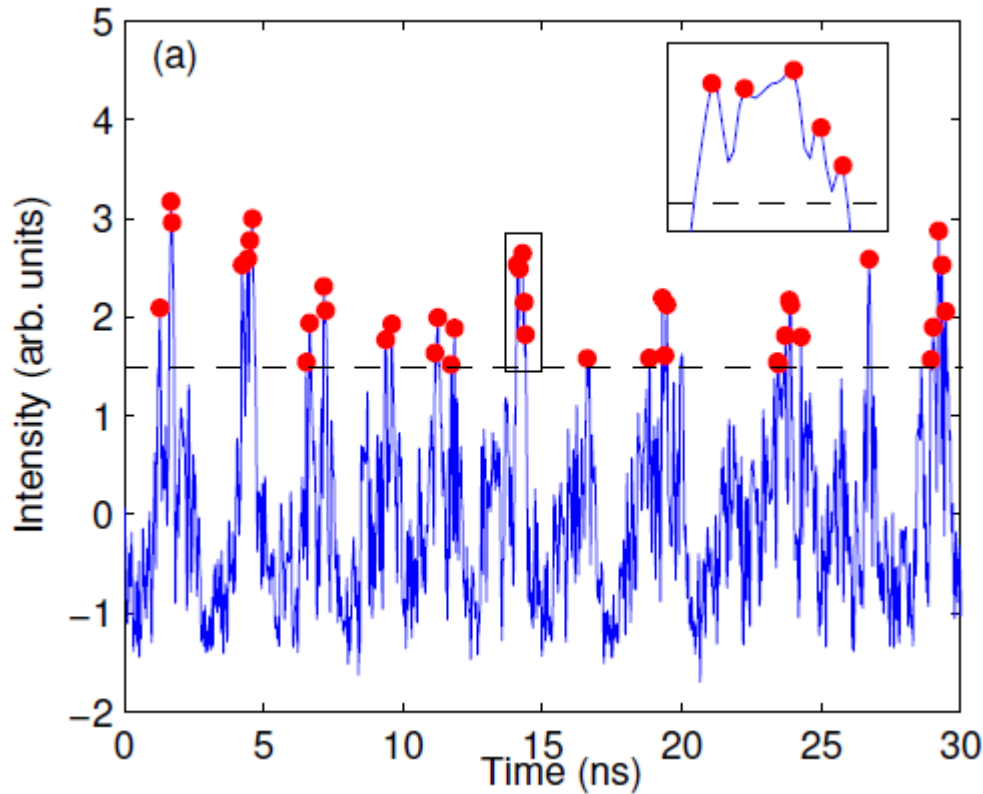
- How to select optimal D? depends on:
 - The length of the data.
 - The length of correlations in the data.

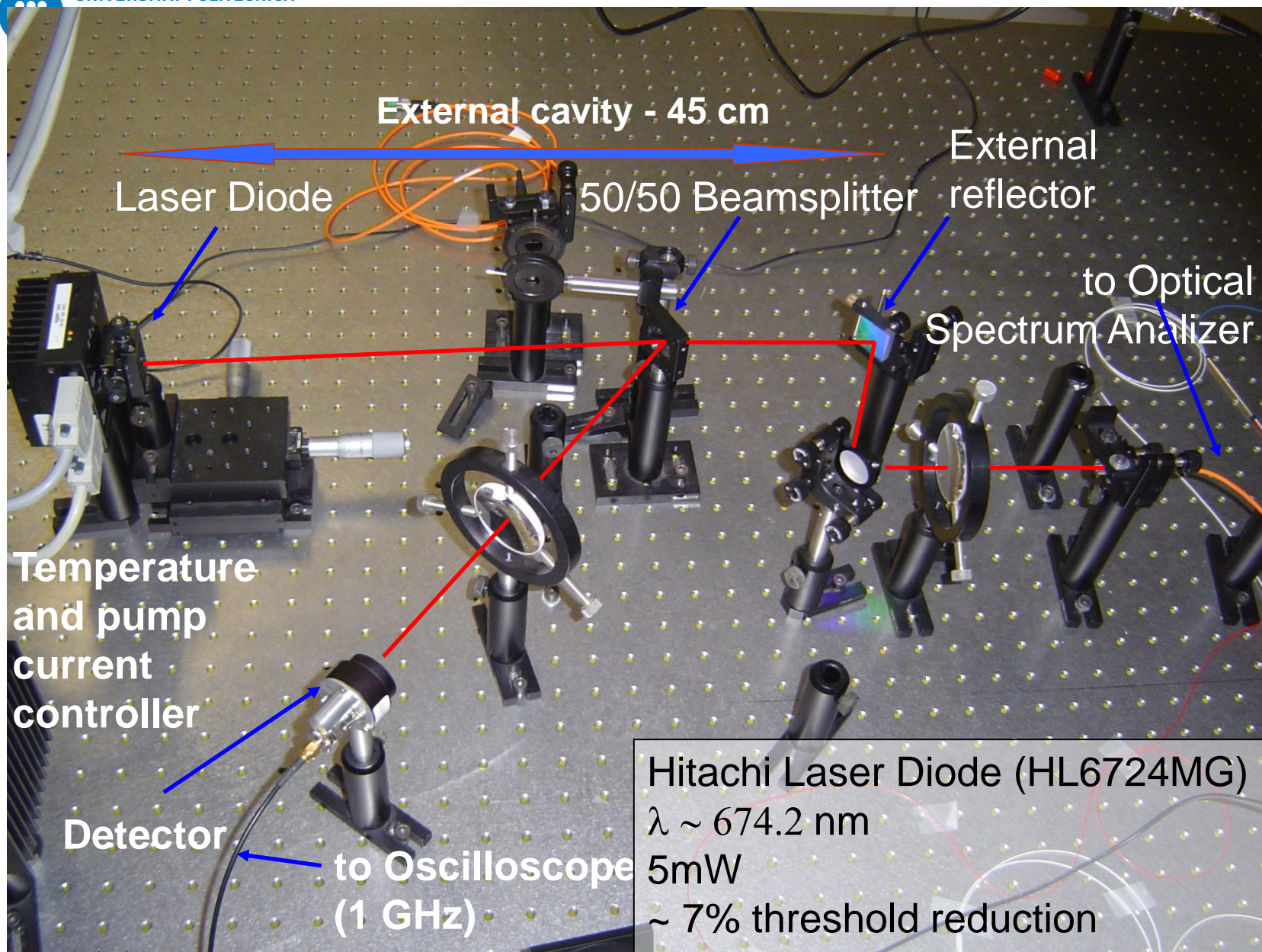
Ordinal analysis has been widely used to analyze the observed output signals of complex systems

- Example: abrupt polarization switching.
- An entropy measure provides an **early warning** of the transition.



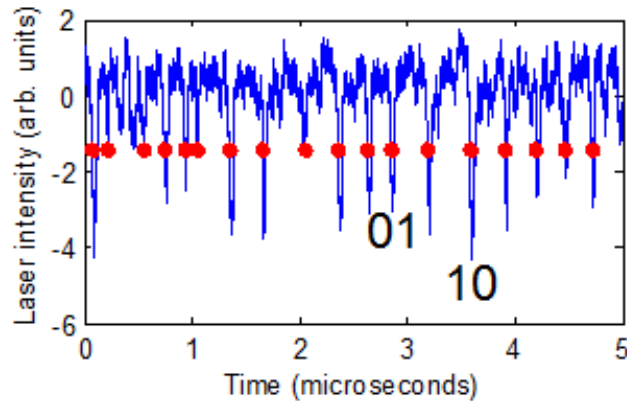
Transition laminar - turbulence in a fiber laser



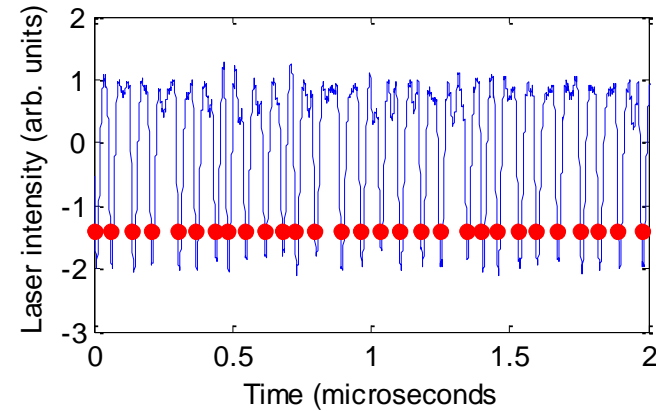


Spiking dynamics

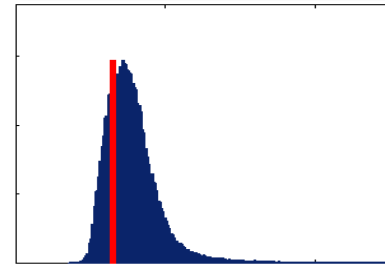
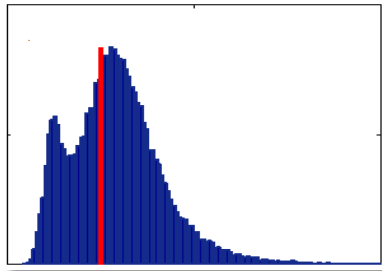
Low pump current



Higher pump current



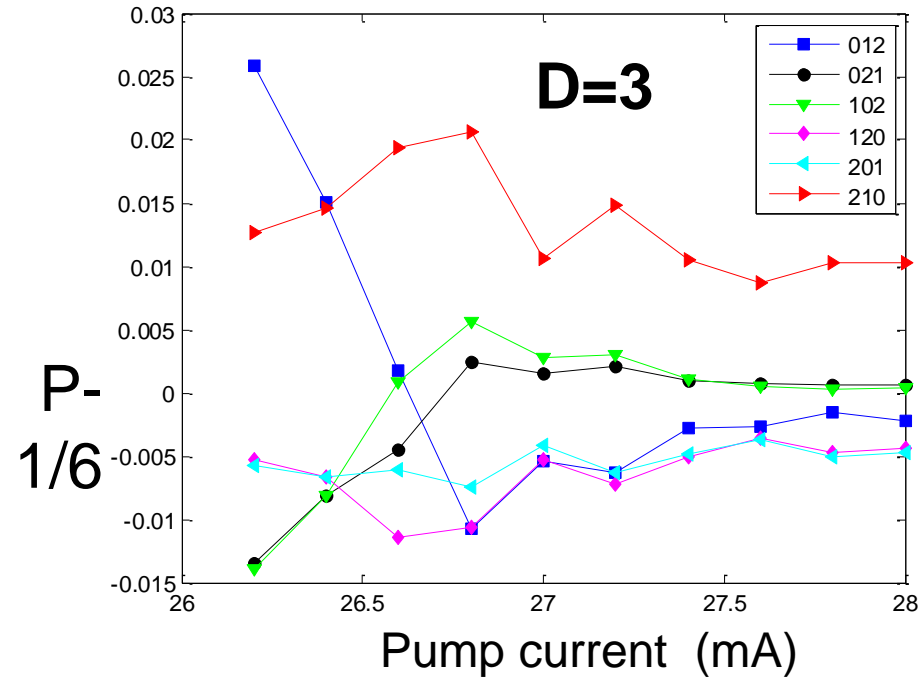
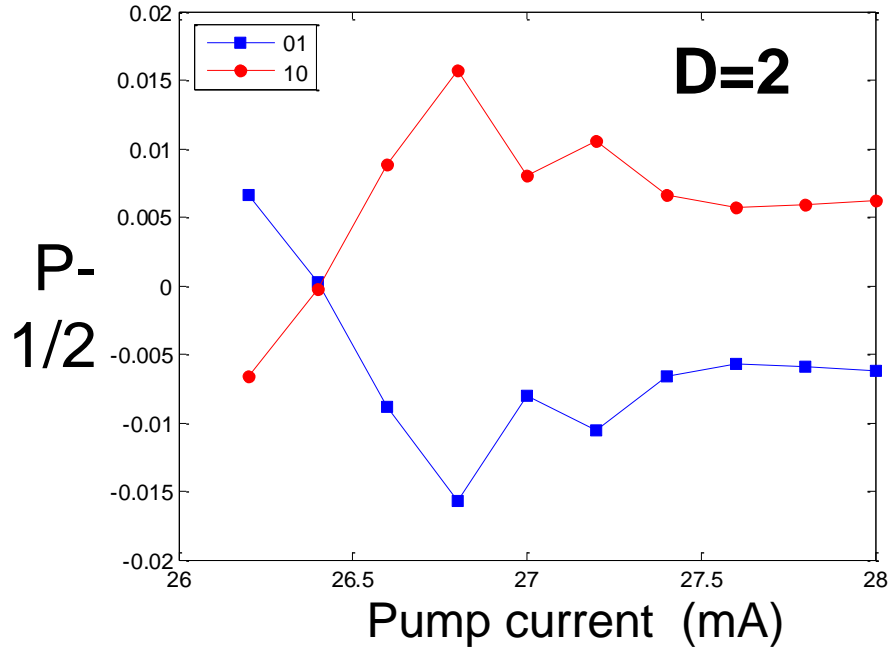
ISI
Histograms



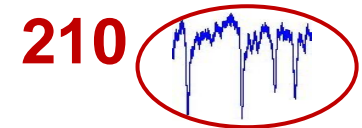
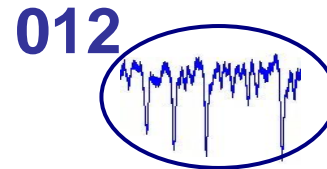
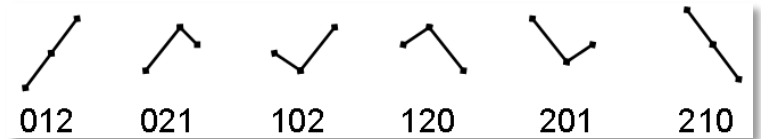
Probabilities of (**01**) and (**10**) uncover 3-spike correlations

Null Hypothesis: random spikes $\Rightarrow P(01) = P(10)$

Are optical spikes random?

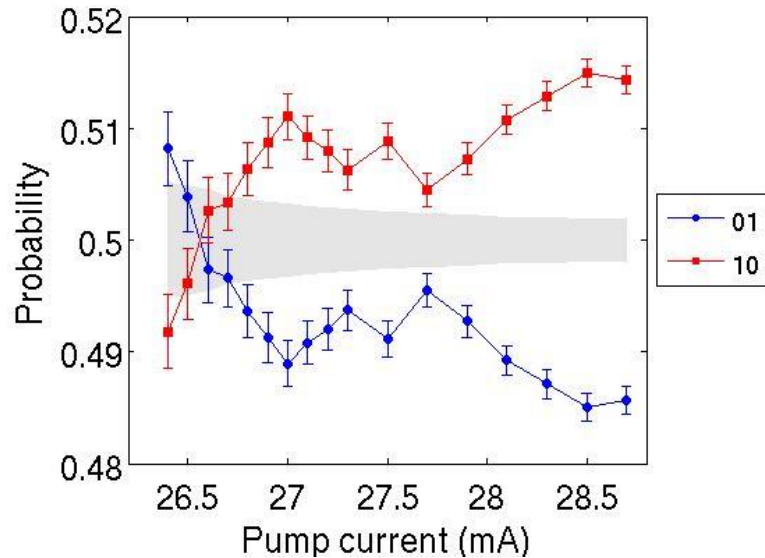


45000 - 220000 spikes

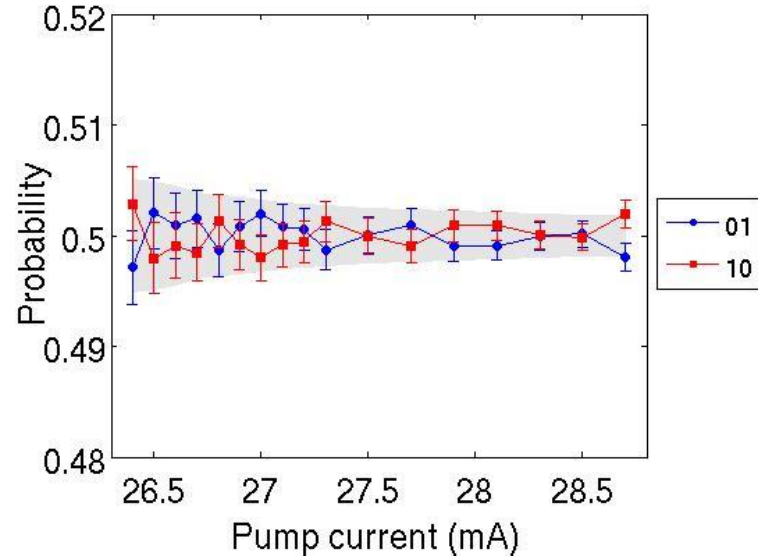


Are the results **significant**?

Recorded data



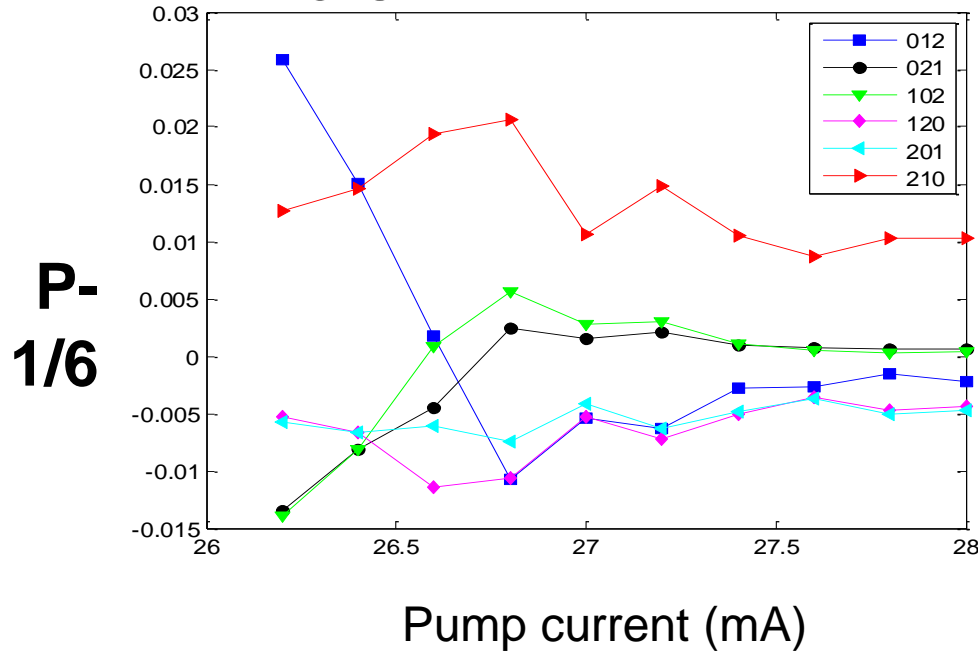
Surrogated data



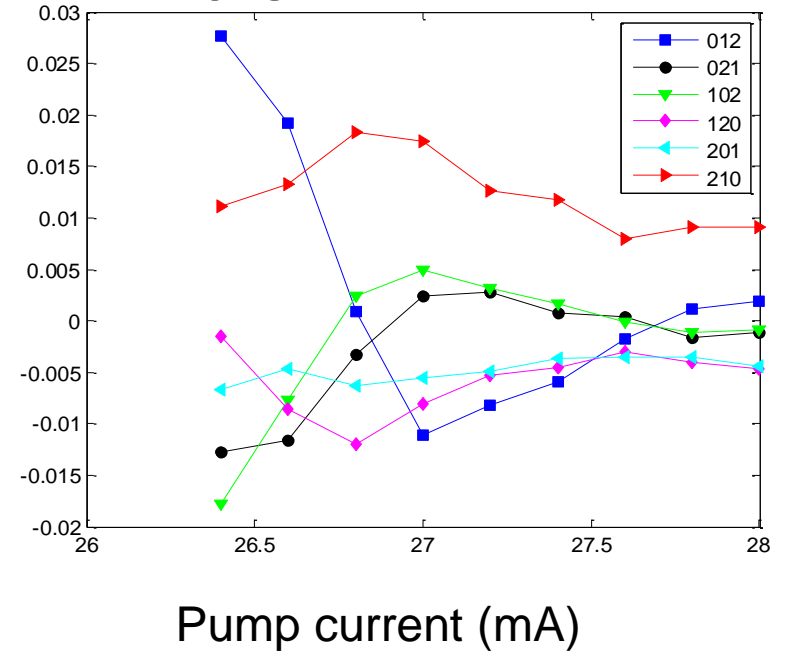
Error bars computed with a binomial test, gray region is consistent with N.H.

Ordinal analysis unveils new information

T=18 C

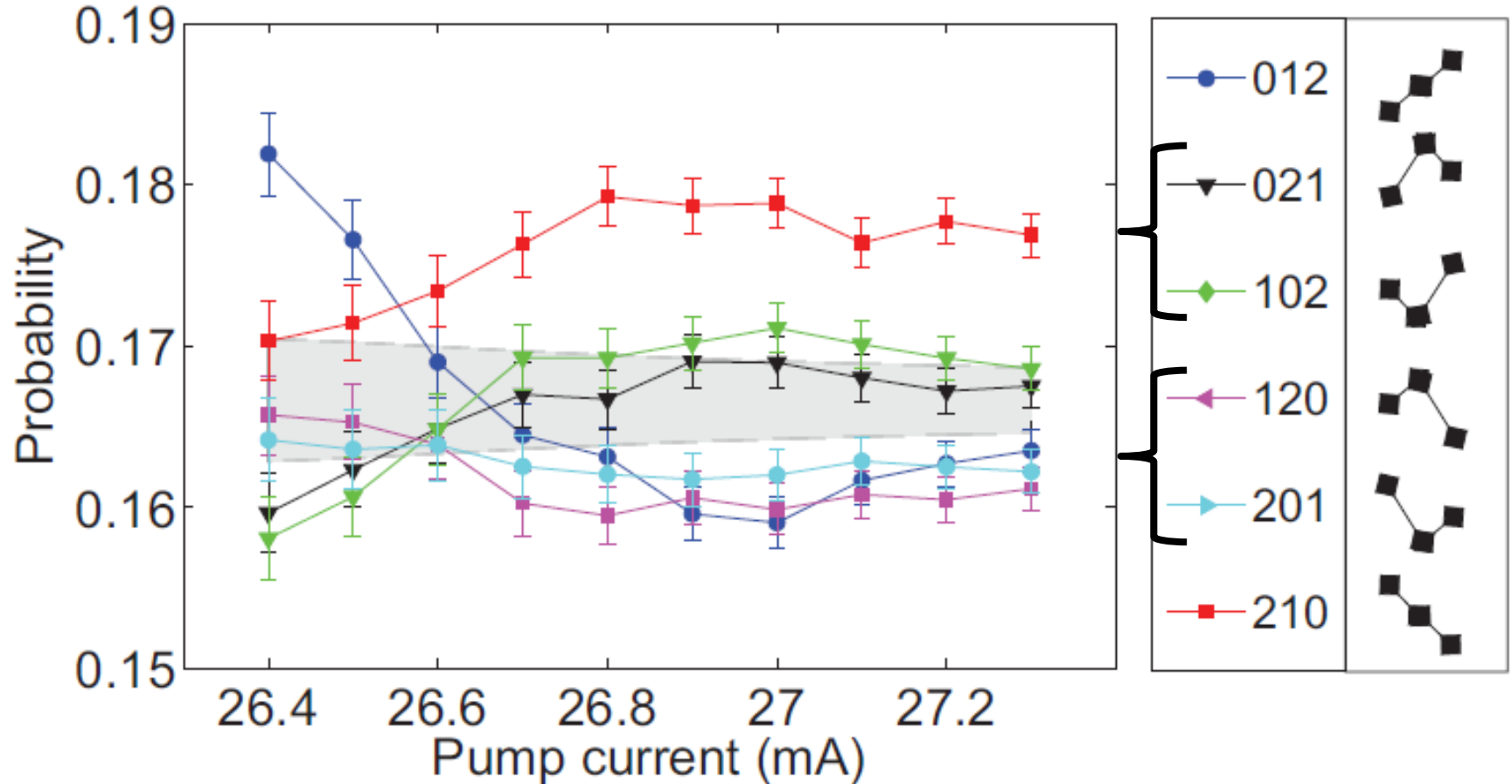


T=20 C



Same transition, hierarchical and clustered organization of the pattern probabilities

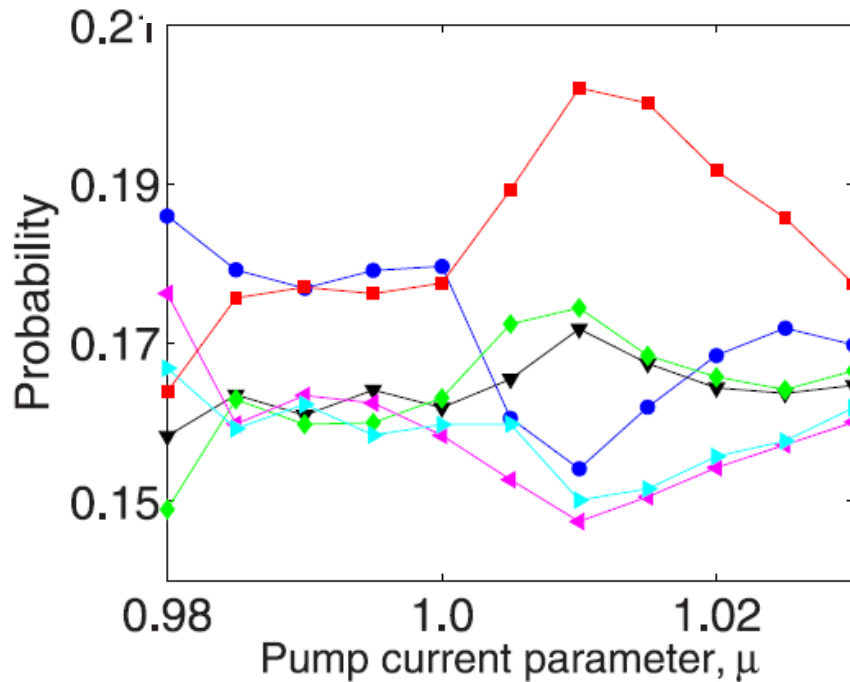
In another experiment: the same transition, hierarchy and clusters



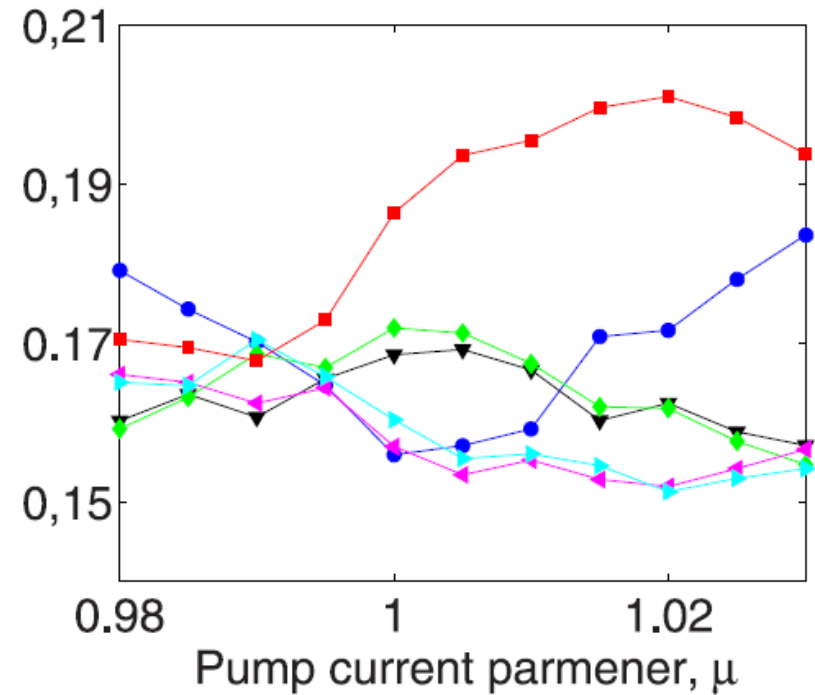
75,000 – 880,000 spikes
(different laser, new oscilloscope)

LK model in good agreement with observations

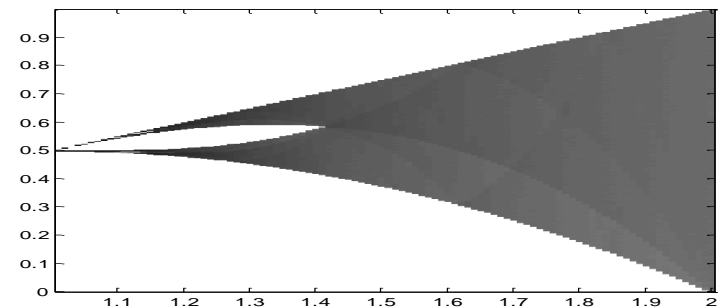
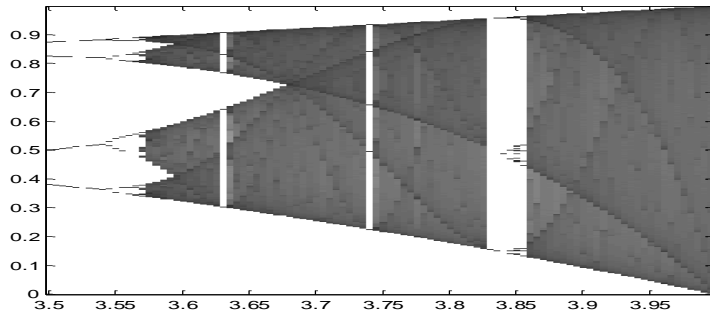
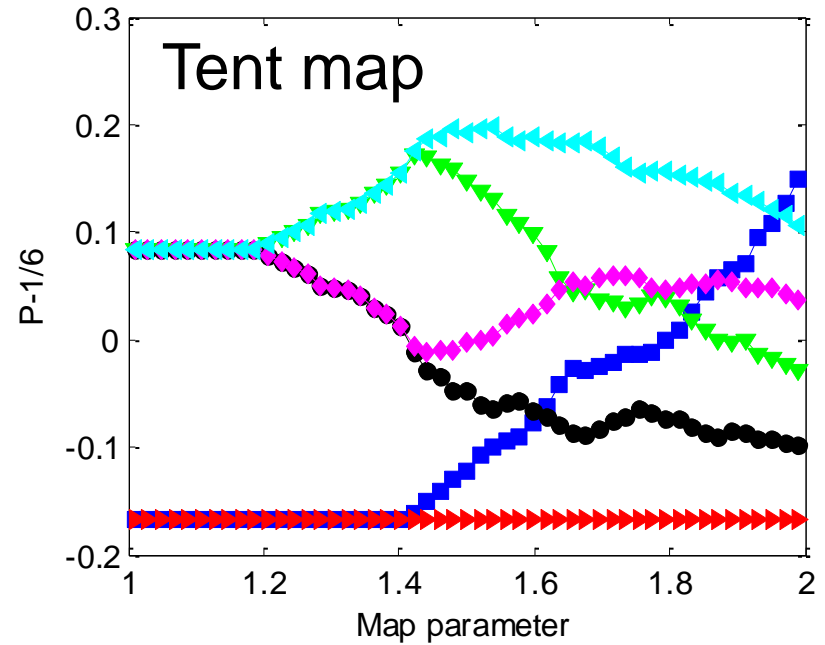
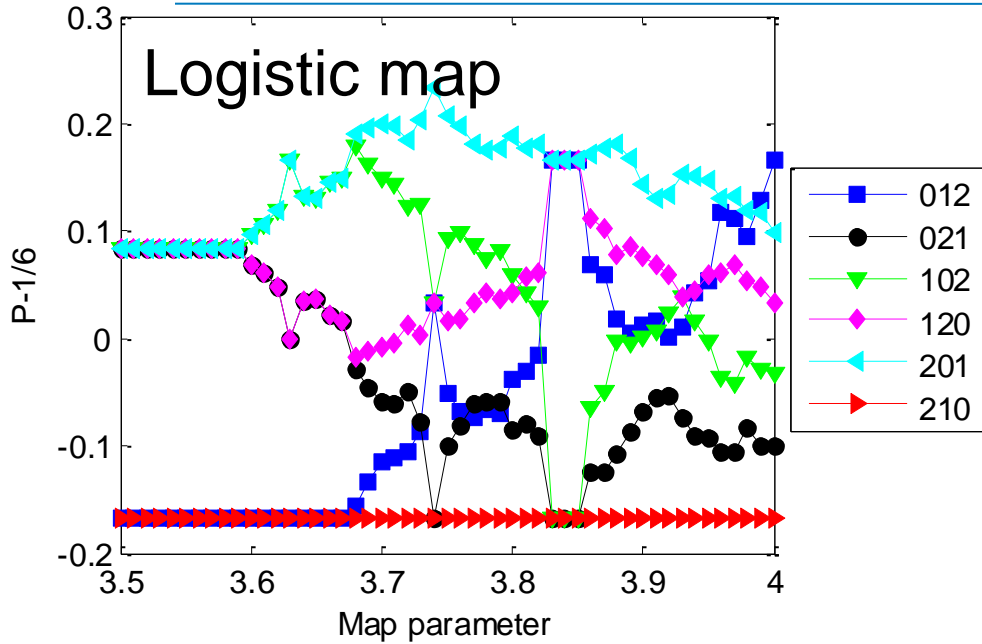
Low feedback



Stronger feedback



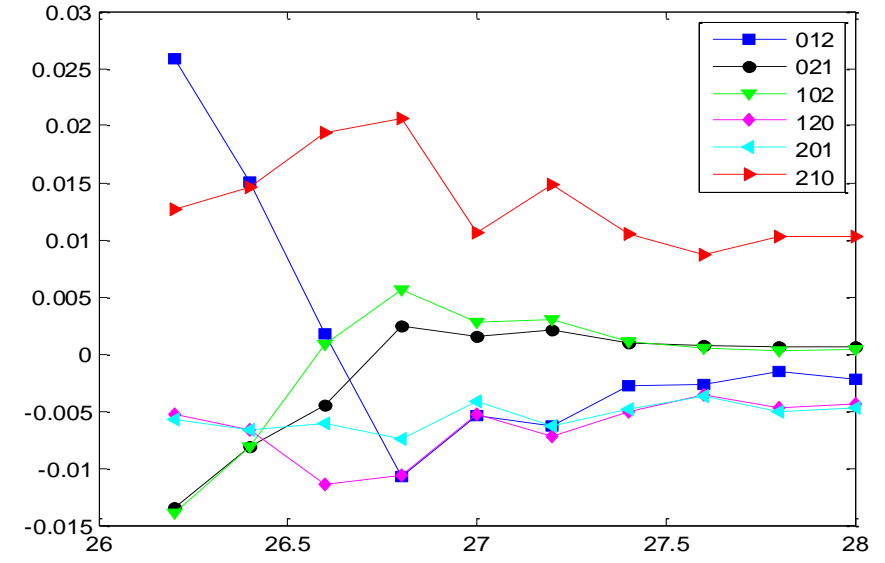
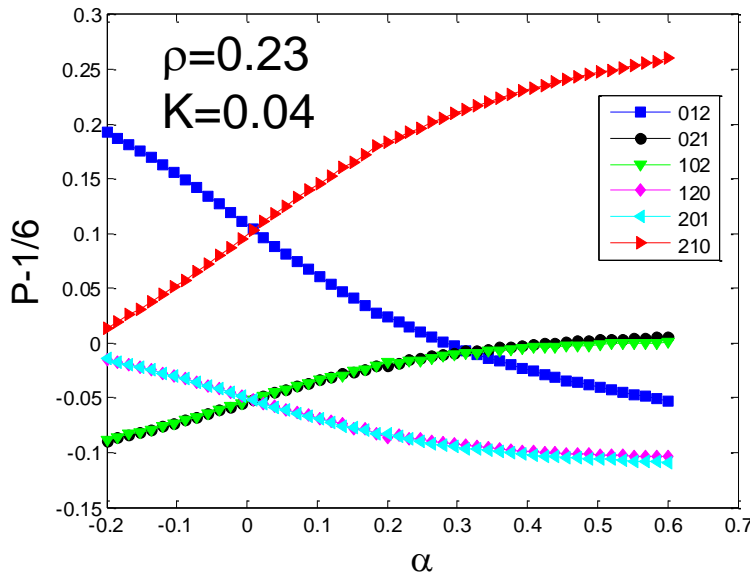
Can we find a minimal model that displays these features?



Modified circle map

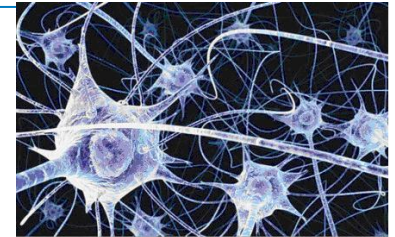
$$\varphi_{i+1} = \varphi_i + \rho + \frac{K}{2\pi} [\sin(2\pi\varphi_i) + \alpha \sin(4\pi\varphi_i)]$$

$$X_i = \varphi_{i+1} - \varphi_i$$



Minimal phenomenological model

- The circle map describes many excitable systems (neurons)
- The modified circle map has been used to describe spike correlations in biological neurons.



Neiman and Russell, *Models of stochastic biperiodic oscillations and extended serial correlations in electroreceptors of paddlefish*, PRE 71, 061915 (2005)

How similar the spikes of lasers and neurons are?

Neuron Interspike Interval (ISI) histogram

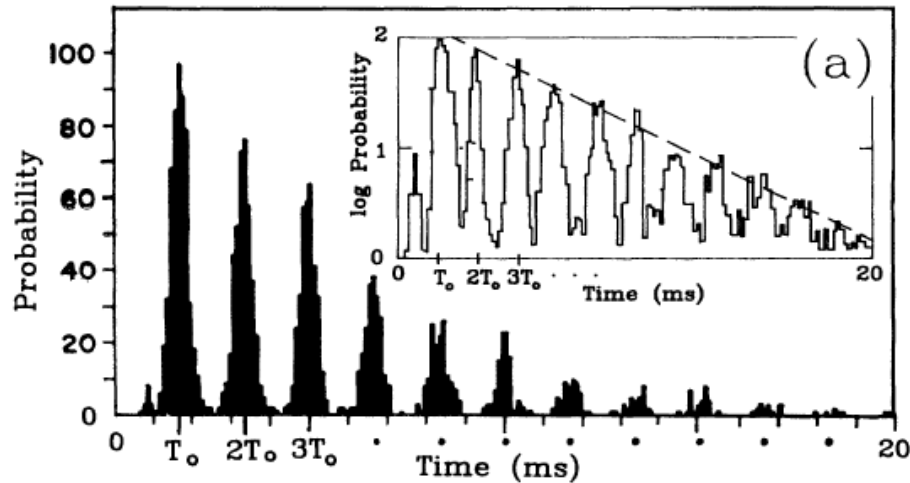
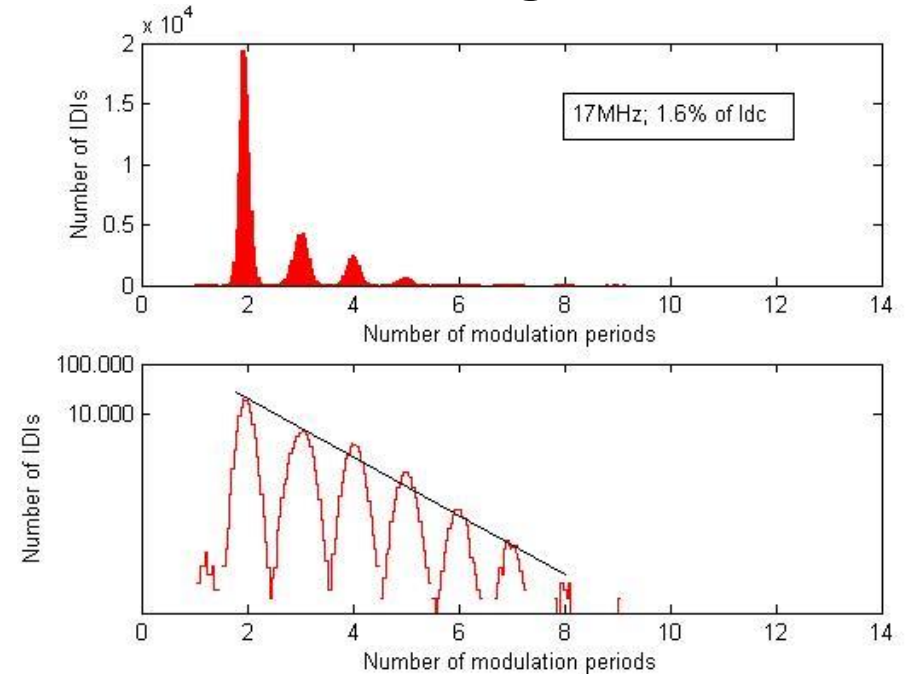


FIG. 1. (a) An experimental ISIH obtained from a single auditory nerve fiber of a squirrel monkey with a sinusoidal 80-dB sound-pressure-level stimulus of period $T_0 = 1.66$ ms applied at the ear. Note the modes at integer multiples of T_0 . Inset:

Laser ISI histogram



With direct current modulation, data recorded in our lab

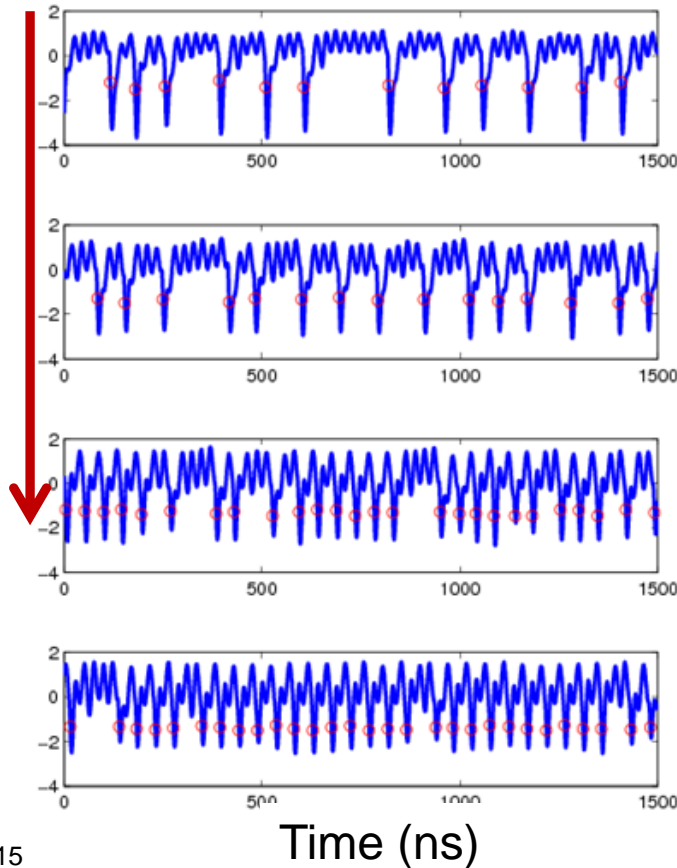
Response to periodic modulation

Relevant for understanding neuronal encoding of external stimuli

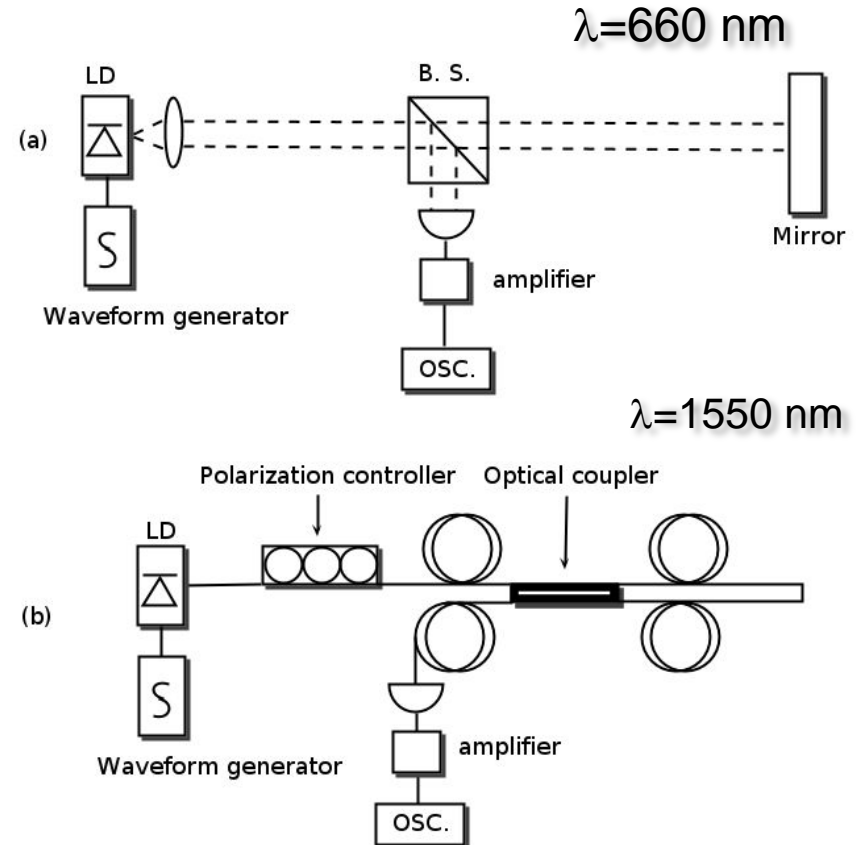
Laser intensity:

Increasing modulation amplitude

Tendency to lock

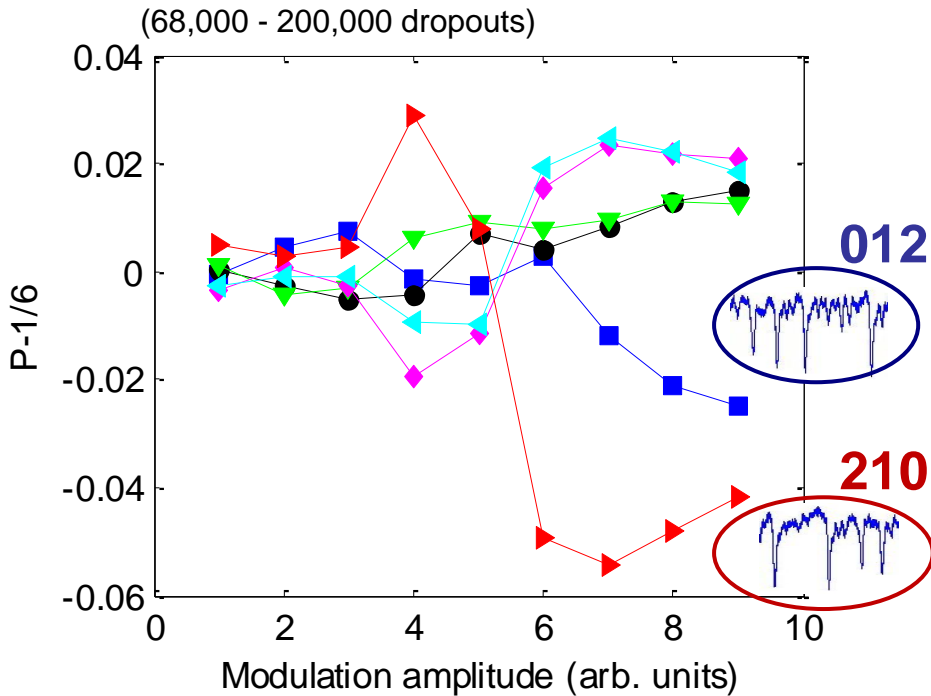


Two experiments:

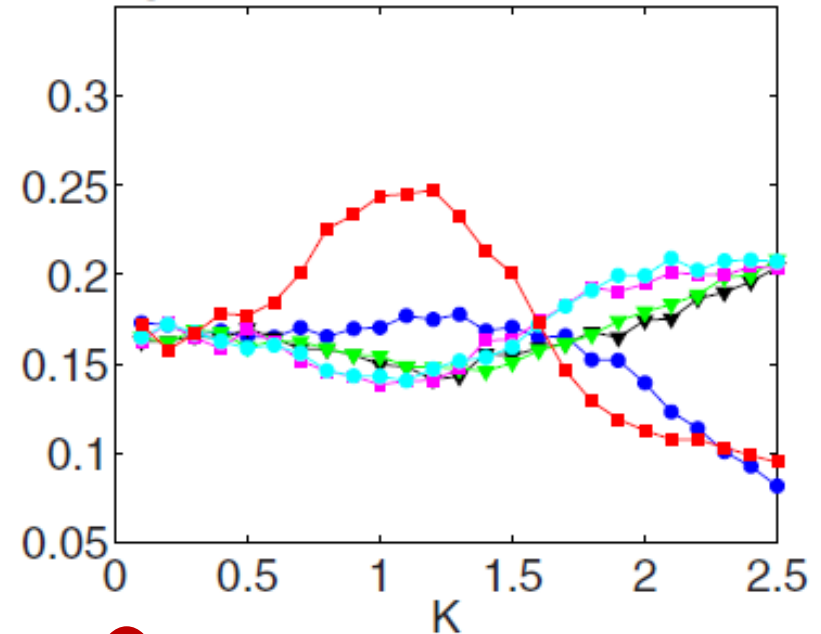


Experiments - minimal model comparison

Experiments @ 660 nm



Circle map



$$\varphi_{i+1} = \varphi_i + \rho + \frac{K}{2\pi} [\sin(2\pi\varphi_i) + \alpha \sin(4\pi\varphi_i)] + D\zeta$$

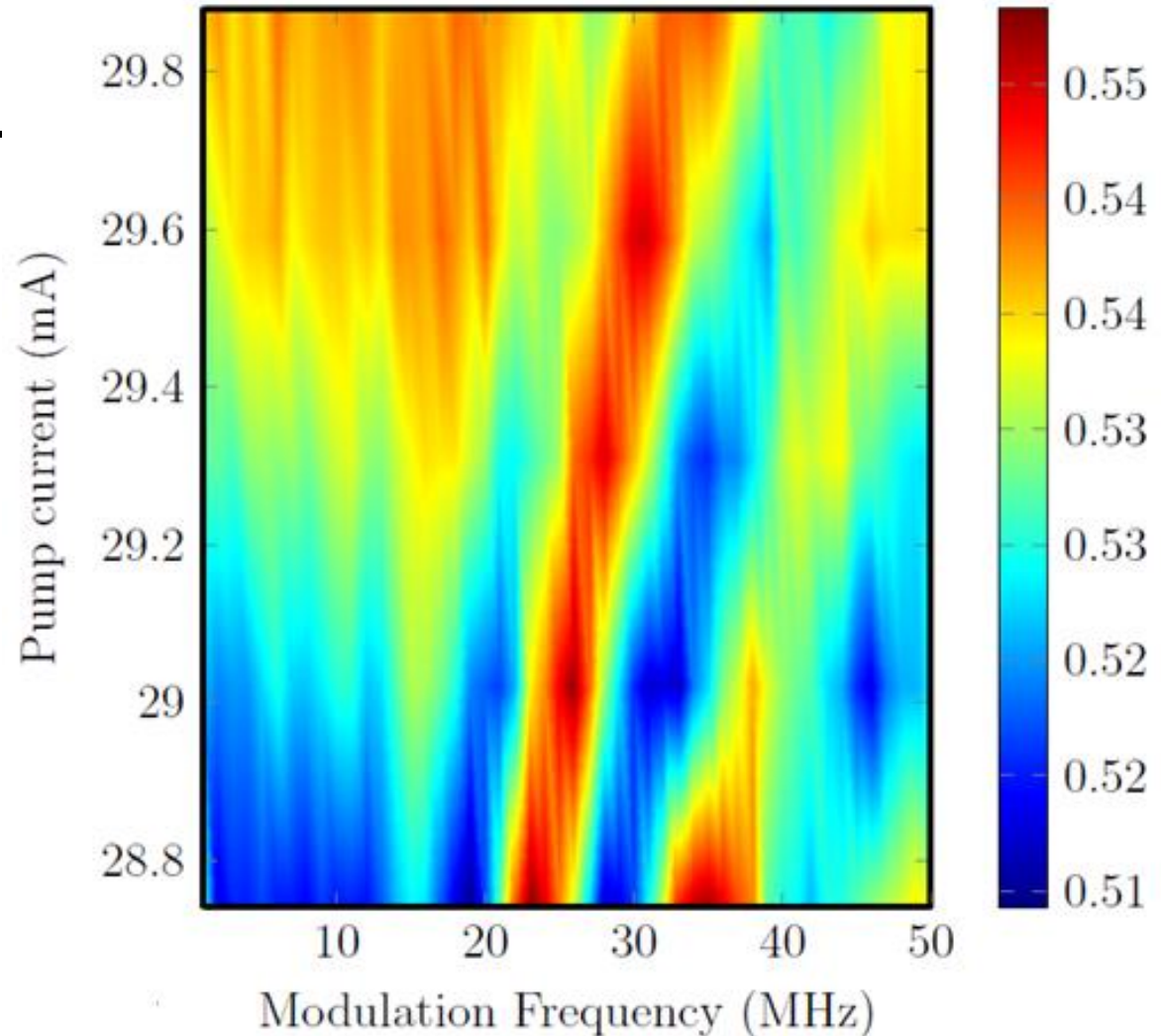
Similar observations @ 1550 nm

Interpretation: locking to external forcing

A valuable tool for identifying noisy locking

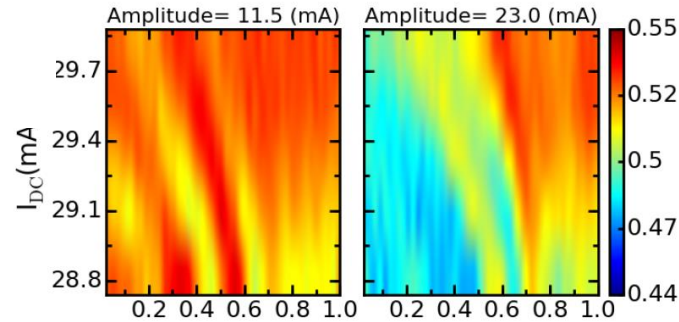
Probability of “10”
from empirical data.

Pump current:
modifies the
natural
(unmodulated)
spike rate

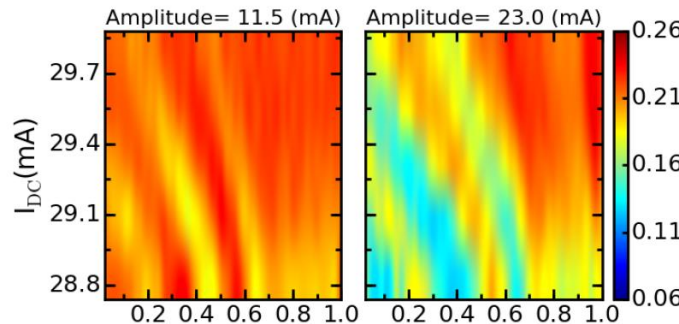


Weak modulation Stronger modulation

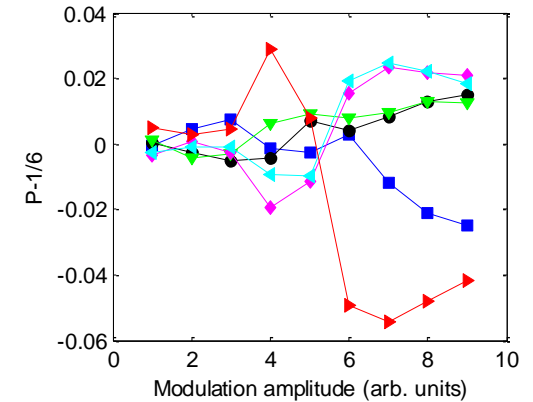
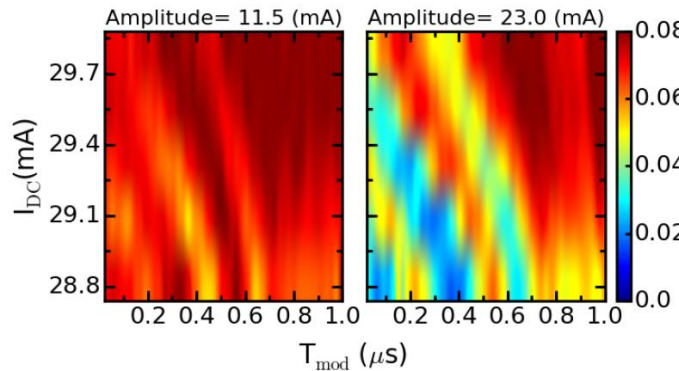
Probability of "10"



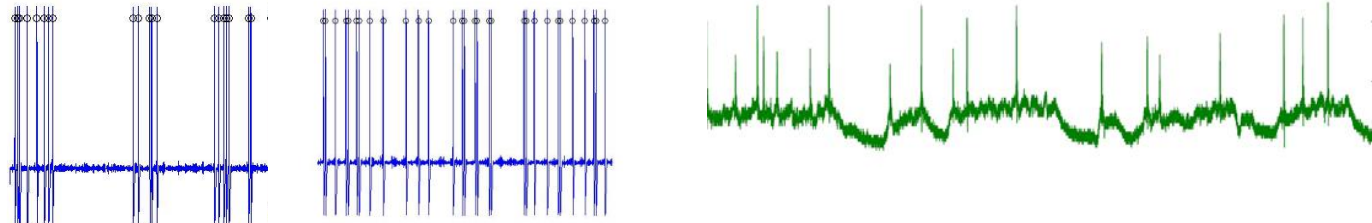
Probability of "210"



Probability of "3210"



- Underlying transition (012→210) in the spiking laser output.
- Hierarchical and clustered organization of pattern probabilities.
- Good agreement with LK model.
- Minimal model identified. Robust under external forcing.
- Present work: can the laser be an “optical neuron”? Detailed comparison with neuronal models & real data.



What is a Rogue Wave?

A “monster wave”, a “freak wave”, an ultra-high wave.

(a) Hokusai's Great Wave



(b) Breaking Wave in the Southern Ocean



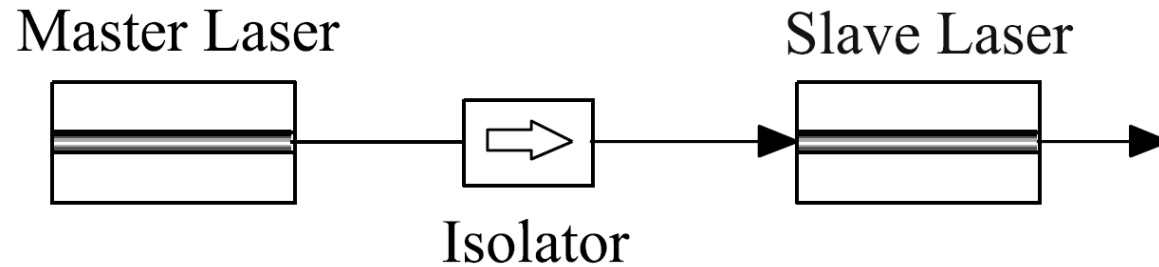
Can develop suddenly even in calm and apparently safe seas.

RWs appear suddenly and vanish without a trace



A challenge for boats and also, for the oil and gas industry, for the design of safe off-shore platforms.

Optically injected semiconductor lasers provide a controllable setup for the study of RWs



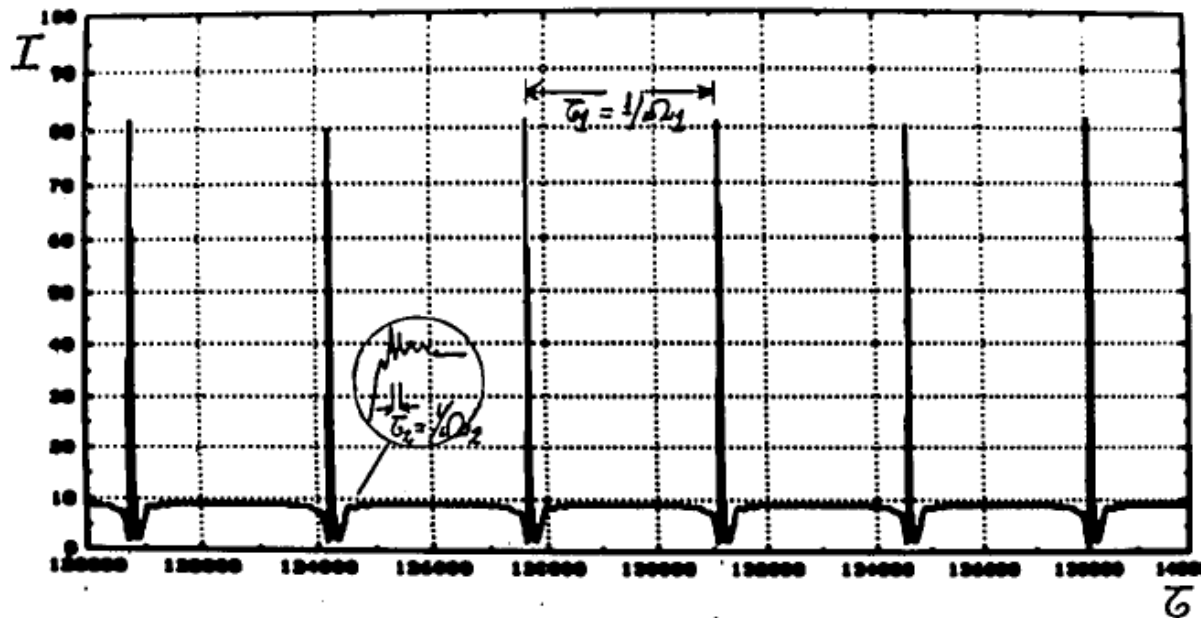
- Parameters:
 - Injection ratio
 - Frequency detuning
- Regimes:
 - Stable emission
 - Periodic oscillations
 - Chaos

S. Wieczorek, B. Krauskopf, T. B. Simpson, and D. Lenstra, "The dynamical complexity of optically injected semiconductor lasers," Phys. Rep. 216 (2005).

Instabilities in lasers with an injected signal

J. R. Tredicce, F. T. Arecchi, G. L. Lippi, and G. P. Puccioni

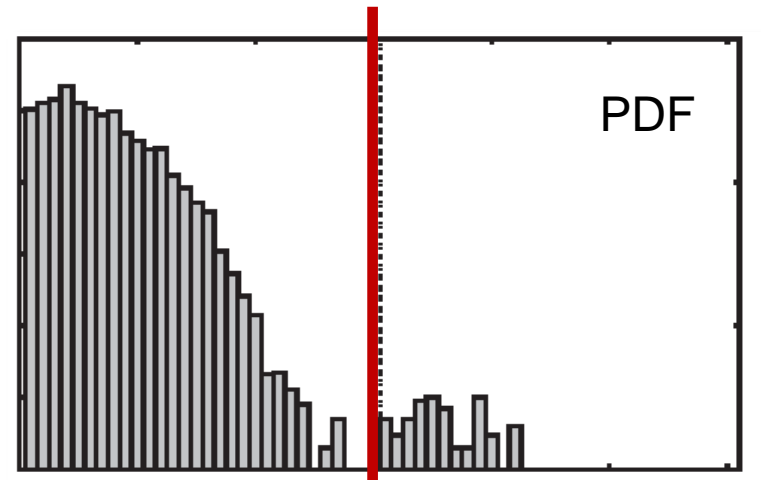
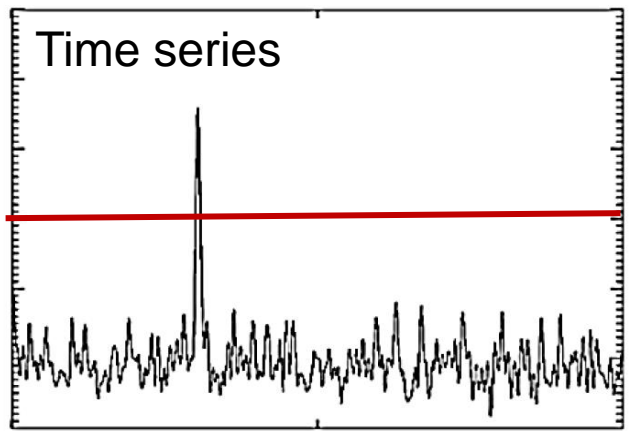
178 J. Opt. Soc. Am. B/Vol. 2, No. 1/January 1985



Experimental observation of extreme pulses in the laser output



- RW definition: pulse above a threshold ($\langle H \rangle + 4-8 \sigma$)



- RWs can be **deterministic**, generated by a crisis-like process.
- RWs can be **predicted** with a certain anticipation time.
- RWs can be **controlled** via noise and/or modulation.
 - C. Bonatto et al, *Deterministic optical rogue waves*, PRL 107, 053901 (2011).
 - J. Zamora-Munt et al, *Rogue waves in optically injected lasers: origin, predictability and suppression*, PRA 87, 035802 (2013).
 - S. Perrone et al, *Controlling the likelihood of RWs in an optically injected semiconductor laser via direct current modulation*, PRA 89, 033804 (2014).
 - J. Ahuja et al, *Rogue waves in injected semiconductor lasers with current modulation: role of the modulation phase*, Optics Express 22, 28377 (2014).

Governing equations

- Complex field, **E**
- Carrier density, **N**

$$\frac{dE}{dt} = \frac{1}{2\tau_p} (1 + i\alpha)(N - 1)E + \underbrace{i\Delta\omega + \sqrt{P_{inj}}}_{\text{optical injection}} + \underbrace{\sqrt{2\beta_{sp} / \tau_N} \xi(t)}_{\text{spontaneous emission noise}}$$

$$\frac{dN}{dt} = \frac{1}{\tau_N} (\mu - N - N|E|^2)$$

Solitary laser
parameters: α τ_p τ_N μ

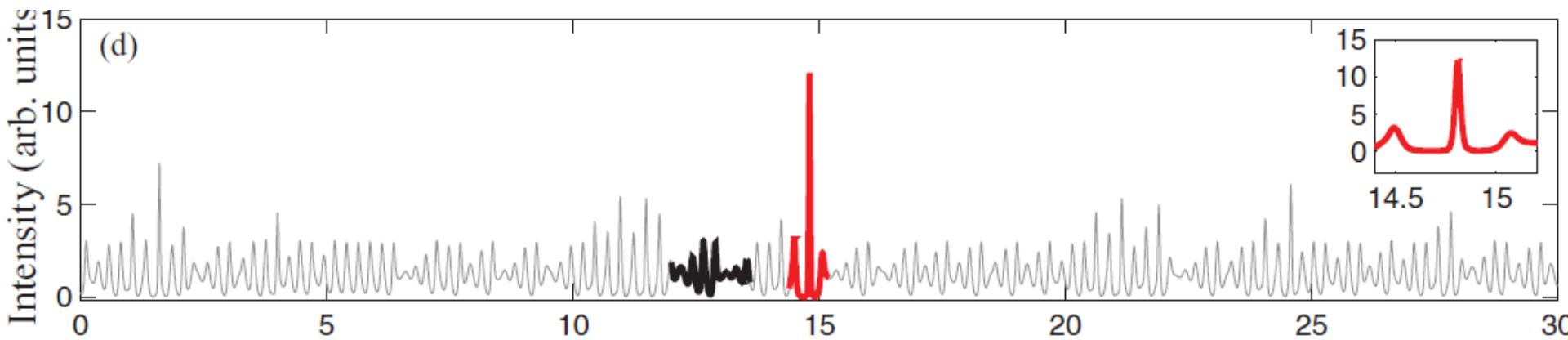
optical injection
 η : injection strength
 $\Delta\omega = \omega_s - \omega_m$: detuning

spontaneous
emission
noise

Typical parameter values:

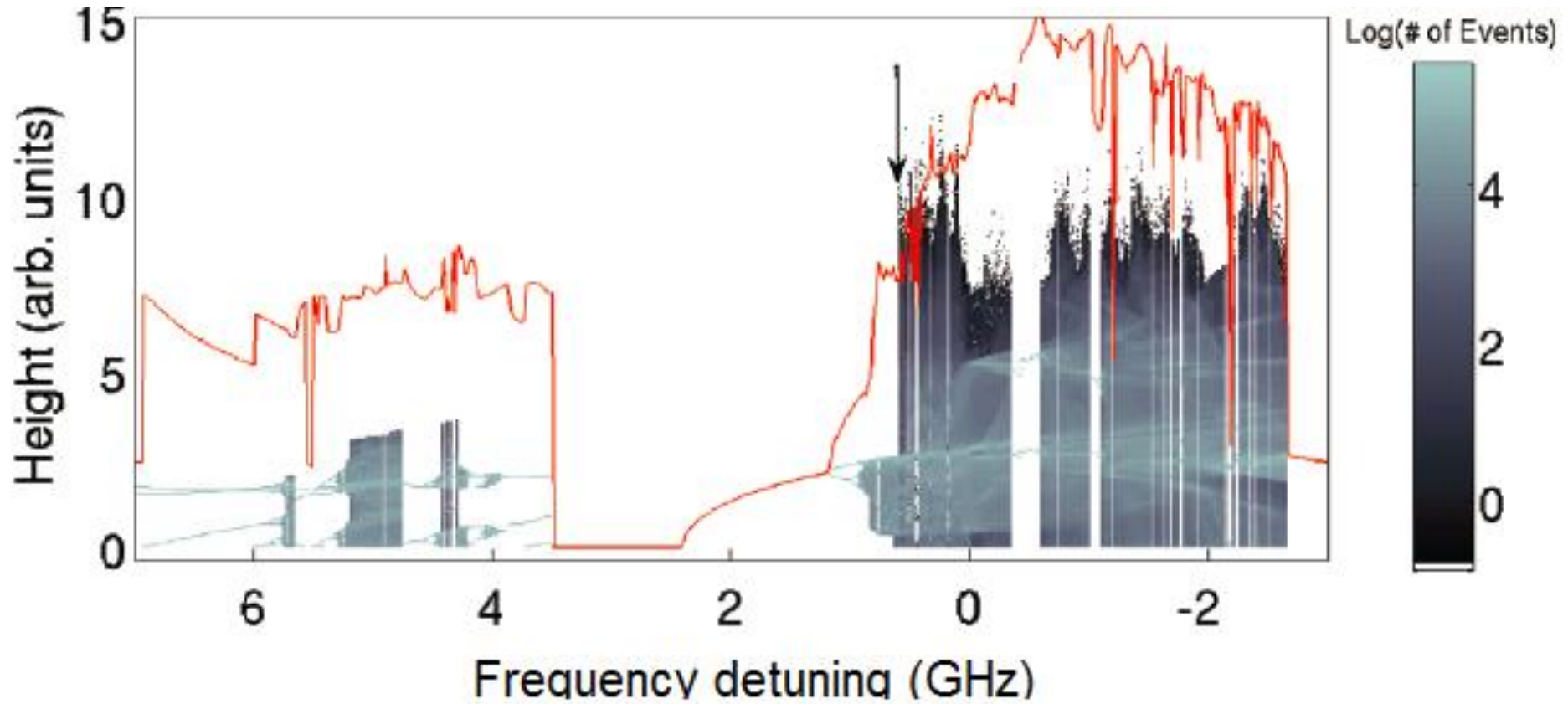
$\alpha = 3$, $\tau_p = 1$ ps, $\tau_N = 1$ ns

μ : normalized pump current parameter



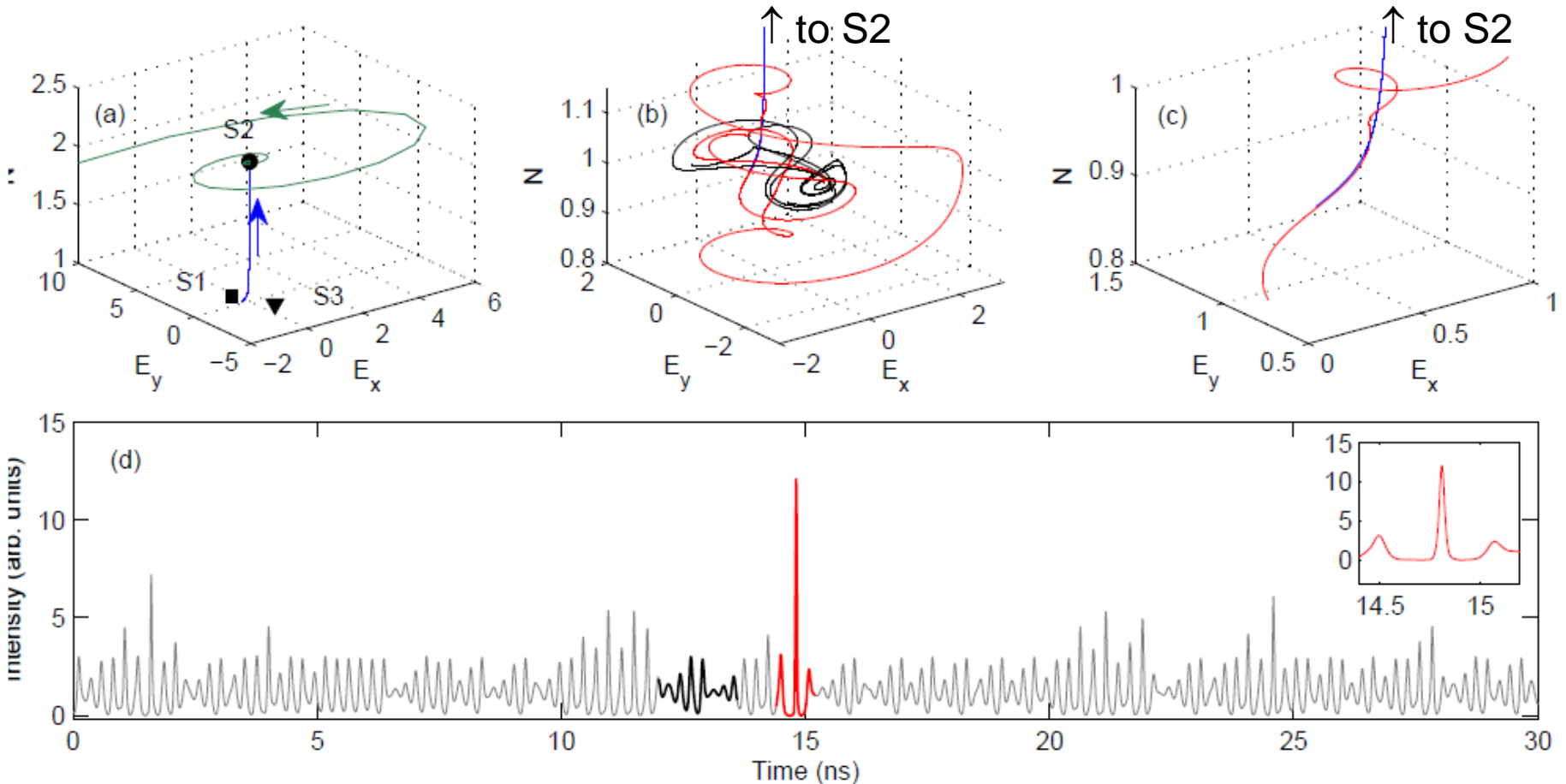
Bifurcation diagram

Threshold: $\langle H \rangle + 8\sigma$



$$\Delta\omega = \omega_s - \omega_m$$

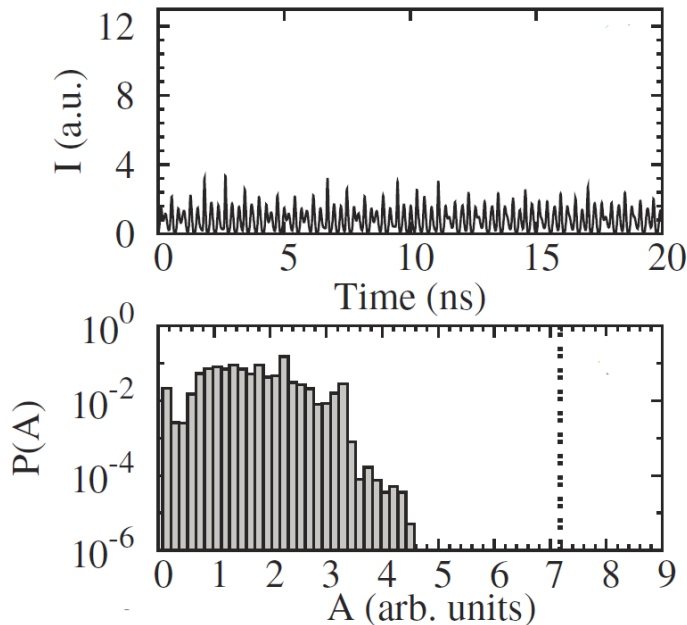
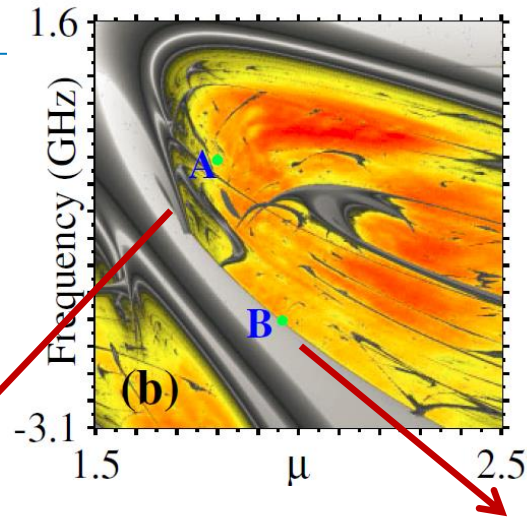
A narrow channel: the RW "door"



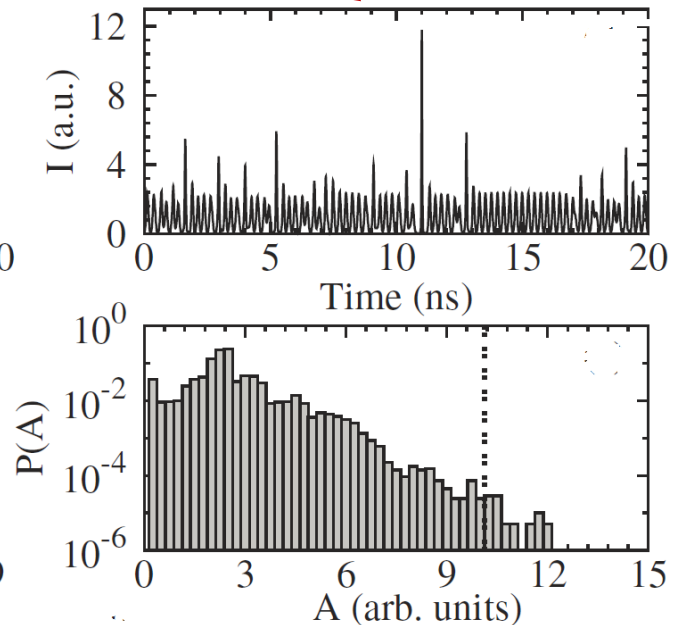
A **RW** is triggered whenever the trajectory closely approaches **the stable manifold of S2**

Deterministic simulations

Lyapunov diagram
(detuning, pump current)

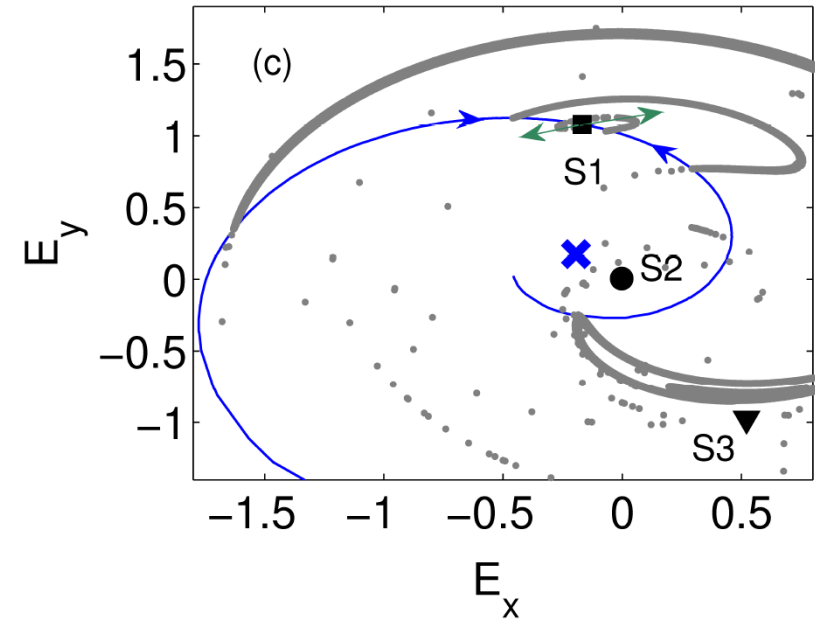
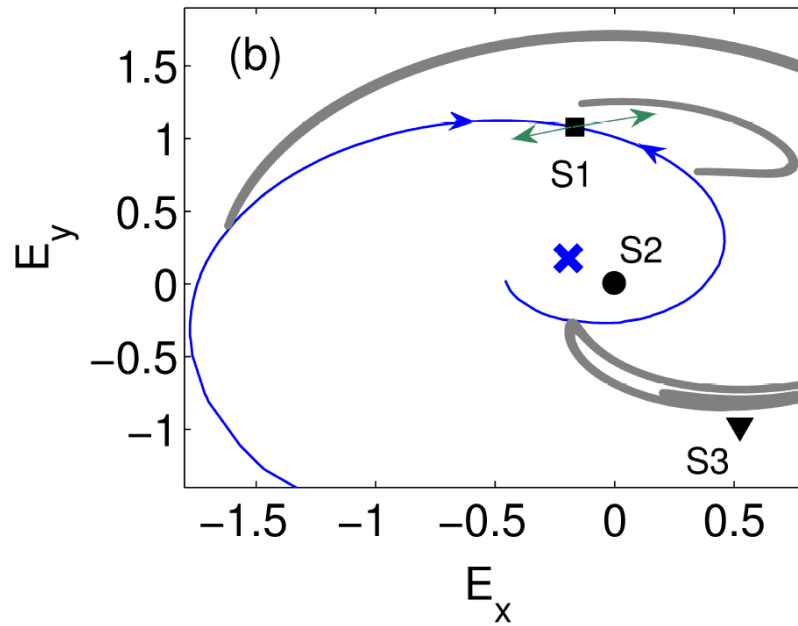


Chaos without RWs



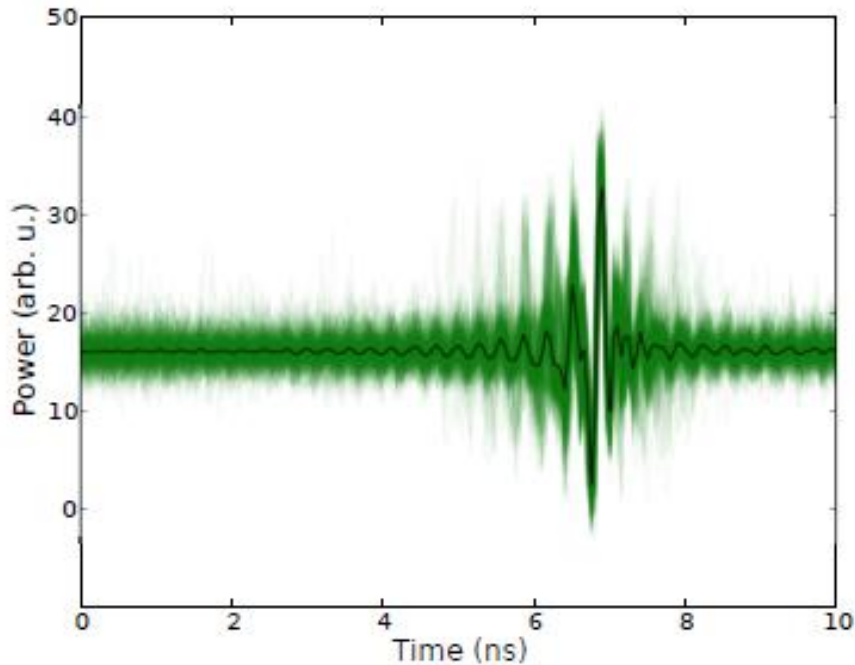
Chaos with RWs

Why chaos with RWs and chaos without them?

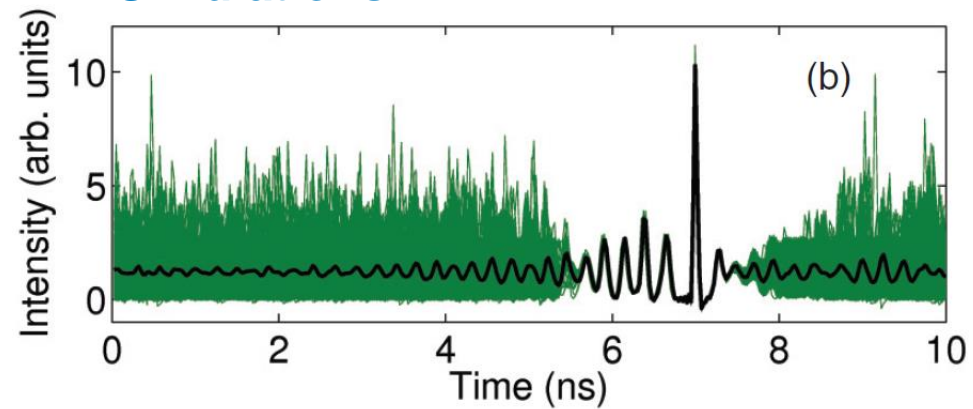


An **external crises-like** process enables access to the region of phase space where the **stable manifold of S2 (x)** is.

Experiments



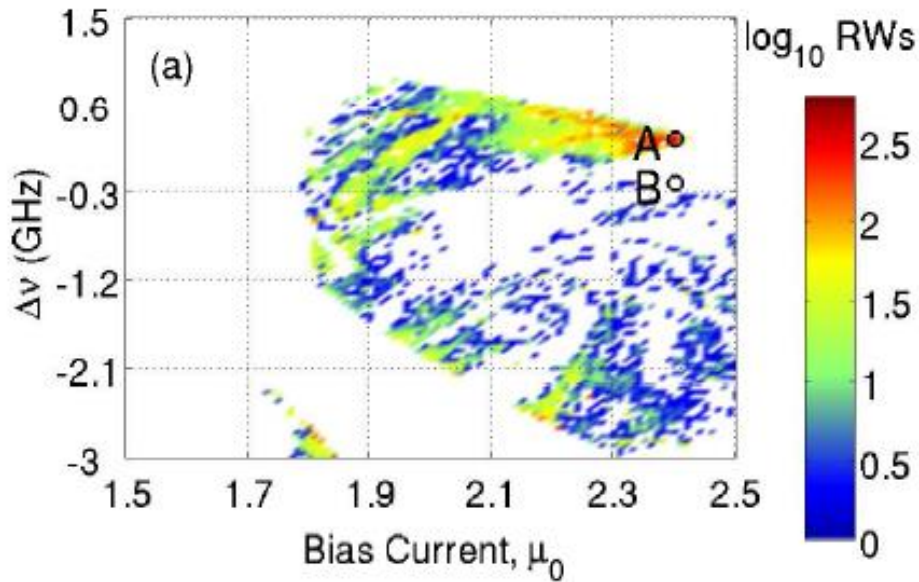
Simulations



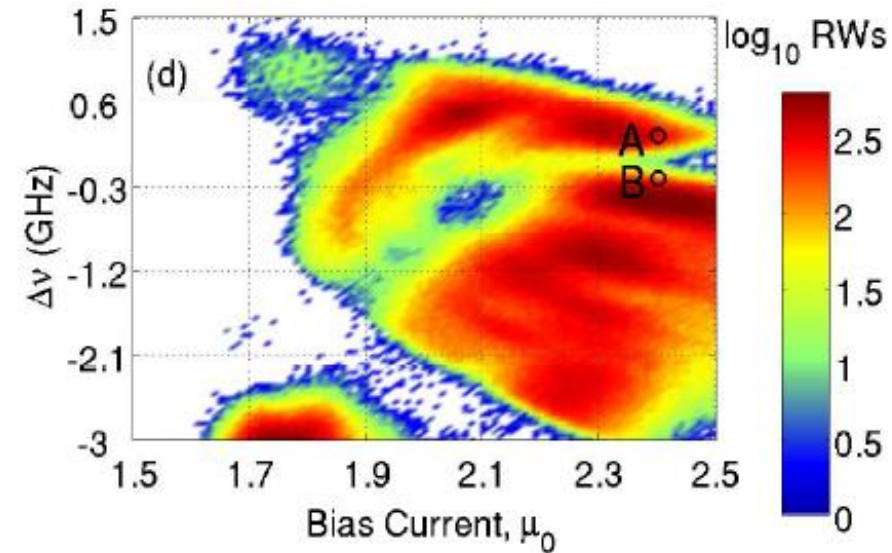
Superposition of 500 time series at the RW peak

Role of noise: number of RWs vs (pump current, detuning)

Deterministic RWs ($\beta_{sp}=0$)

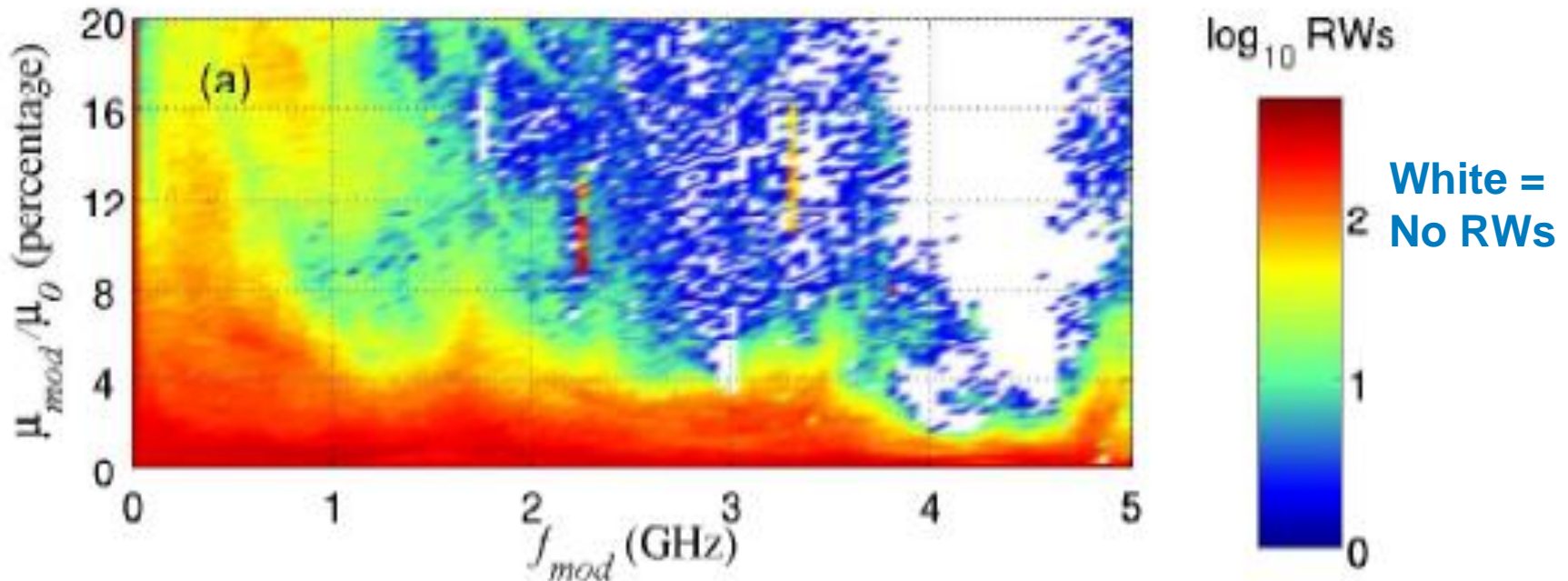


Stochastic simulations ($\beta_{sp}=0.01$)



White = No RWs

$$\mu = \mu_0 + \mu_{\text{mod}} \sin(2\pi f_{\text{mod}} t)$$

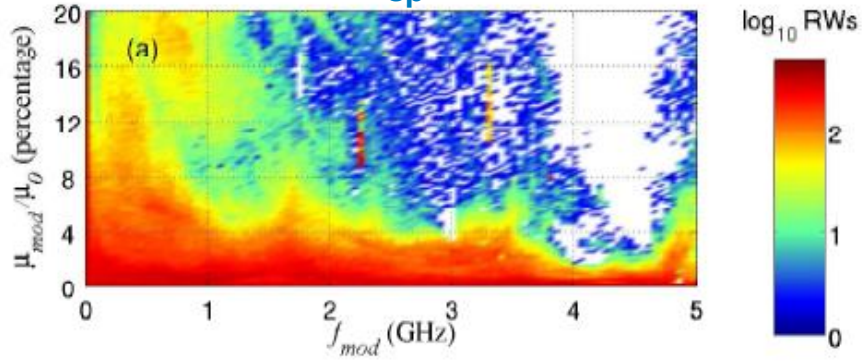


Current modulation with appropriated amplitude and frequency completely suppresses RWs.

RW control in Point A: influence of noise

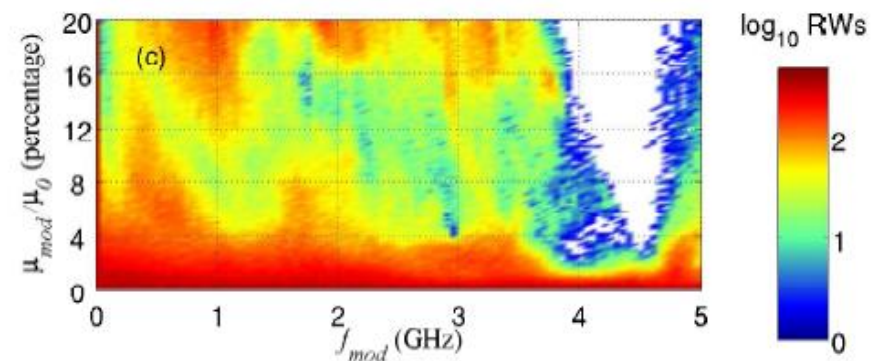
$$\mu = \mu_0 + \mu_{\text{mod}} \sin(2\pi f_{\text{mod}} t)$$

No noise ($\beta_{\text{sp}}=0$)



White = No RWs

Stochastic simulations ($\beta_{\text{sn}}=0.01$)

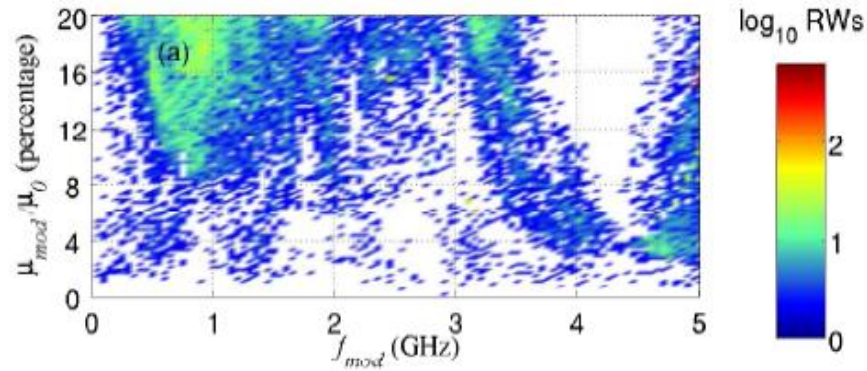


White = No RWs

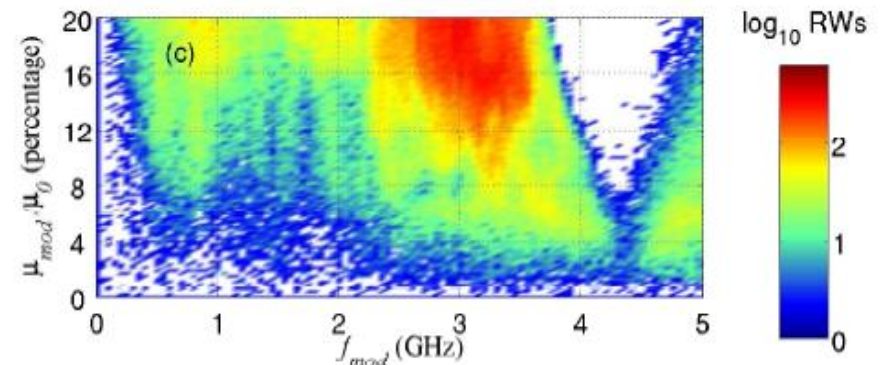
“safe parameter region” is robust to the presence of noise.

in Point B (no deterministic RWs)

$\beta_{sp}=0$



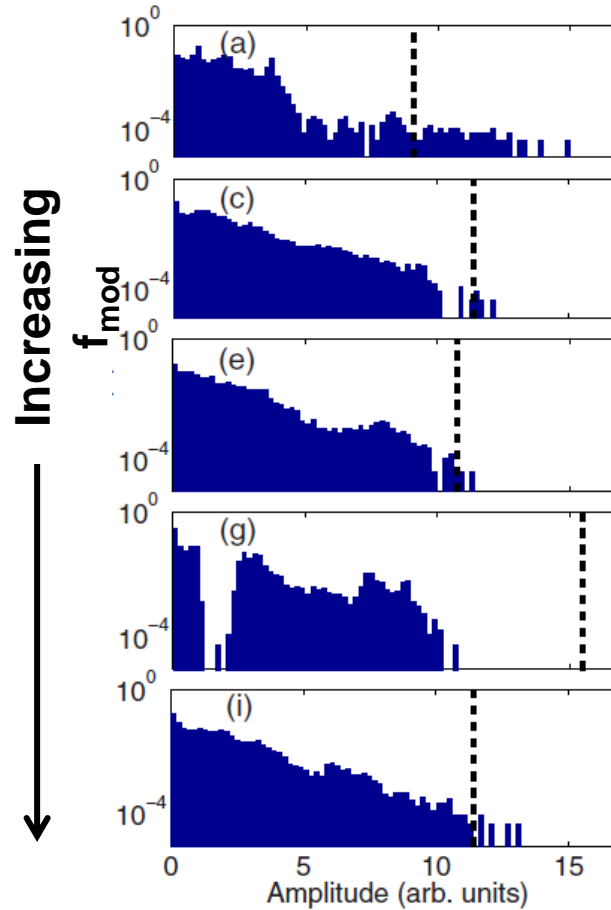
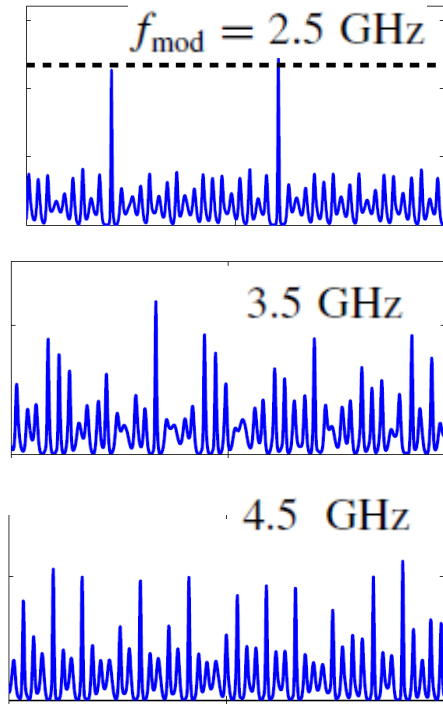
$\beta_{sp}=0.01$



White = No RWs

“safe parameter region” also in point B, also robust to noise

Modulation suppressed RWs

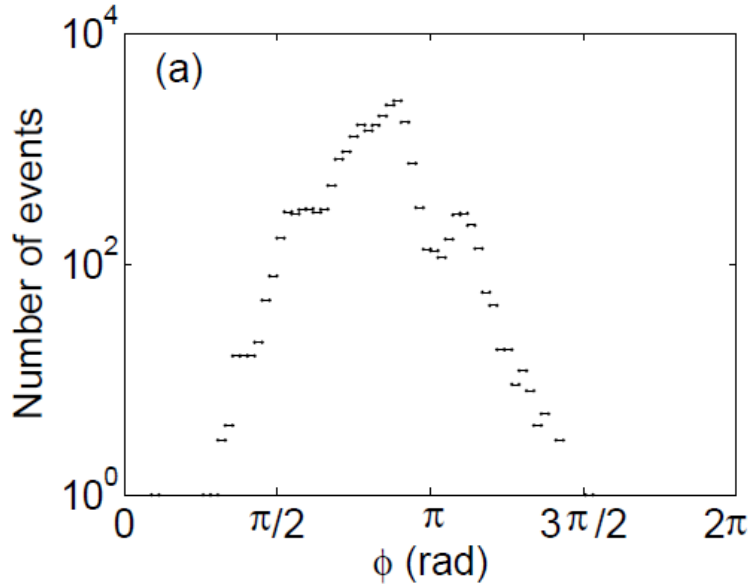


Analogy: avalanche risk
Triggering controlled small avalanches avoids a large and dangerous avalanche.

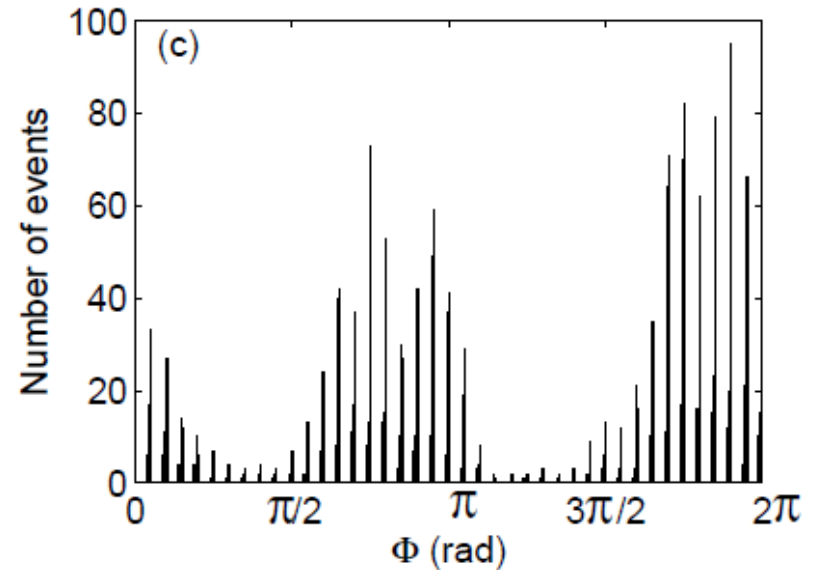


When RWs are not suppressed: role of the modulation phase

Slow modulation



Fast modulation



RWs occur during the first $\frac{3}{4}$ of the modulation cycle.

- **RW Control:** noise and modulation strongly affect the likelihood of RWs.
- **RW Predictability:**
 - RWs can be predicted with some anticipation.
 - with modulation, they occur at certain values of the modulation phase.

Papers @ www.fisica.edu.uy/~cris

- C. Bonatto et al, PRL 107, 053901 (2011).
- J. Zamora-Munt et al, PRA 87, 035802 (2013).
- S. Perrone et al, Phys. Rev. A 89, 033804 (2014).
- J. Ahuja et al, Optics Express 22, 28377 (2014).



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