

Extreme optical pulses: origin, predictability and control

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Collaborators



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- Jose Rios Leite (Universidade Federal de Pernambuco, Recife, Brasil)

What is a rogue wave?

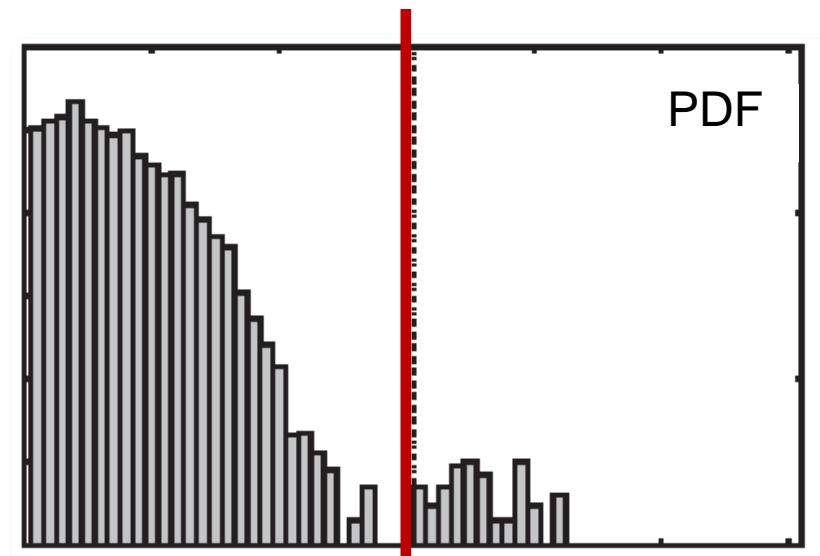
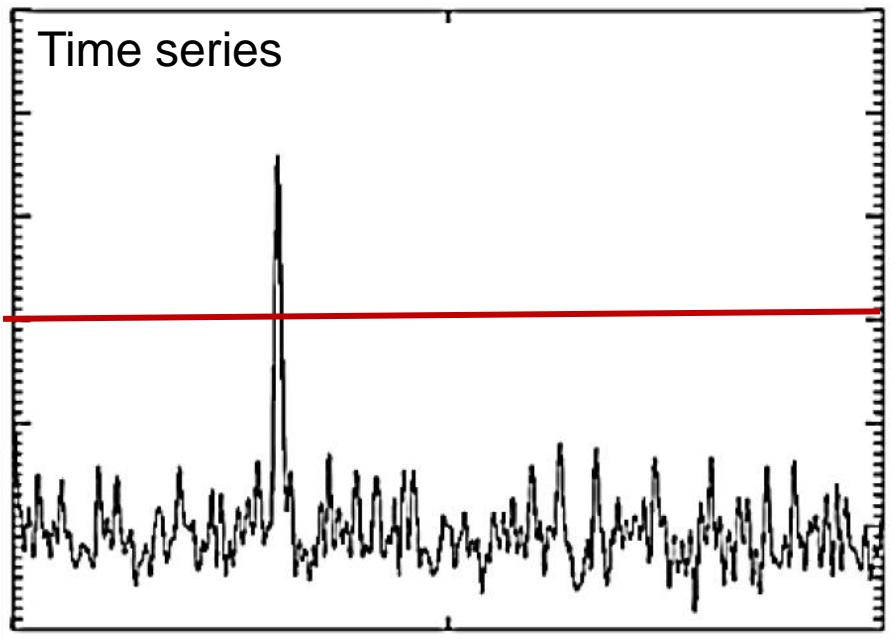
A “monster wave”, ultra-high wave that falls outside (and far from) the main part of a long-tailed probability distribution.



The Great Wave of Kanagawa, Katsushika Hokusai. Source: Wikipedia

Two more precise definitions

- A RW: above **twice** the significant wave height (SWH).
SWH = <highest 1/3 waves> that occur over a certain period of time.
- A RW: above a certain threshold ($\langle H \rangle + 4\text{-}8 \sigma$)



RWs appear suddenly

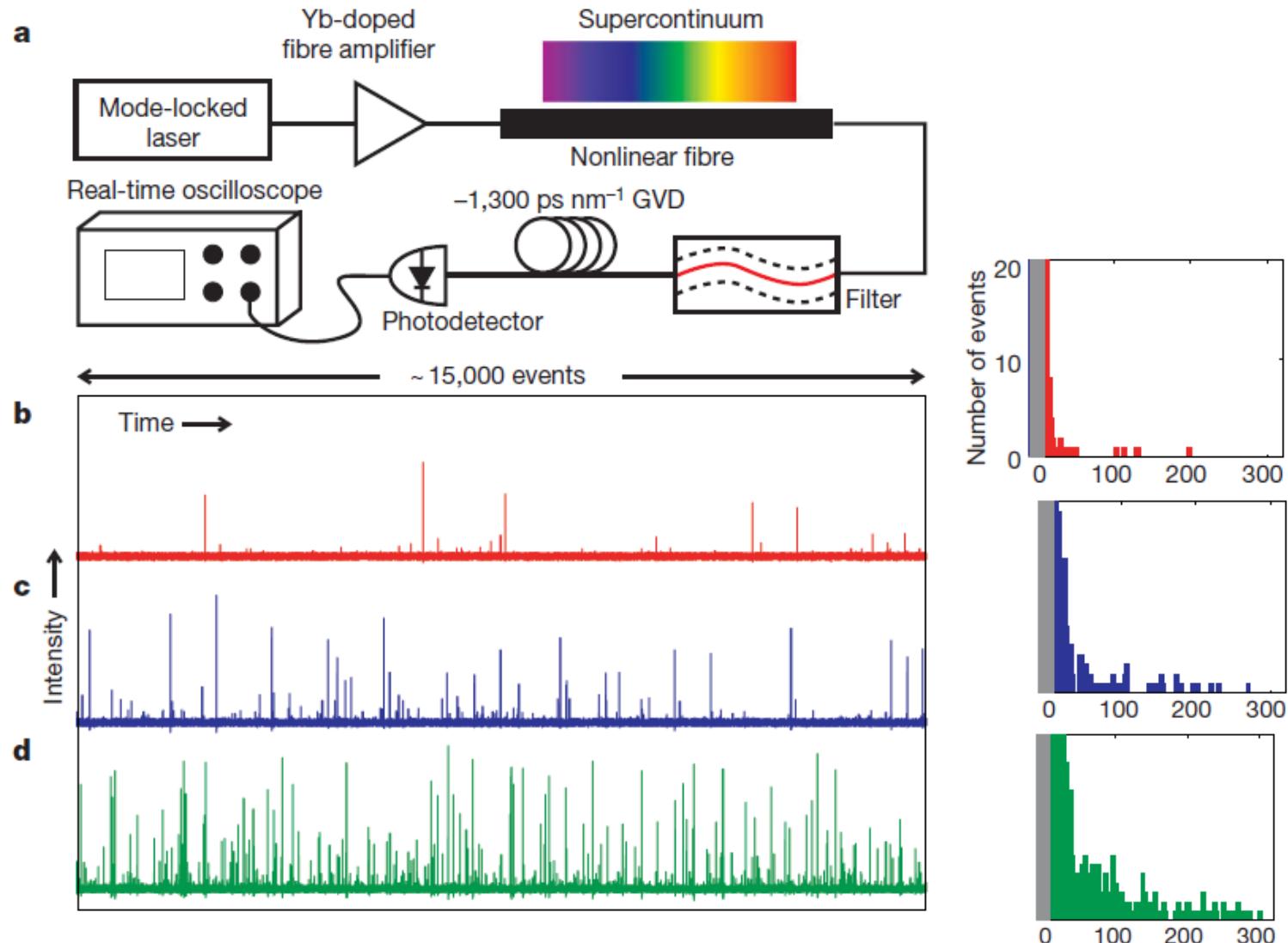


A challenge for the oil and gas industry, for the design of safe off-shore platforms.

Source: National Geographic

Optical RWs: first observation

Campus d'Excel·lència Internacional



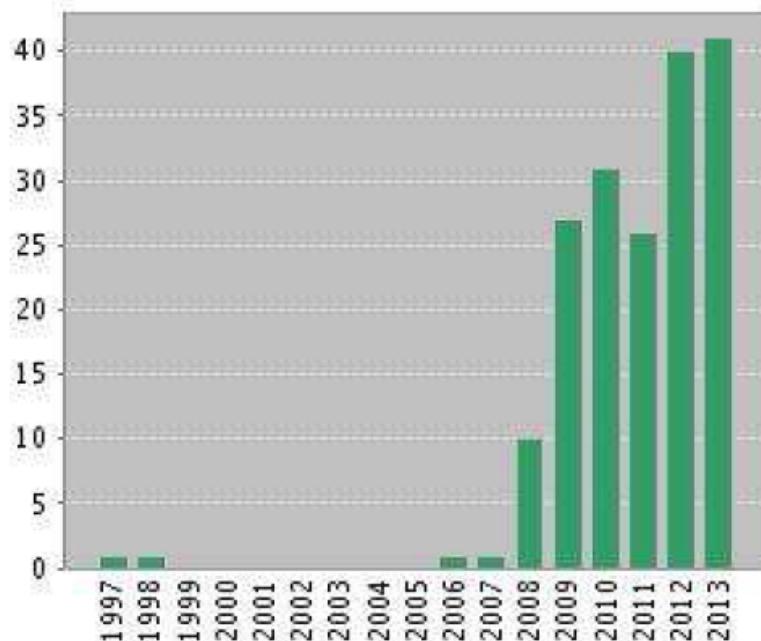
Since 2007: a lot of work

Citation Report Topic=(optical rogue wave)

Timespan=All years. Databases=SCI-EXPANDED, SSCI, A&HCI, CPCI-S, CPCI-SSH.

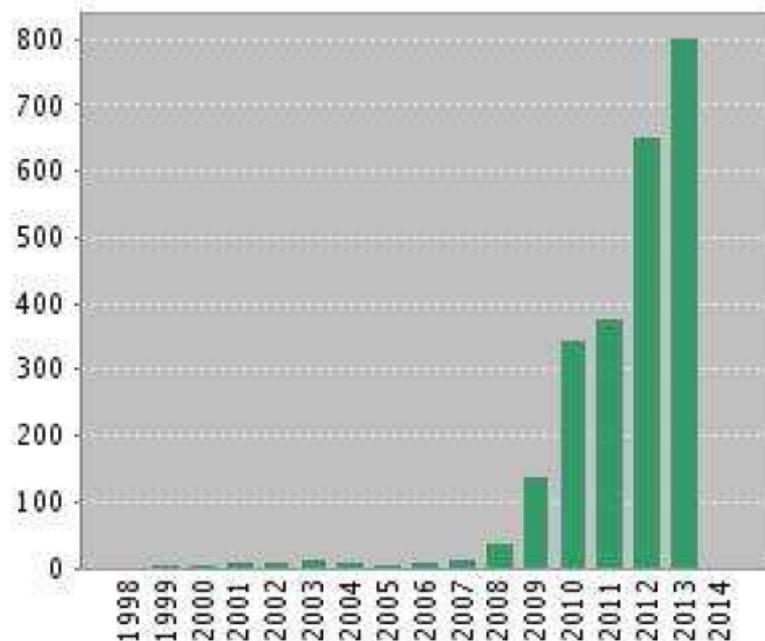
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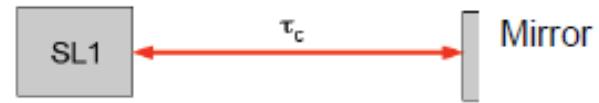
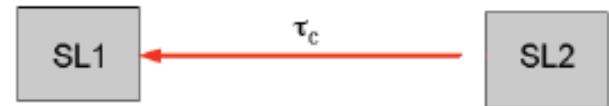


The latest 20 years are displayed.

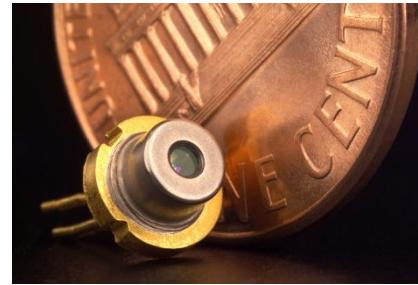
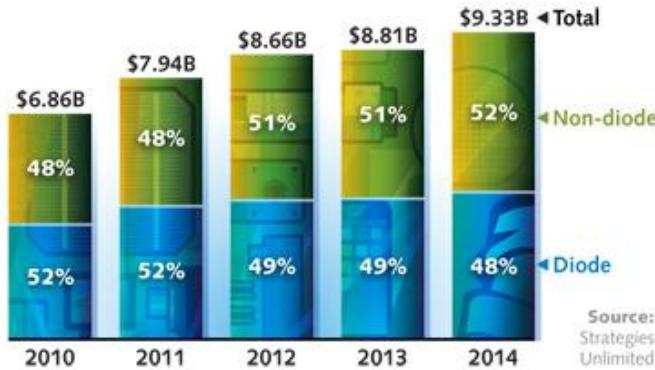
Optical RWs in semiconductor lasers

Two setups:

- With **optical injection**:
observations & numerical results.
- With **optical feedback**:
numerical results.



Why semiconductor lasers? (diode lasers)



- Used in:
 - Telecommunications
 - Optical data storage (CDs, DVDs, Blu rays)
 - Barcode scanners, printers, mouse
 - Biomedical applications (imaging, sensing, etc)

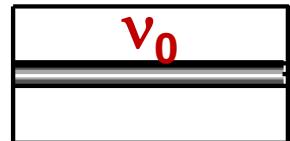
Semiconductor lasers provide an inexpensive and controllable setup for the study of optical rogue waves

What we have learned

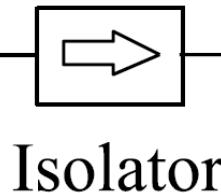
- RWs can be **deterministic**, generated by a crisis-like process.
- RWs can be **predicted** with a certain anticipation time.
- RWs can be **controlled** via noise and/or modulation.
- C. Bonatto et al, *Deterministic optical rogue waves*, PRL 107, 053901 (2011).
- J. Zamora-Munt et al, *Rogue waves in optically injected lasers: origin, predictability and suppression*, PRA 87, 035802 (2013).
- J. M. Reinoso et al, PRE 87, 062913 (2013).
- S. Perrone et al, PRA 89, 033804 (2014).

First setup: an optically injected laser

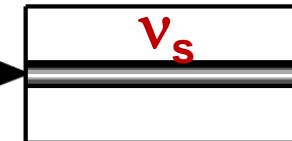
Master Laser



Tunable SCL



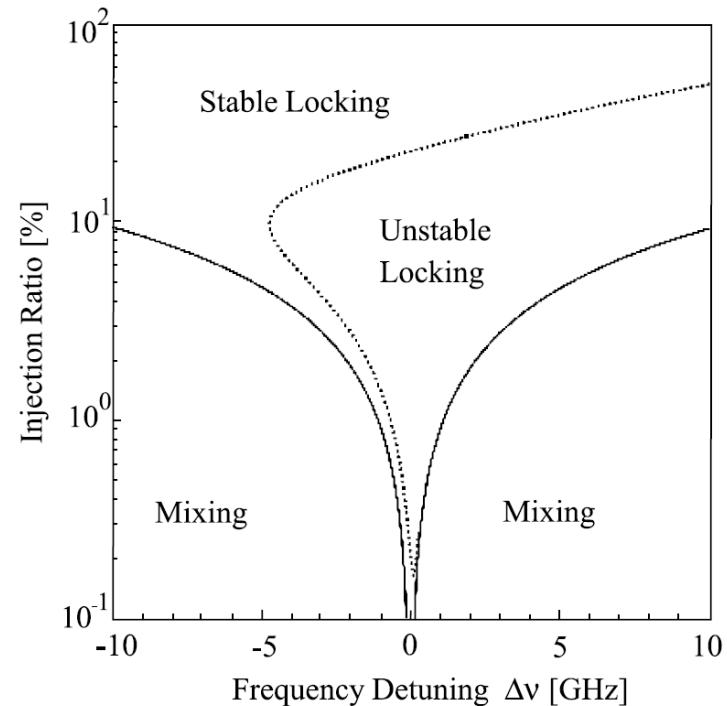
Slave Laser



VCSEL 980 nm

Detection system
(photo detector,
oscilloscope,
spectrum
analyzer)

- Parameters:
 - Injection ratio
 - Frequency detuning $\Delta\nu = \nu_s - \nu_0$
- Dynamical regimes:
 - Injection locking (cw output)
 - Period-one oscillation
 - Period-two oscillation
 - Chaos



Labyrinth bifurcations in optically injected diode lasers

V. Kovanis¹, A. Gavrielides², and J.A.C. Gallas^{3,4,5,a}

¹ Air Force Research Laboratory, 2241 Avionics Circle, Wright-Patterson AFB, Dayton OH 45433, USA

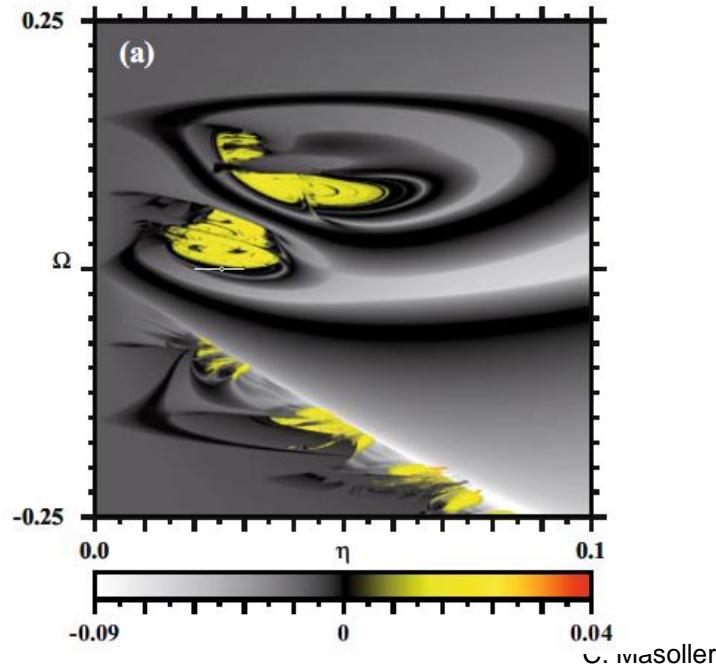
² USAF, Research Laboratory, High Power Solid State Lasers Branch, Kirtland AFB, NM 87117, USA

³ TecEdge, Wright Brothers Institute, 5100 Springfield Street, Dayton OH 45431, USA

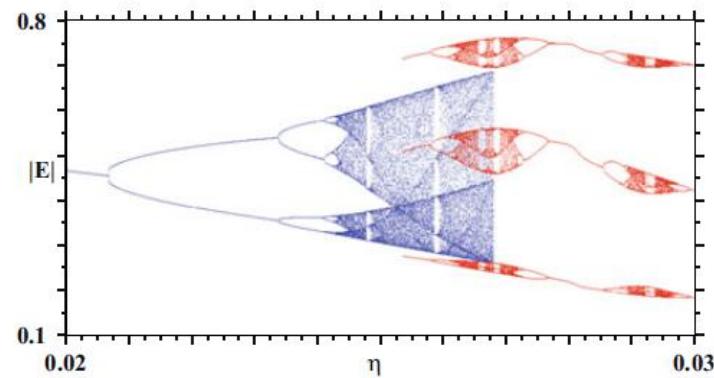
⁴ Departamento de Física, Universidade Federal da Paraíba, 58051-970 João Pessoa, Brazil

⁵ Instituto de Física, Universidade Federal do Rio Grande do Sul, 91501-970 Porto Alegre, Brazil

Lyapunov diagram



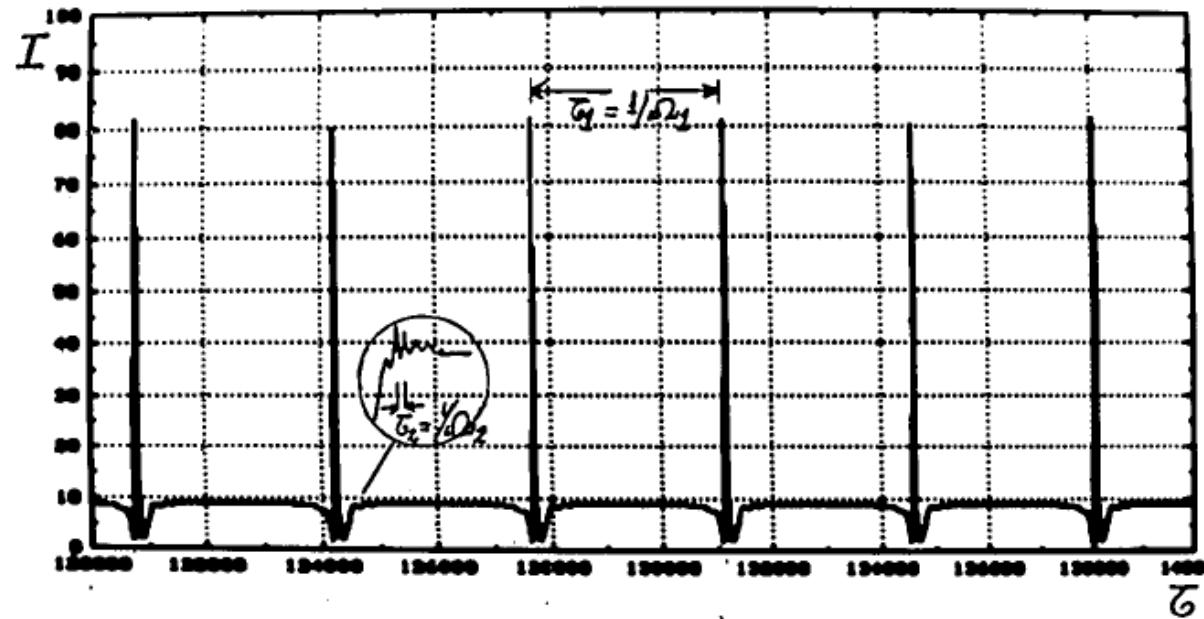
Bifurcation diagram



Instabilities in lasers with an injected signal

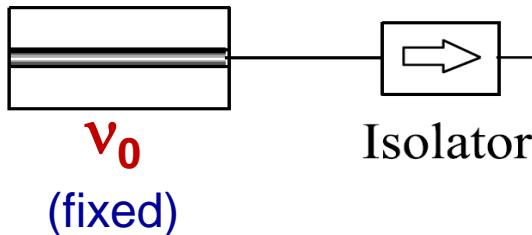
J. R. Tredicce, F. T. Arecchi, G. L. Lippi, and G. P. Puccioni

178 J. Opt. Soc. Am. B/Vol. 2, No. 1/January 1985



Experimental control parameters

Master Laser



Slave Laser

- Photo detector: 9.5GHz
- Oscilloscope: 6 GHz, 20GS/s
- Spectrum analyzer: 72GHz

The **frequency detuning** between the lasers, $\Delta\nu = \nu_s - \nu_0$, is controlled by the slave laser pump current, **I**

When **I** increases:

- Joule heating
- the temperature modifies the cavity refractive index
- decreases the cavity resonance frequency

$$\nu_s = g(\text{Temp}) = f(I)$$

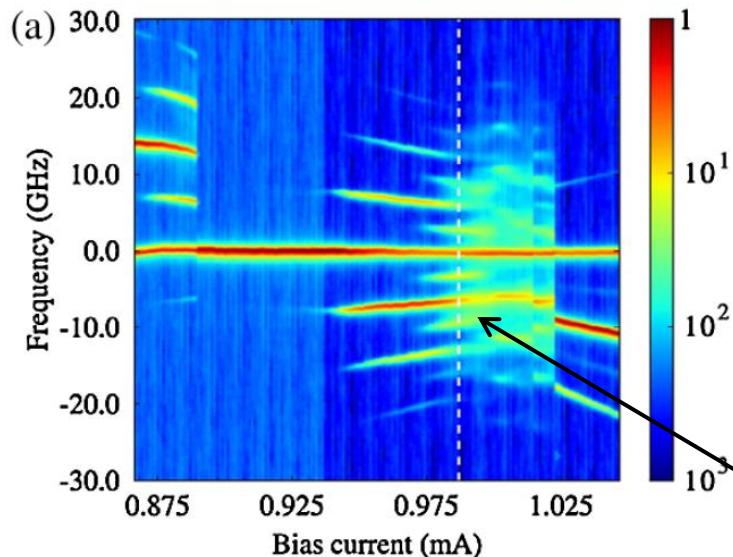
(f approximately linear)

Observed laser output:

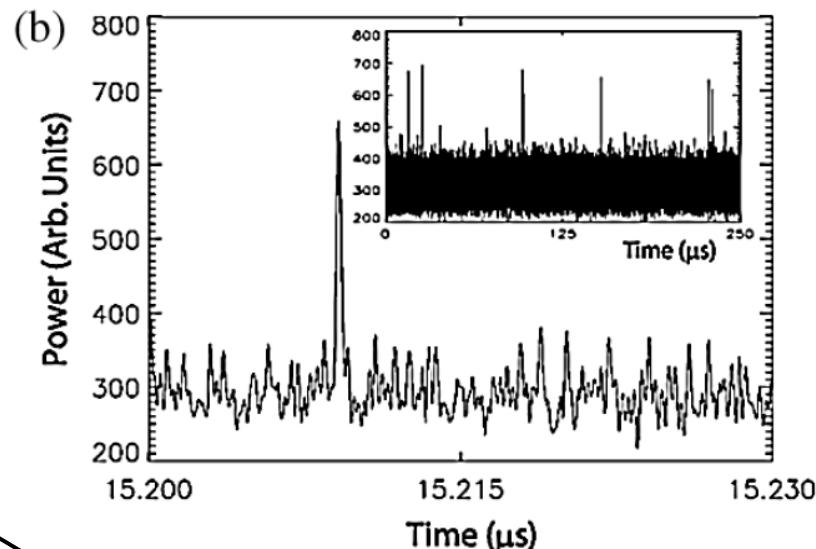
- Intensity time - series (oscilloscope)
- Intensity Fourier spectrum (spectrum analyzer)

Observations

Fourier spectrum of the laser intensity



Time series of the laser intensity



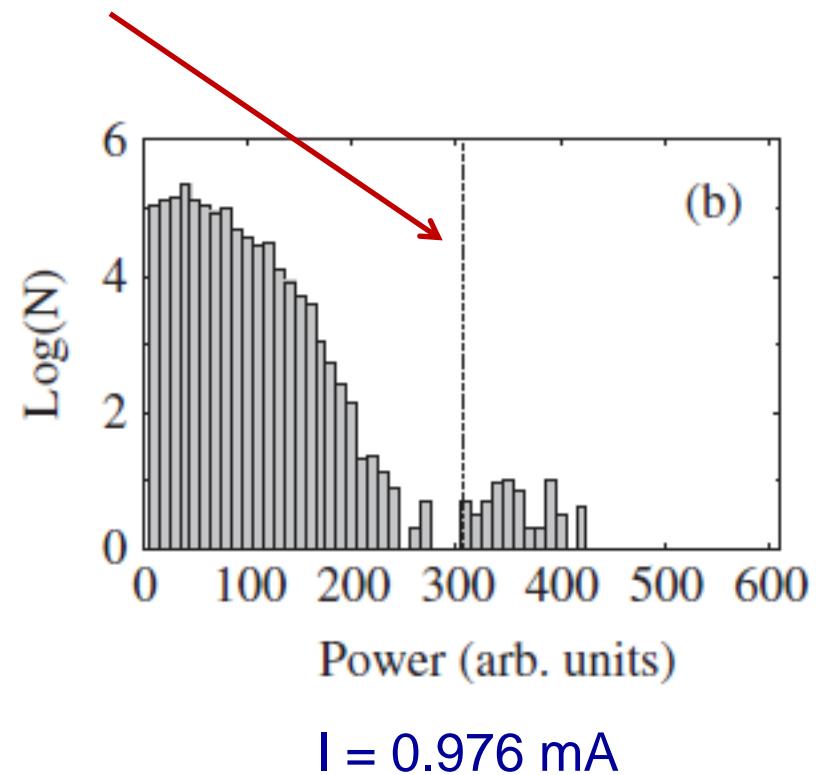
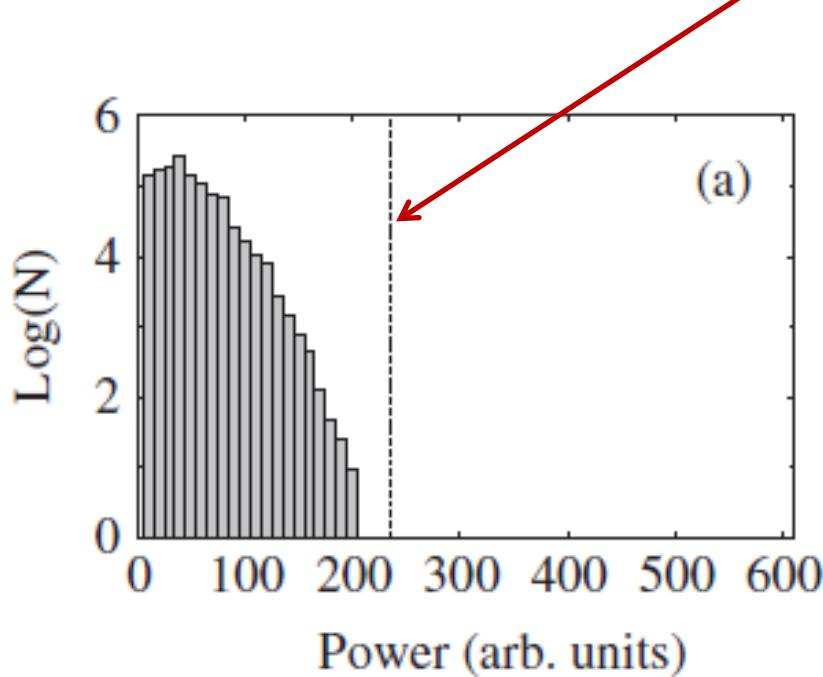
Five regions:

- Beating (independent lasers)
- Locking
- Period 2 of the beat note
- Chaos
- Beating (independent lasers again)

(In the chaotic region,
 $I = 0.976 \text{ mA}$, $\Delta\nu = -1.34 \text{ GHz}$)

Histograms of pulse amplitude

RW threshold = mean value + **8** σ



Governing equations

- Complex field, \mathbf{E} (photon number $\propto |\mathbf{E}|^2$)
- Carrier density, \mathbf{N}

$$\frac{dE}{dt} = \frac{1}{2\tau_p} (1 + i\alpha)(N - 1)E + \underbrace{i\Delta\omega + \sqrt{P_{\text{inj}}}}_{\text{optical injection}} + \underbrace{\sqrt{2\beta_{sp}/\tau_N}\xi(t)}_{\text{spontaneous emission noise}}$$

$$\frac{dN}{dt} = \frac{1}{\tau_N} (\mu - N - N|E|^2)$$

Solitary laser
parameters: α τ_p τ_N μ

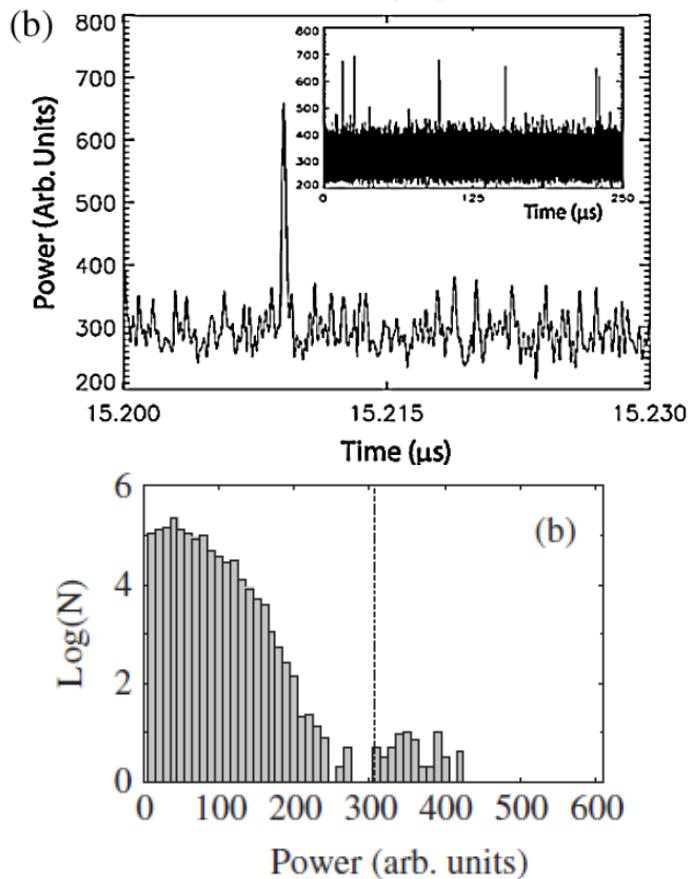
η : injection strength
 $\Delta\omega$: frequency detuning

spontaneous
emission
noise

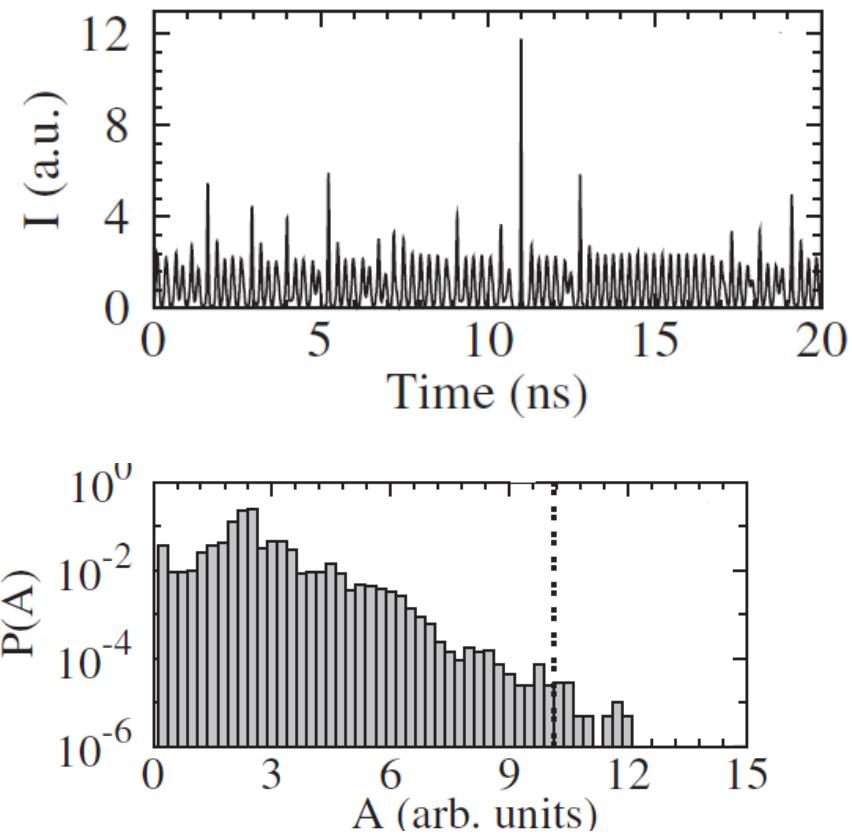
Typical parameter values:
 $\alpha = 3$, $\tau_p = 1 \text{ ps}$, $\tau_N = 1 \text{ ns}$

μ : normalized pump current parameter

Experiments

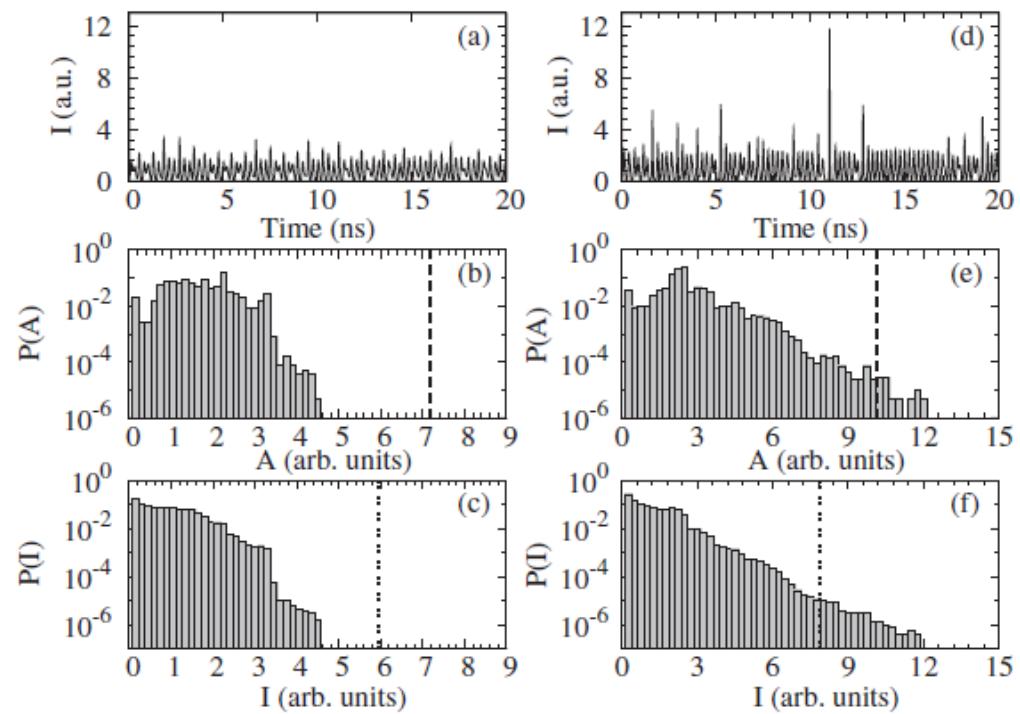
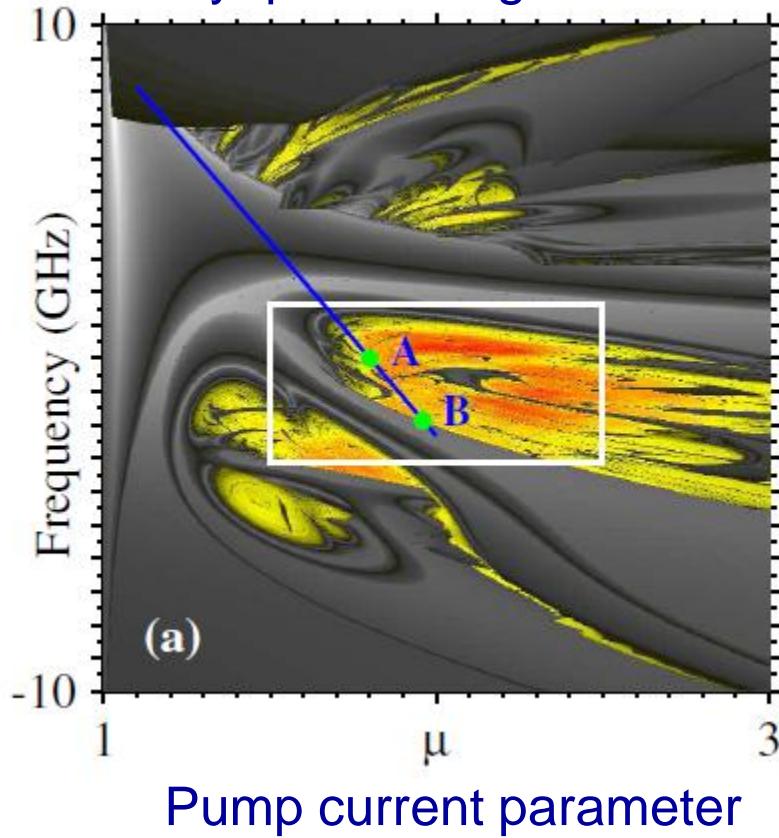


Simulations



Deterministic simulations $(\beta_{sp}=0)$

Lyapunov diagram

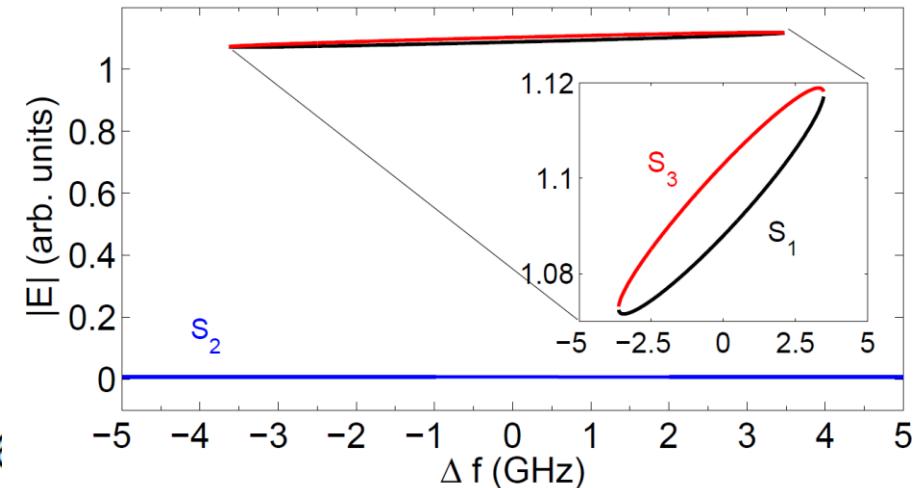
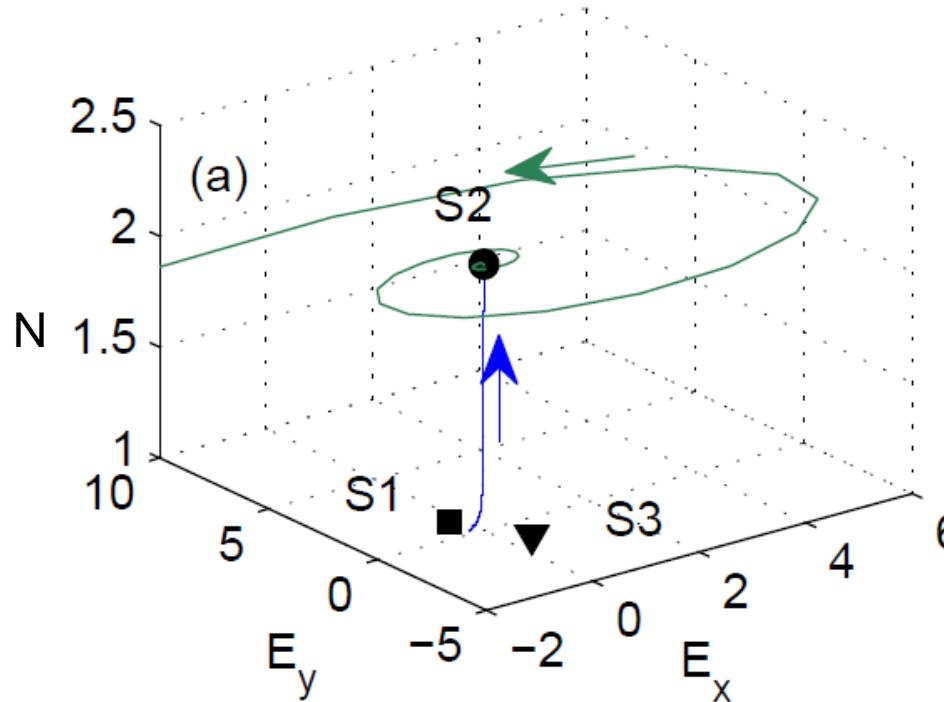


Point A: No RWs

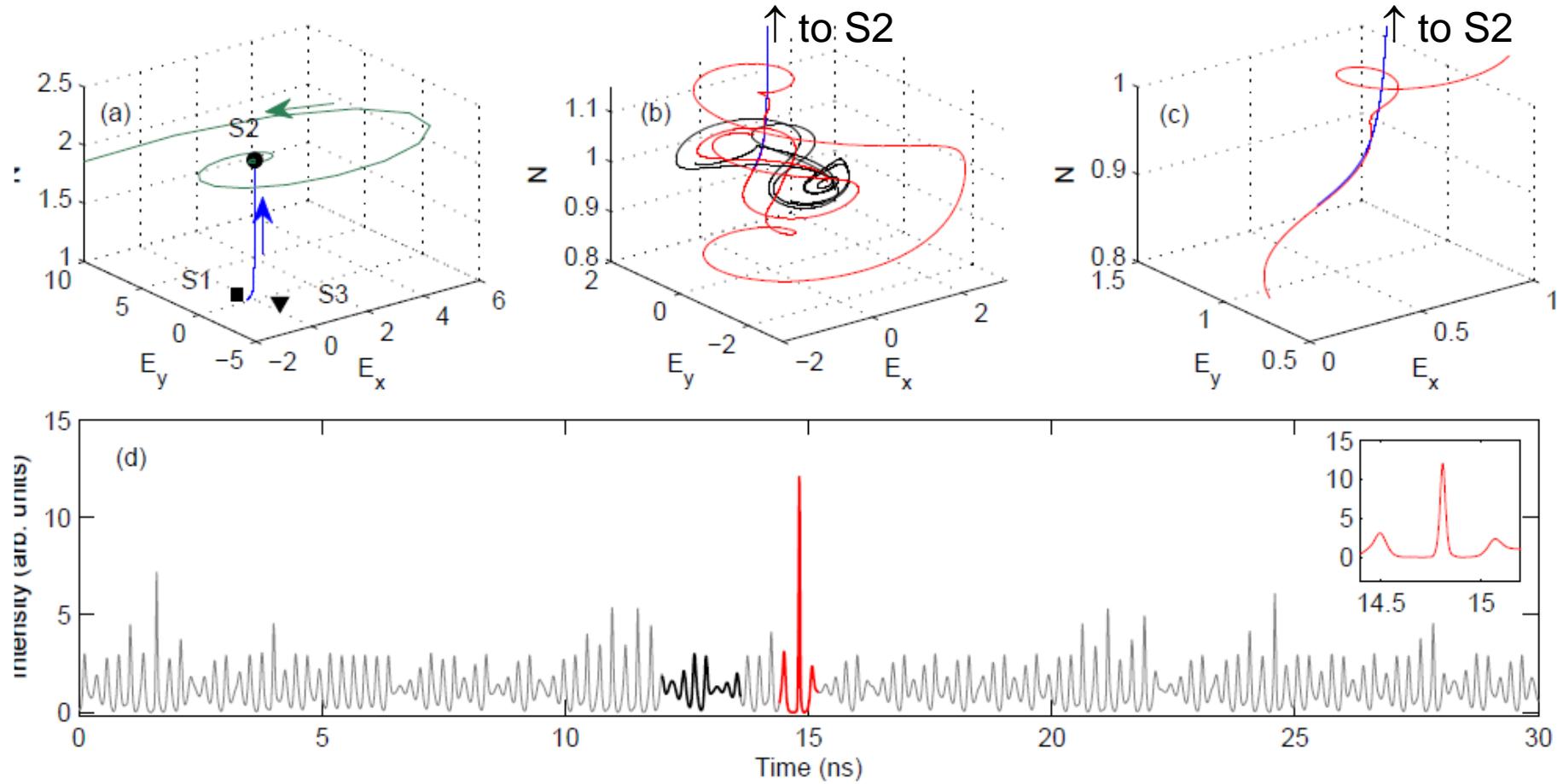
Point B: RWs

What triggers a RW?

Three fixed points in the phase space

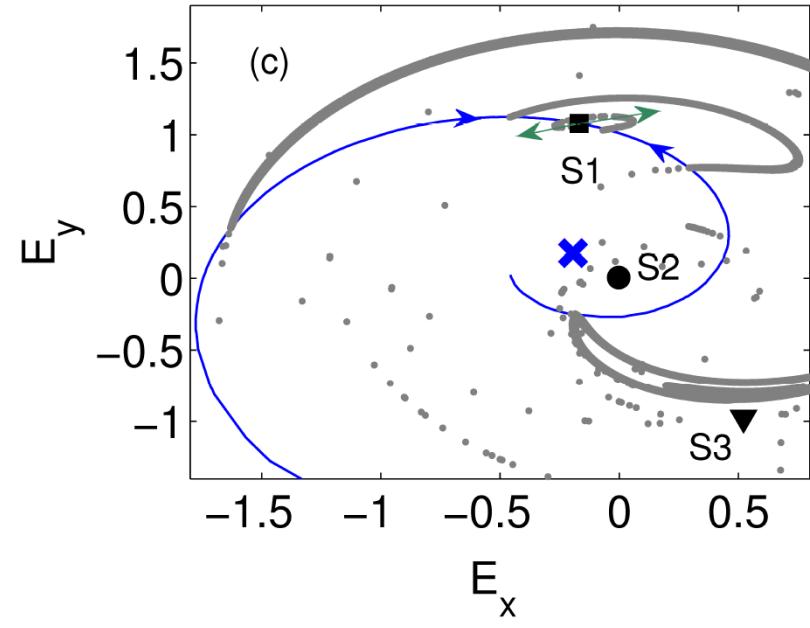
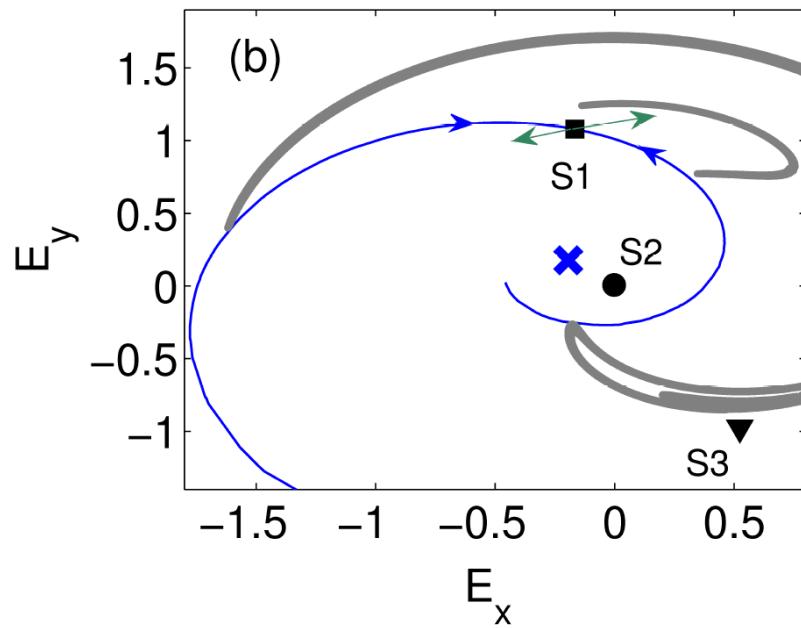


A narrow channel: the RW “door”



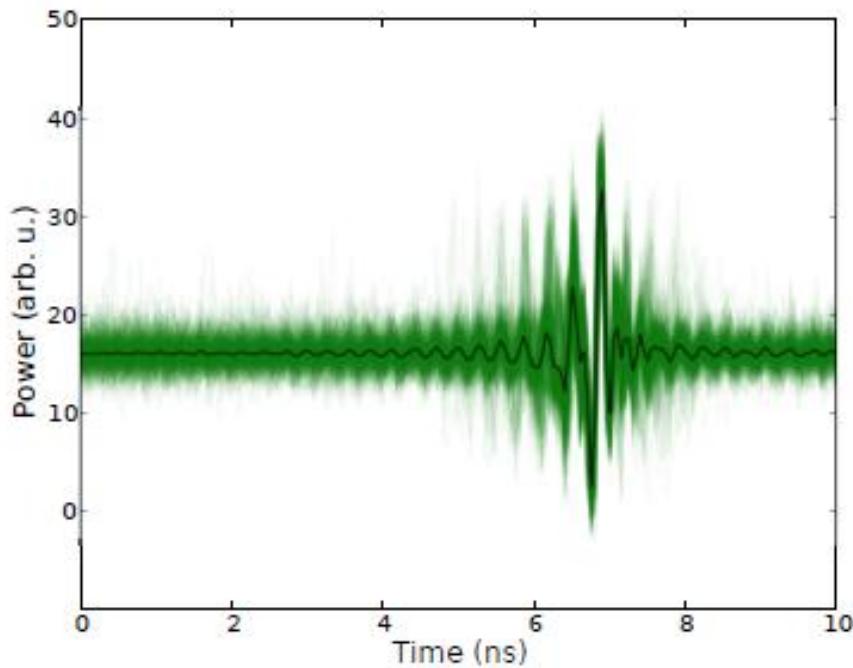
A **RW** is triggered whenever the trajectory closely approaches the stable manifold of S2

Why chaos with RWs and chaos without them?

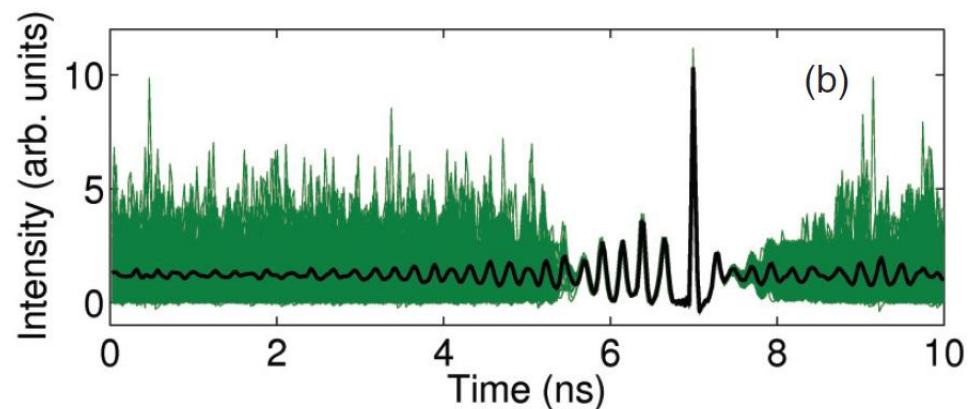


An **external crises-like** process enables access to the region of phase space where the **stable manifold of S2 (x)** is.

Experiments

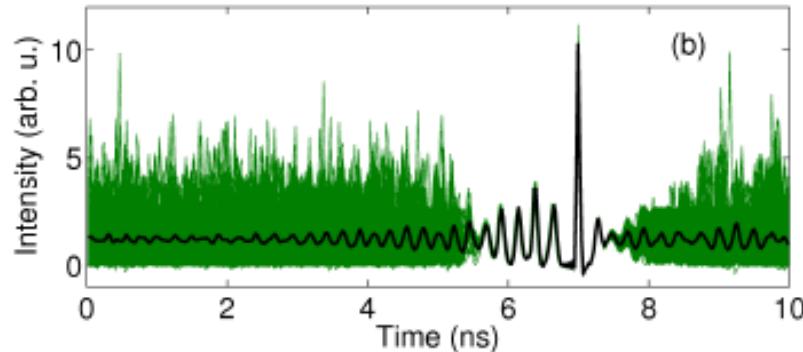


Simulations

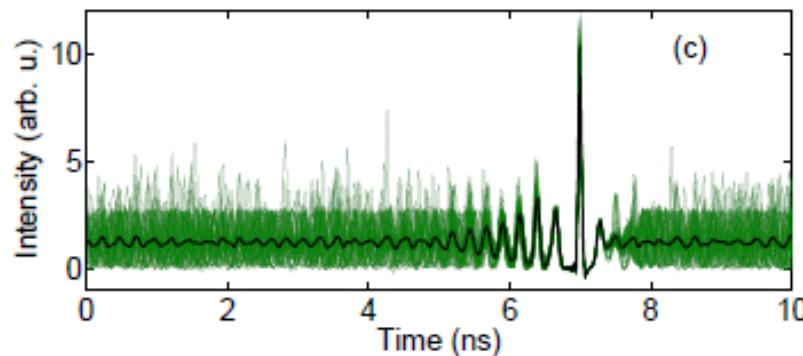


Superposition of 500 time series at the RW peak

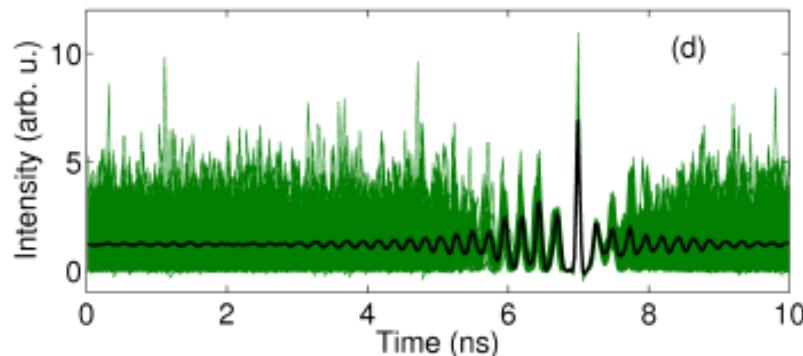
Influence of noise & RW threshold



$\langle H \rangle + 8\sigma$
 $\beta_{sp} = 0$
459 RWs



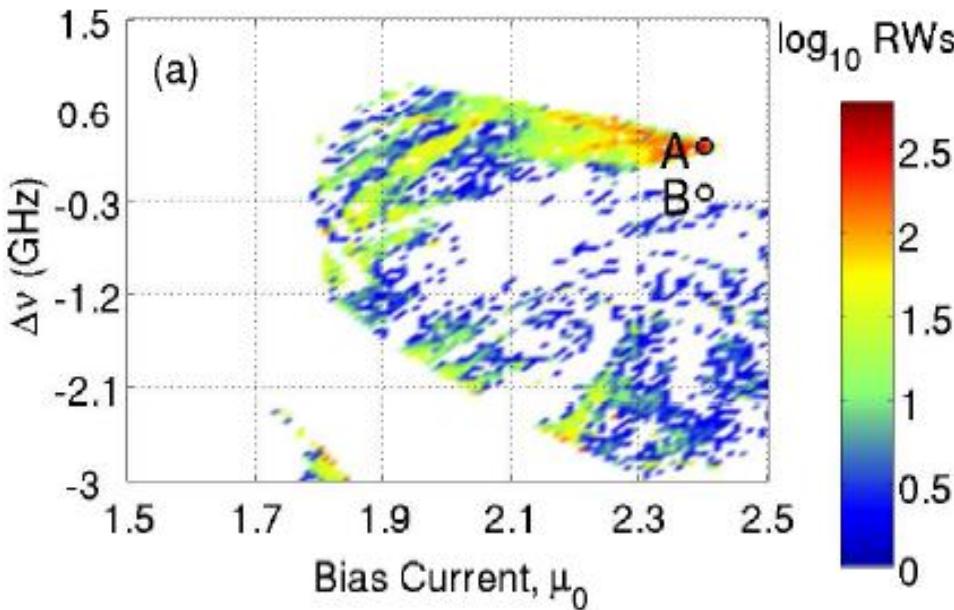
$\langle H \rangle + 8\sigma$
 $\beta_{sp} = 10^{-2}$
53 RWs



$\langle H \rangle + 4\sigma$

Number of RW in the parameter space (pump current, detuning)

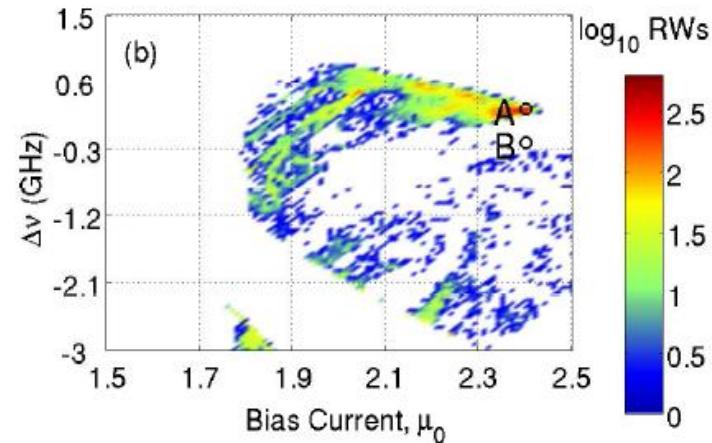
Deterministic RWs ($\beta_{sp}=0$)



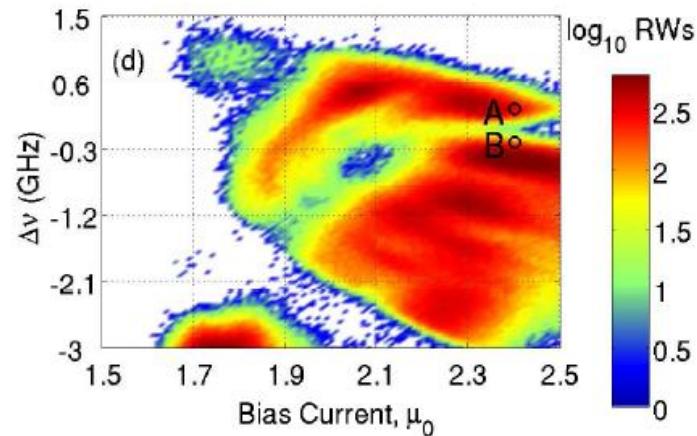
White = No RWs

Weak noise reduces the number
of RWs, but strong noise
increases the number of RWs

Weak noise ($\beta_{sn}=0.0001$)



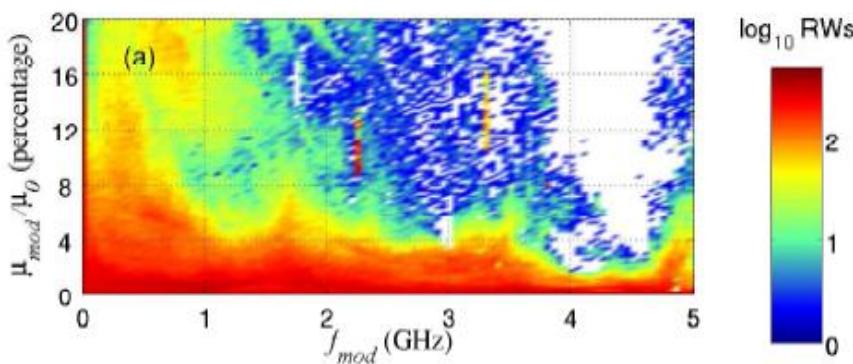
Stronger noise ($\beta_{sp}=0.01$)



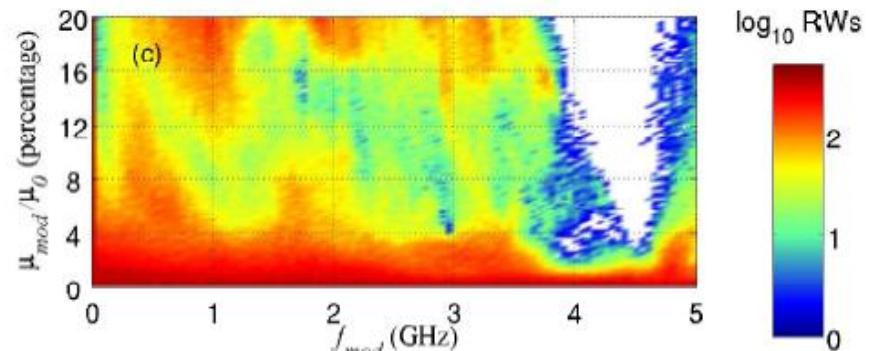
RW control in Point A (deterministic RWs)

$$\mu = \mu_0 + \mu_{\text{mod}} \sin(2\pi f_{\text{mod}} t)$$

$\beta_{\text{sp}}=0$



$\beta_{\text{sp}}=0.01$

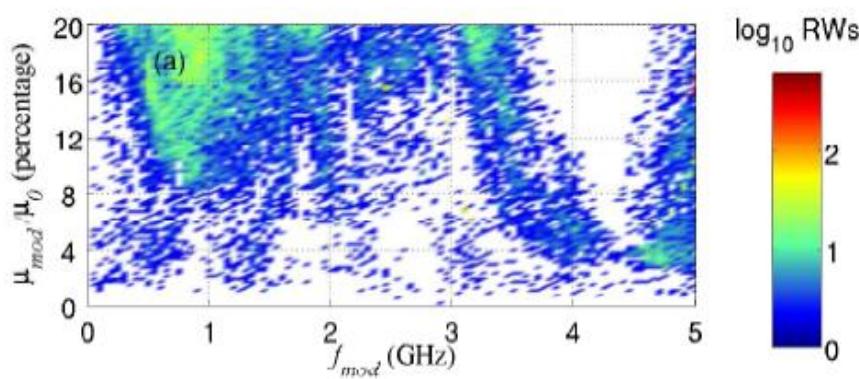


White = No RWs

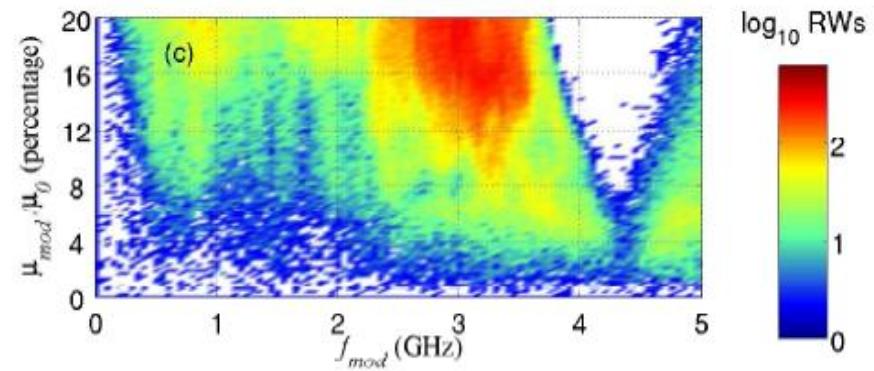
Current modulation with appropriated amplitude and frequency can completely suppress the RWs.

Controlling noise-induced RWs in point B (no deterministic RWs)

$\beta_{sp}=0$



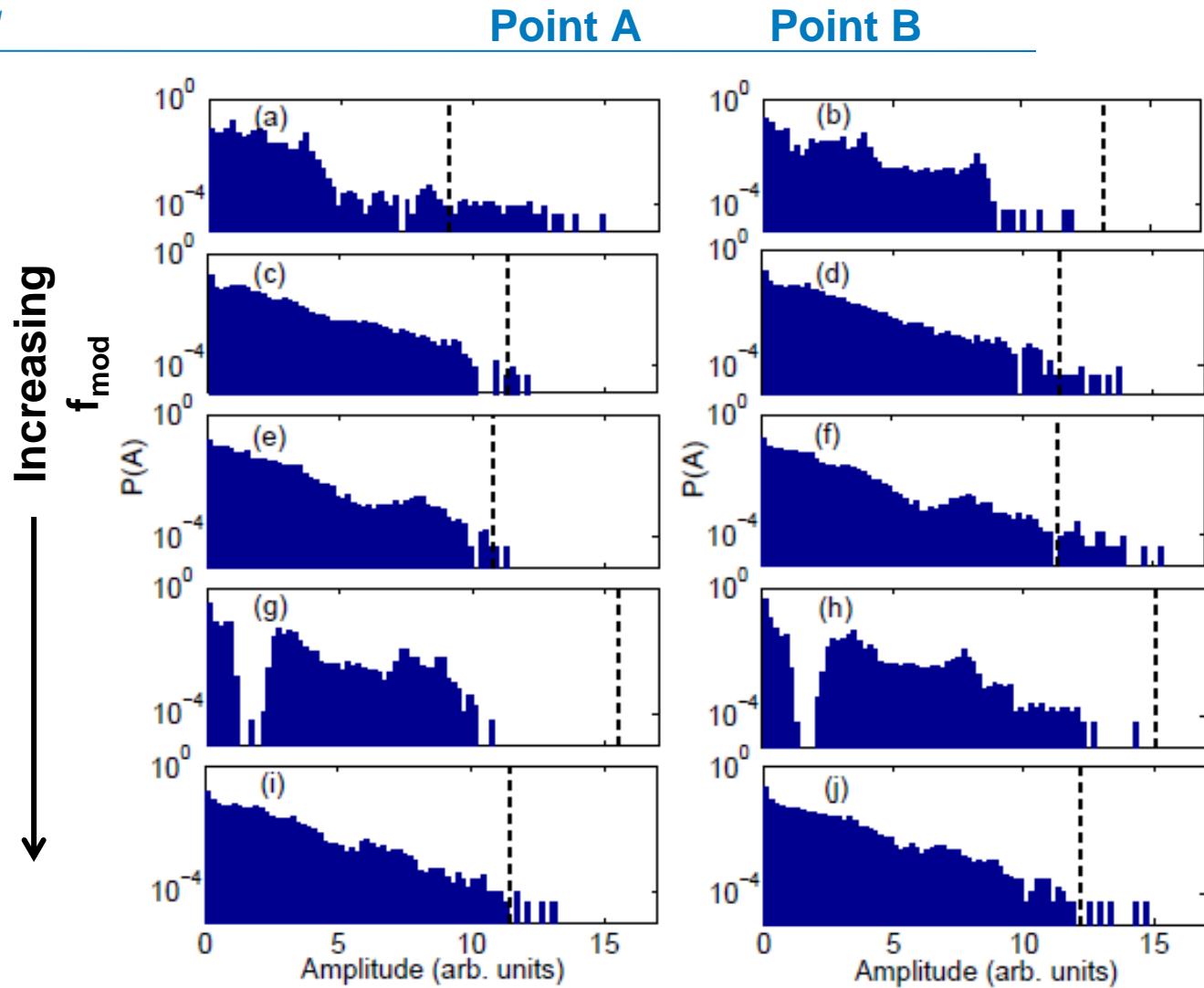
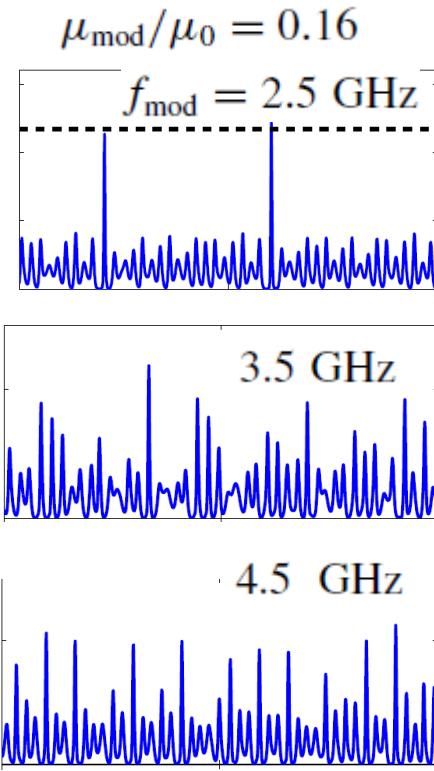
$\beta_{sp}=0.01$



White = No RWs

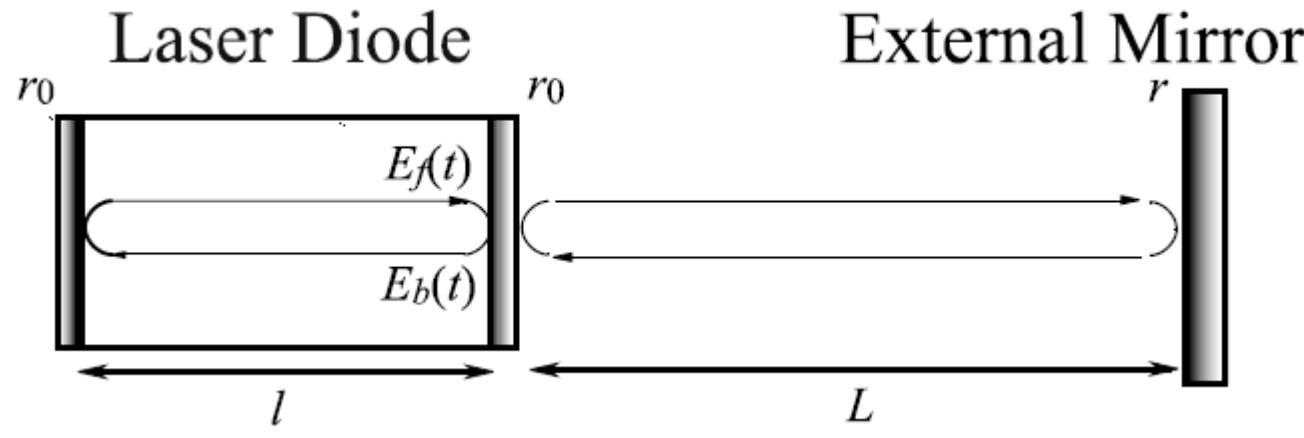
Current modulation provides a “**safe**” parameter region (amplitude, frequency) where no RWs occur.

Histograms of pulse amplitudes



RWs are suppressed because high (but not ultra high) pulses are frequent

Second setup: a laser diode with optical feedback



$$\tau = 2L/c$$

Short EC:
 $f_{\text{ext}} = 1/\tau > f_{\text{ro}}$

$$dE/ds = (1 + i\alpha)NE(s) + \underline{\eta e^{-i w \theta} E(s - \theta)} + \beta \xi,$$

$$TdN/ds = J - N - (1 + 2N)|E(s)|^2.$$

Typical parameter values:

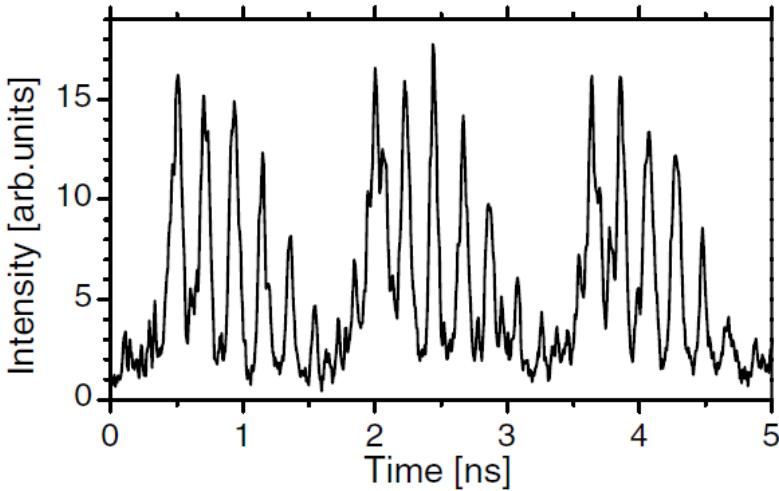
$\alpha = 5, T = 1710, \theta = 70, J = 1.155$

$$s = t / \tau_p \quad \theta = \tau / \tau_p$$

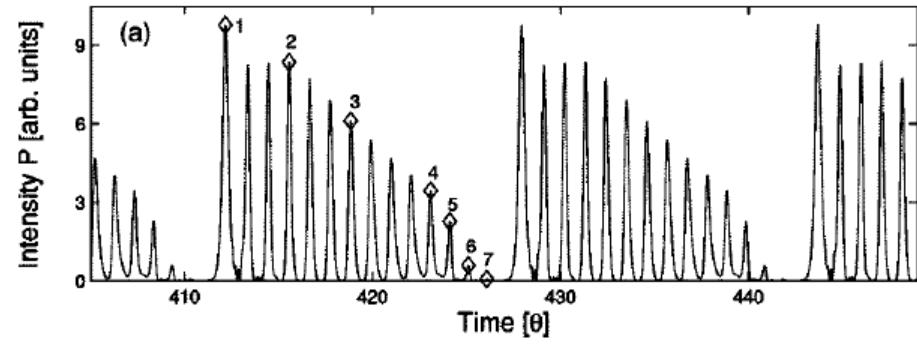
$$T_{\text{RO}} = \pi \sqrt{2T/J} = 171$$

Dynamics: Regular Pulse Packages (RPPs)

■ Experiments



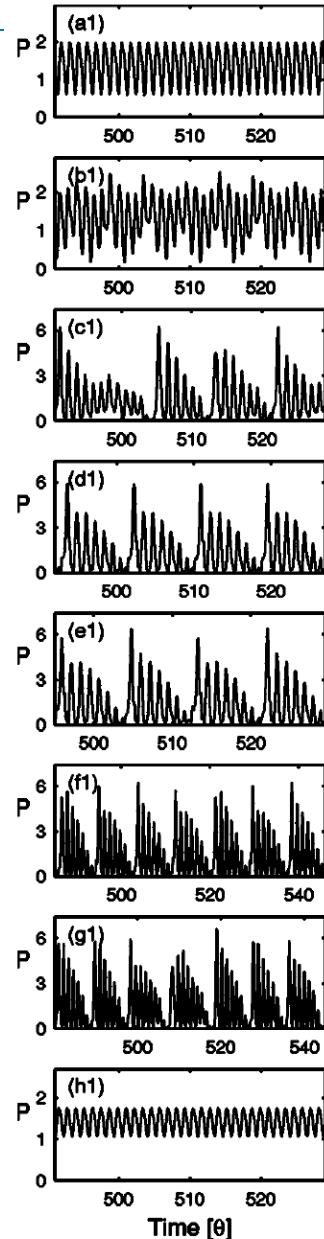
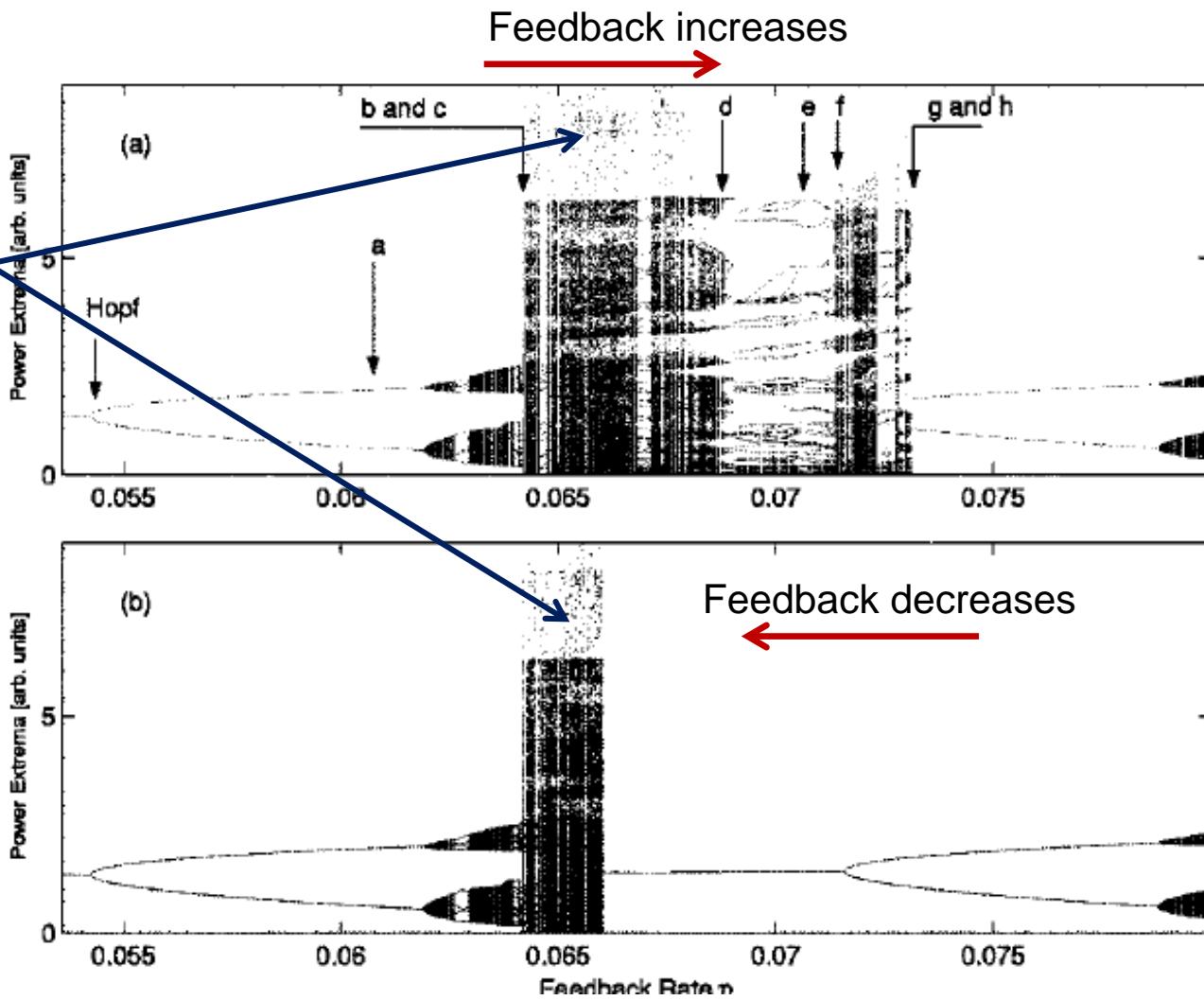
■ Simulations



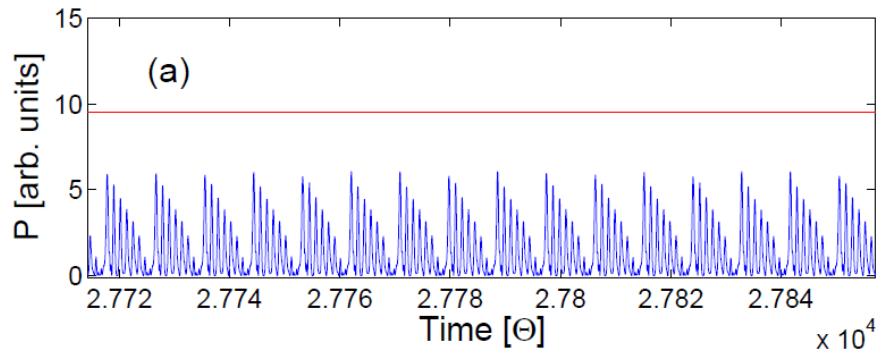
T. Heil et al, PRL 87, 243901 (2001)
A. Tabaka, et al. PRE 70, 036211 (2004)

Numerical bifurcation diagram

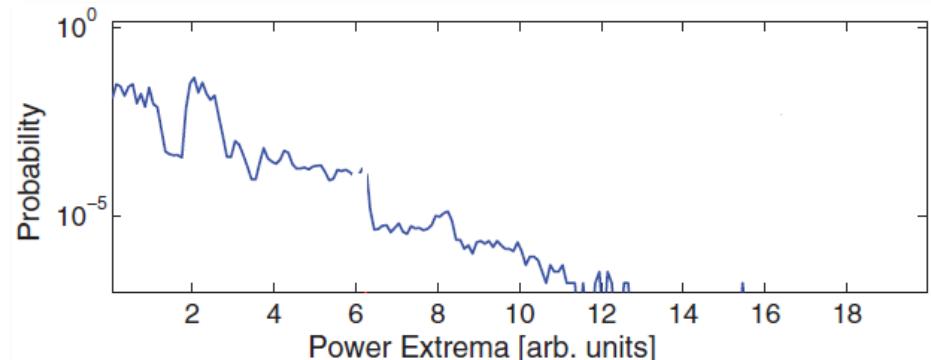
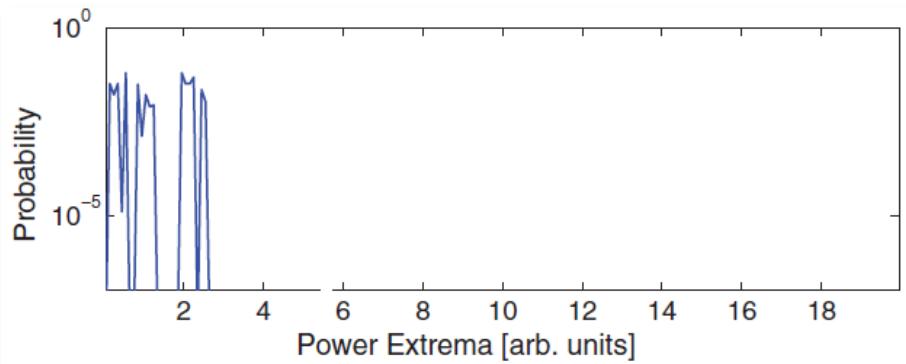
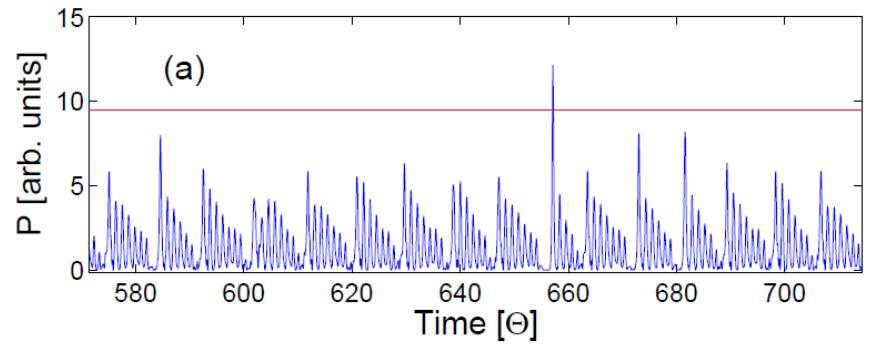
Extreme pulses?



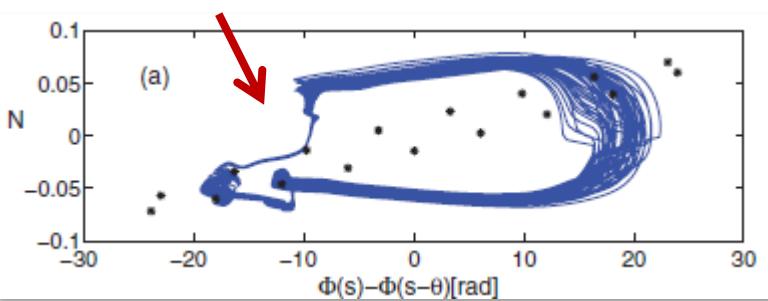
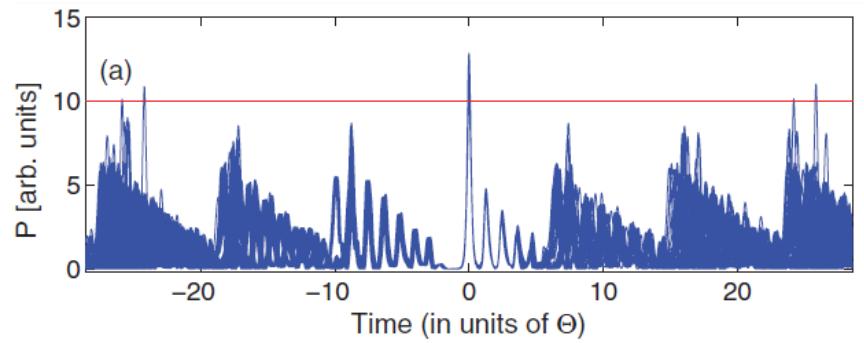
Regular pulse packages



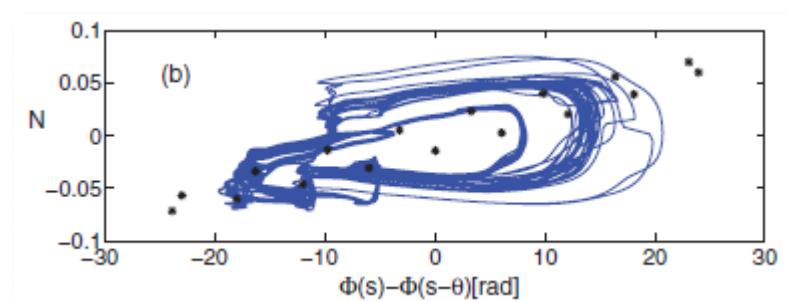
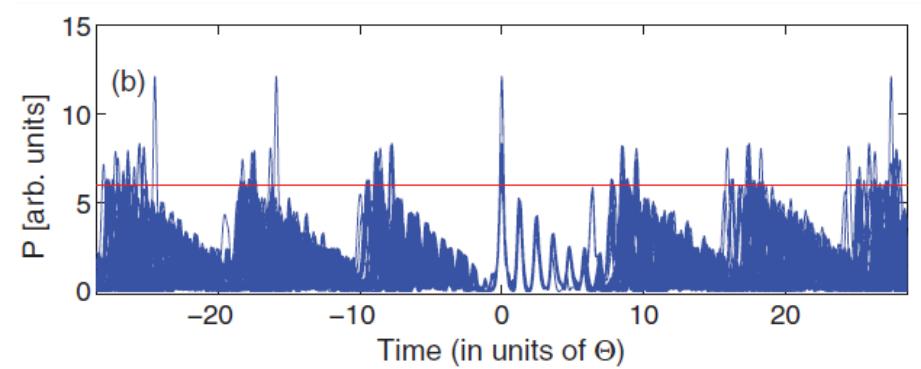
Extreme pulses



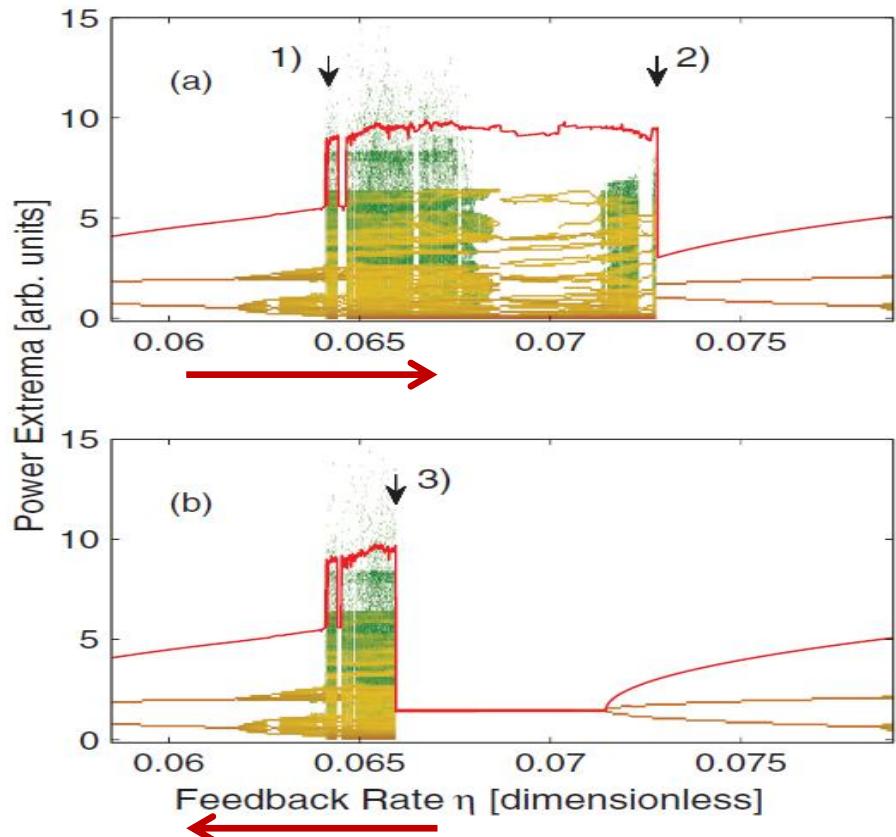
Using a high threshold



Using a lower threshold



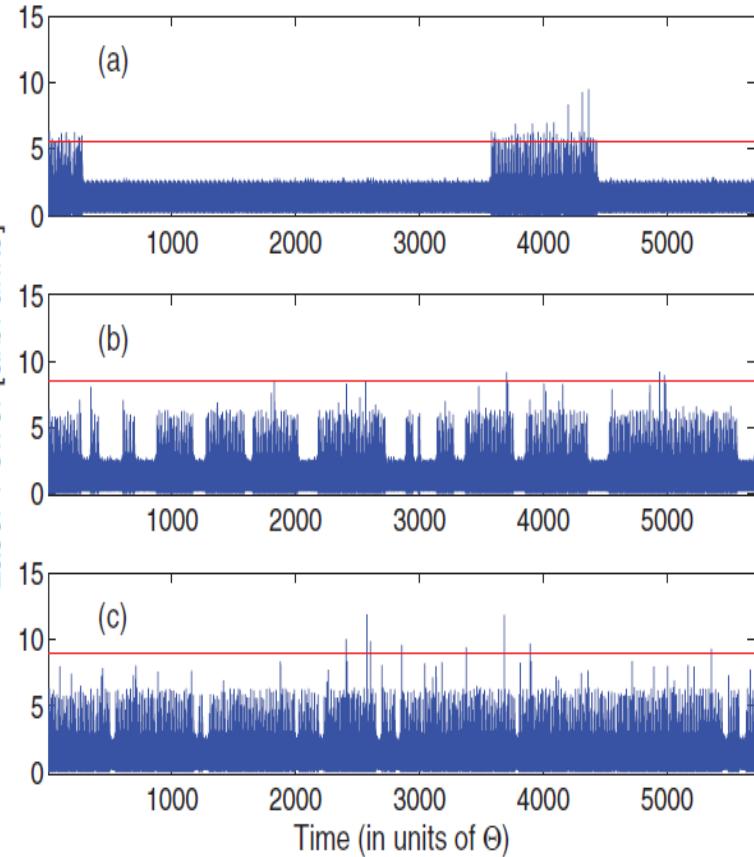
Deterministic intermittency



Feedback increases

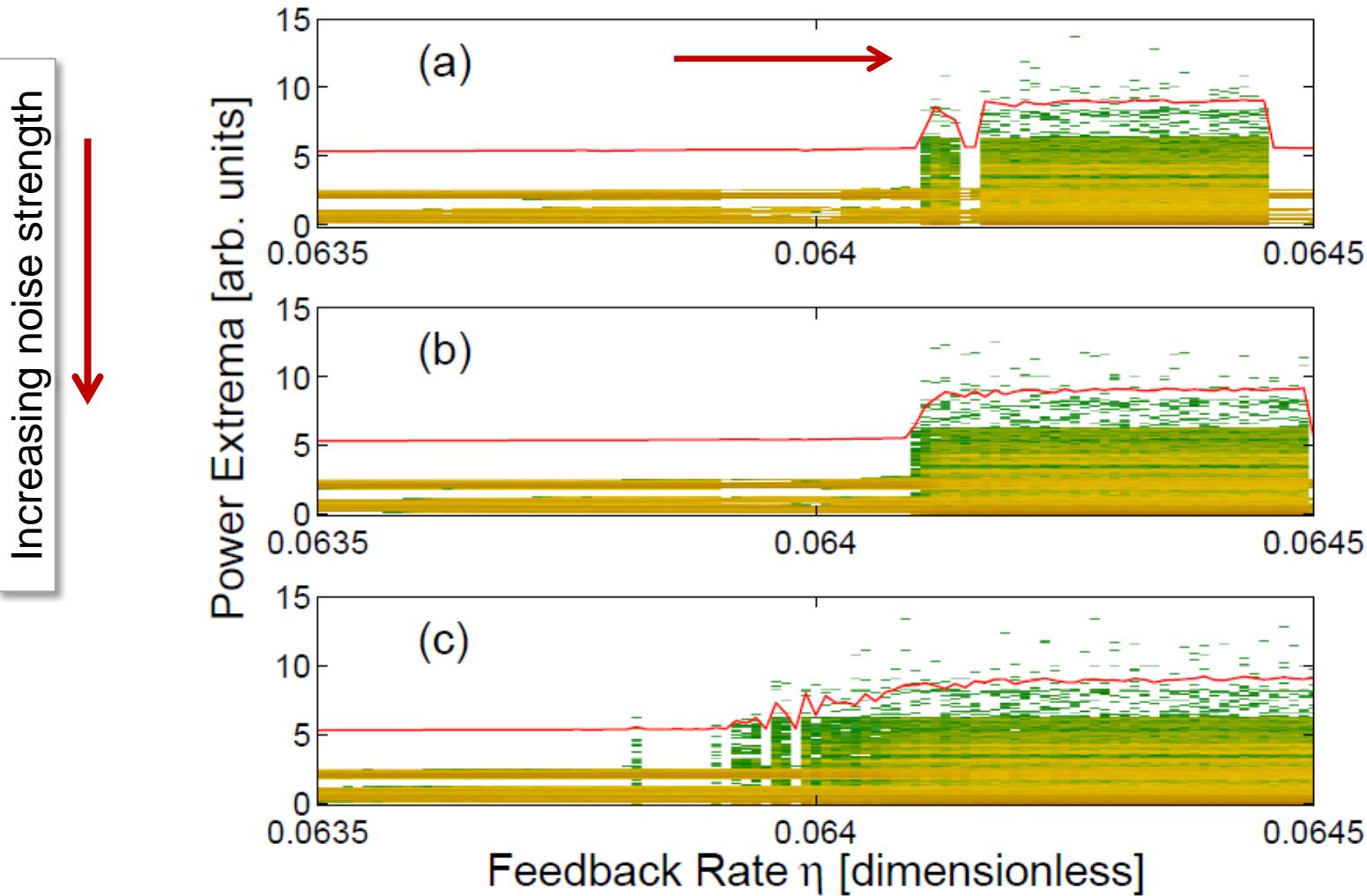
Feedback decreases

At transition 1:



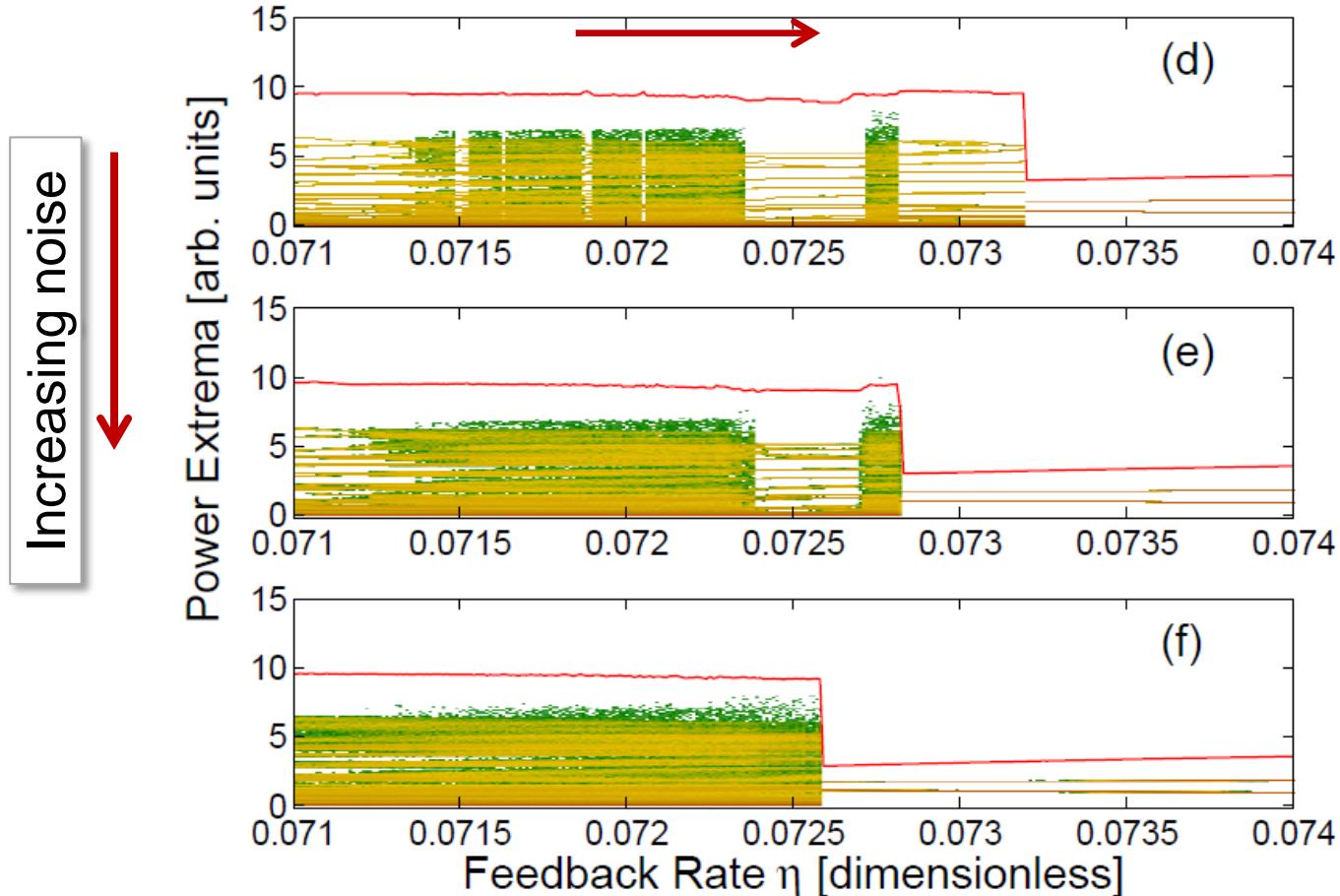
Influence of noise

- Transition 1: noise induced extreme pulses



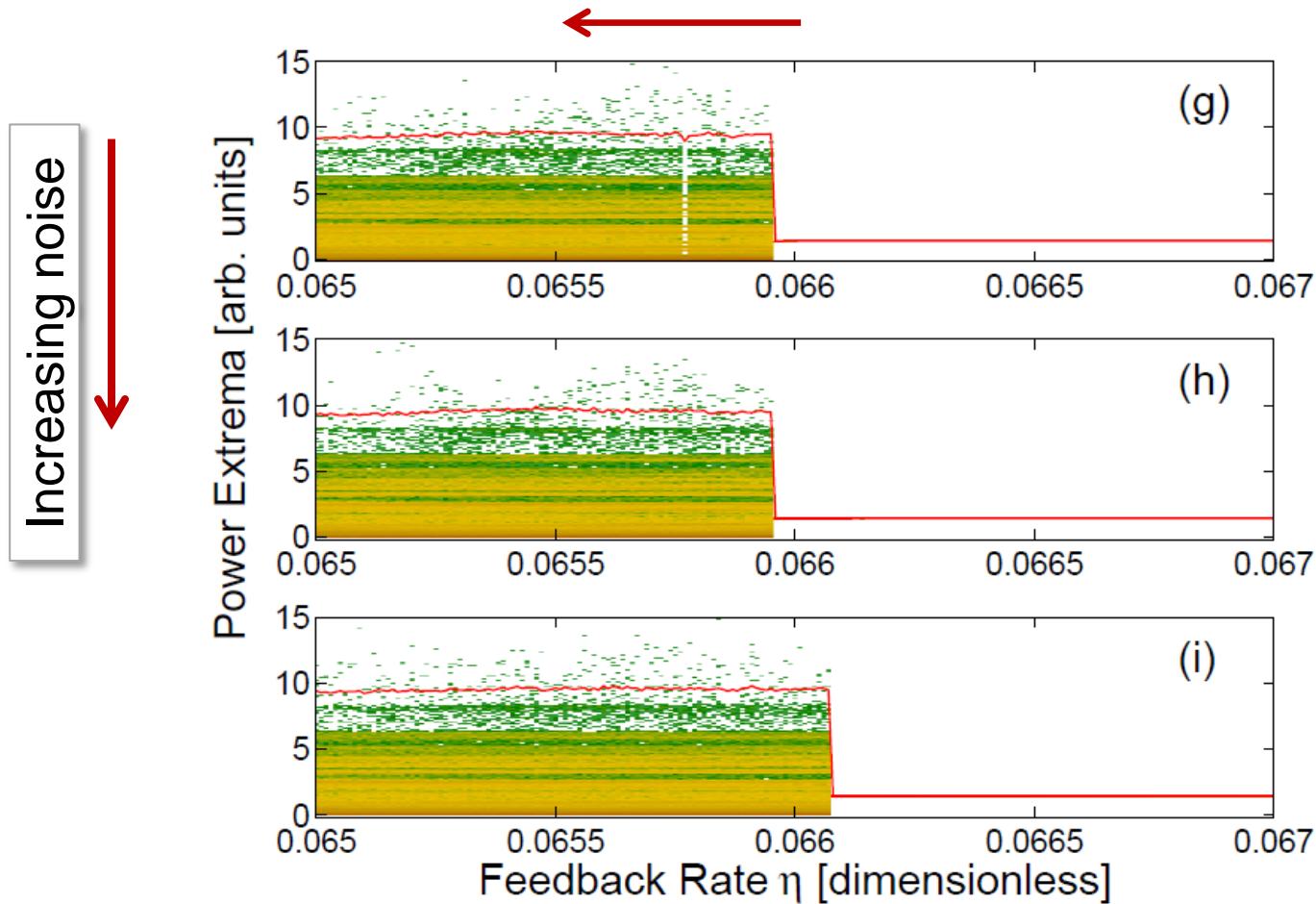
Influence of noise

- Transition 2: noise advances the switching



Influence of noise

- Transition 3: noise advances the switching



Laser with optical injection:

- **Origin** of RWs: an external crises-like process enables access to the region in phase space where RWs can be triggered.
- **Predictability**: RWs can be predicted with some anticipation.
- **Control**: noise and modulation strongly affects their probability of occurrence.

Lasers with optical feedback:

- Similar results, intermittency is the route to extreme pulses.

Ongoing work: modulation-induced periodicity in the RW waiting times?



THANK YOU FOR YOUR ATTENTION !

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