

Extreme pulses in optically injected semiconductor lasers: characterization, prediction, and control

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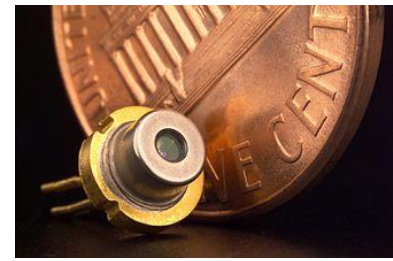
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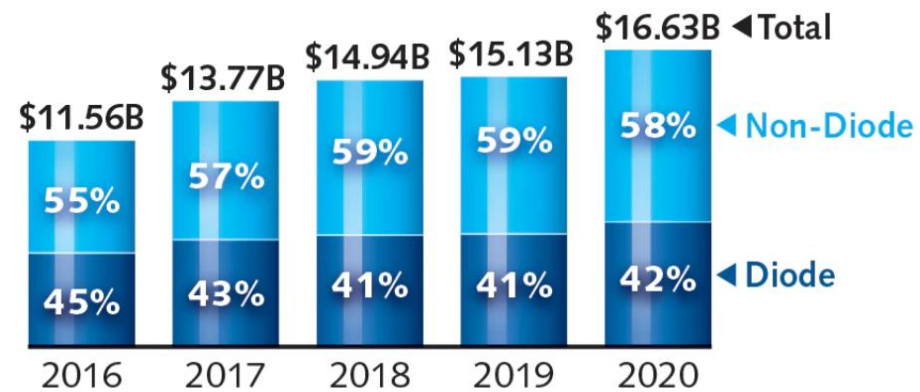


Extreme Waves, MPIPKS, Dresden, August 28, 2023

Semiconductor lasers play a crucial role in photonic technologies

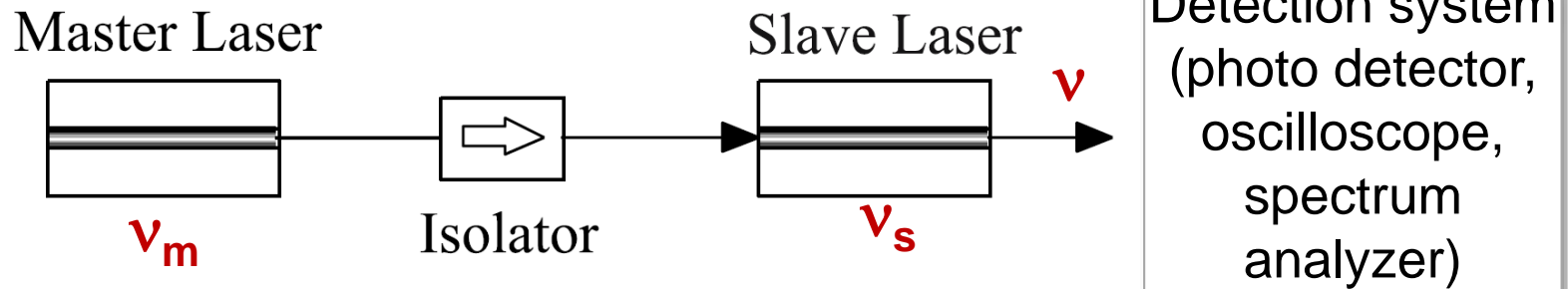


- Inexpensive, compact, efficient
- Emit a wide range of wavelengths (optical communications, biomedical applications),
- Emit a wide range of powers (μ Ws-KWs).

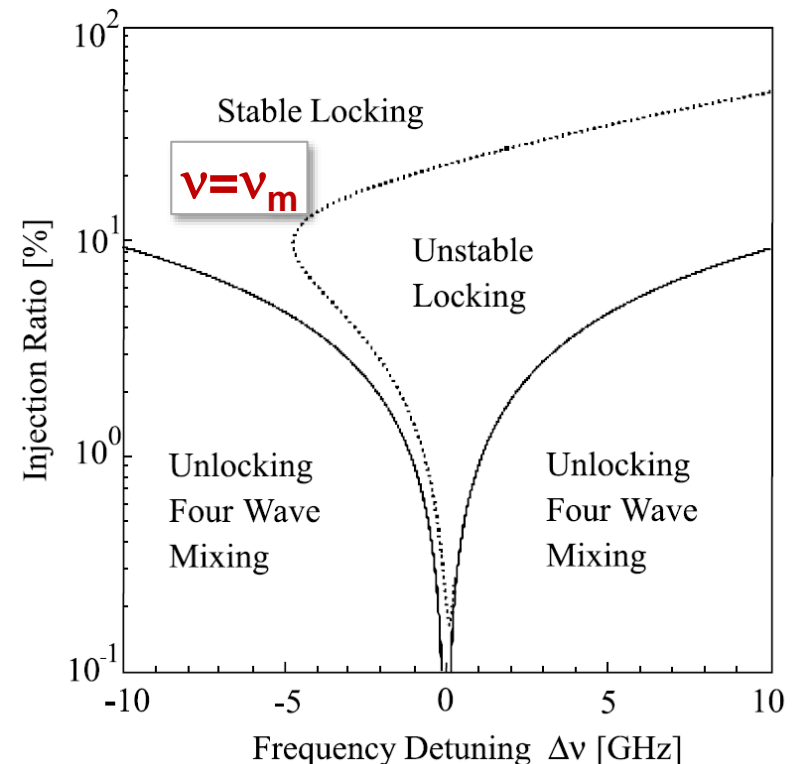


Source: Strategies Unlimited

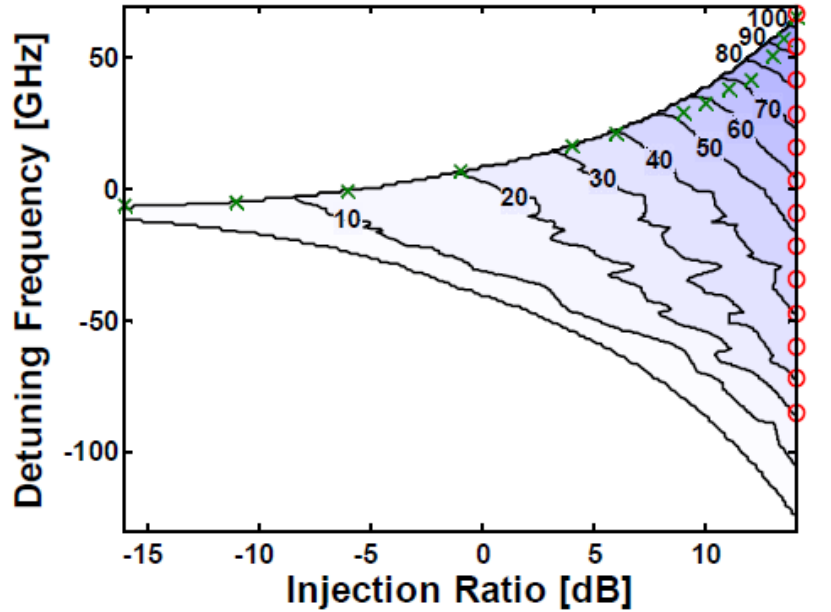
Optically injected lasers are nonlinear dynamical systems



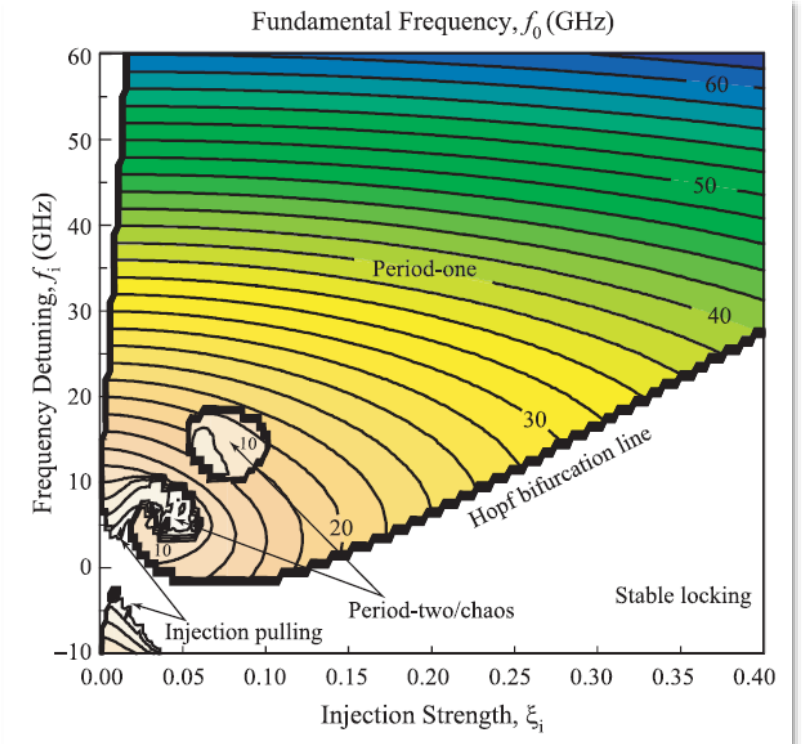
- Two Parameters:
 - Injection ratio
 - Frequency detuning $\Delta\nu = \nu_s - \nu_m$
- Dynamical regimes:
 - Stable locking (cw output)
 - Periodic oscillations
 - Chaos
 - Beating (no interaction)



Injection locking increases the modulation bandwidth; outside the locking region: intensity oscillations

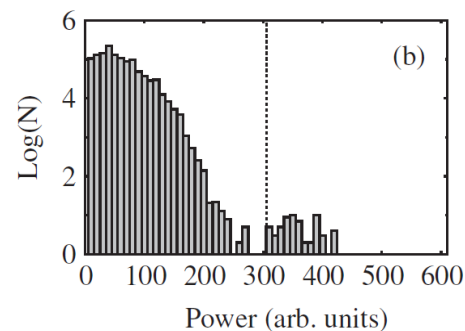
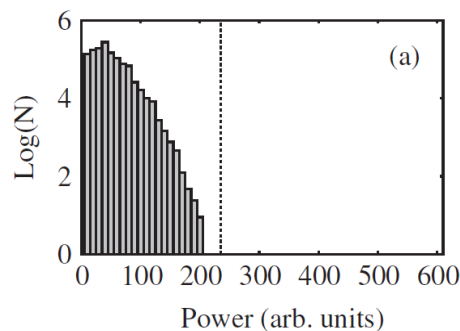
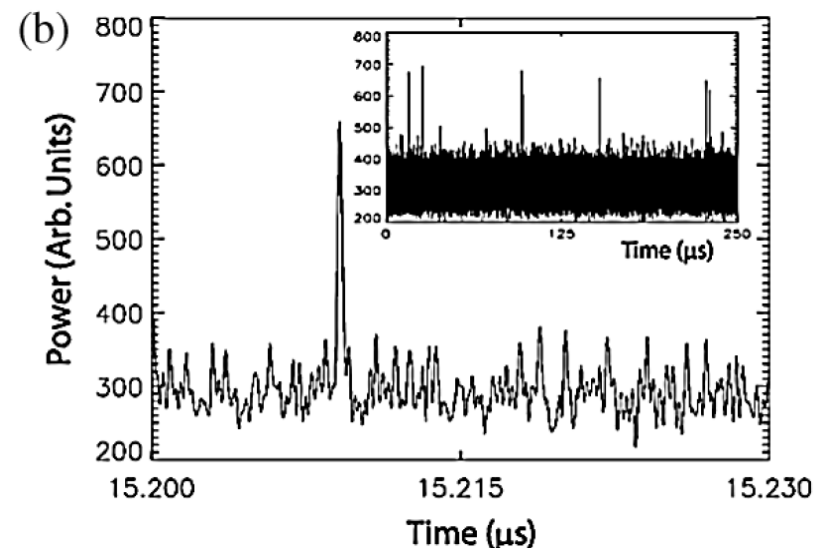
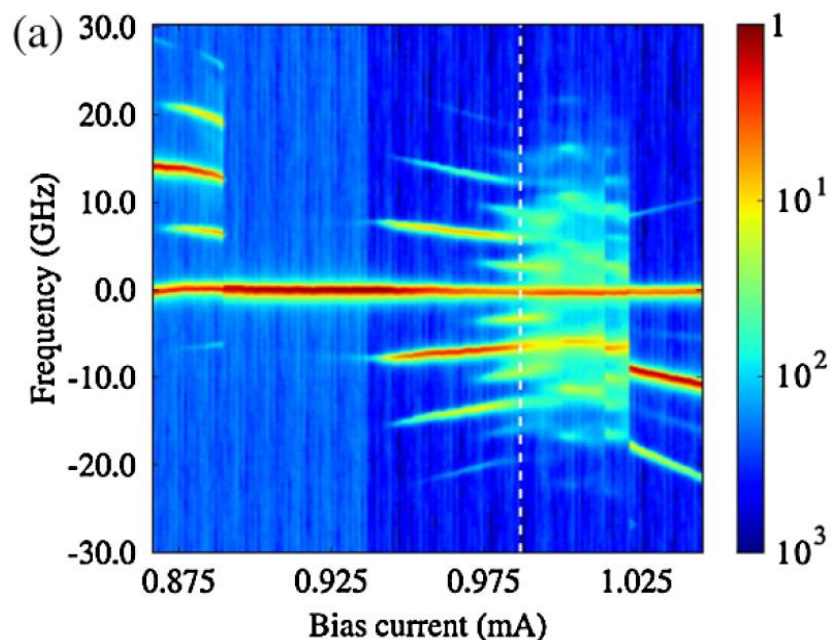


E. K. Lau et. al, Opt. Express 16, 6609 (2008)



S-C Chan et. al, Optics Express 15, 14921 (2007)

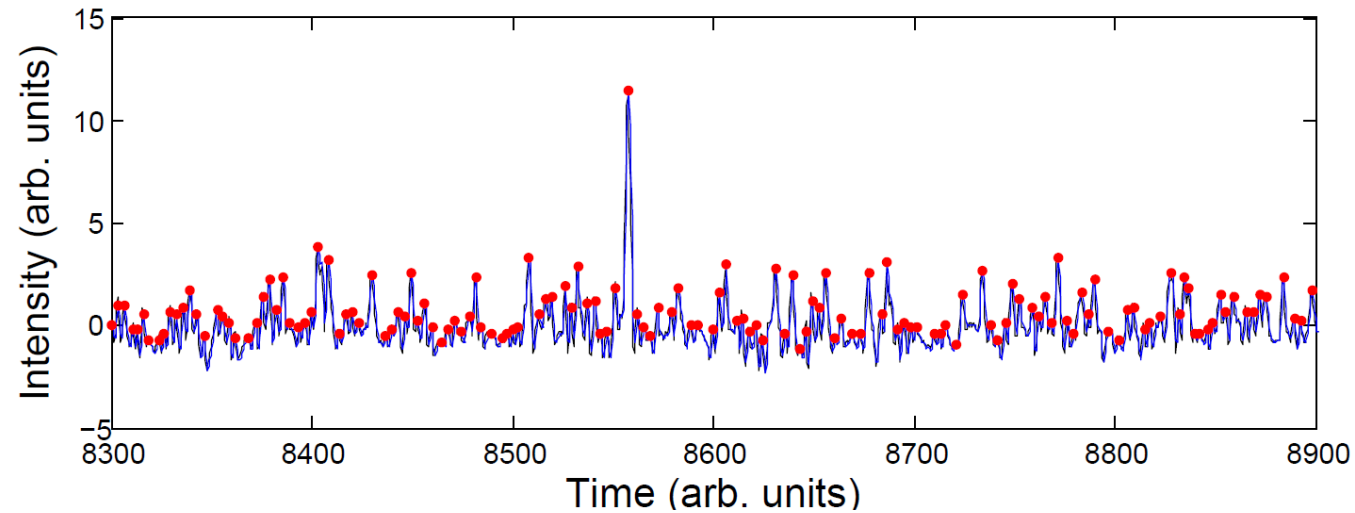
In a narrow parameter region (outside the locking region): extreme pulses



Extreme pulse:
amplitude above $\langle I \rangle + 8 \sigma$

C. Bonatto et al, PRL 107, 053901 (2011), Optics & Photonics News February 2012, Research Highlight in Nature Photonics DOI:10.1038/nphoton.2011.240

Questions



In our system:

- Which mechanisms induce extreme pulses?
- Role of noise?
- Can they be suppress?
- Can they be generated “on demand”?
- Can they be predicted?

Governing equations

Complex field, **E** –Laser intensity $\sim |E|^2$

Carrier density, **N**

$$\frac{dE}{dt} = \frac{1}{2\tau_p} (1 + i\alpha)(N - 1)E + \underbrace{i\Delta\omega + \sqrt{P_{inj}}}_{\text{optical injection}} + \underbrace{\sqrt{2\beta_{sp} / \tau_N} \xi(t)}_{\text{spontaneous emission noise}}$$

$$\frac{dN}{dt} = \frac{1}{\tau_N} (\mu - N - N|E|^2)$$

optical injection
 η : injection strength
 $\Delta\omega = \omega_s - \omega_m$: detuning

spontaneous emission noise

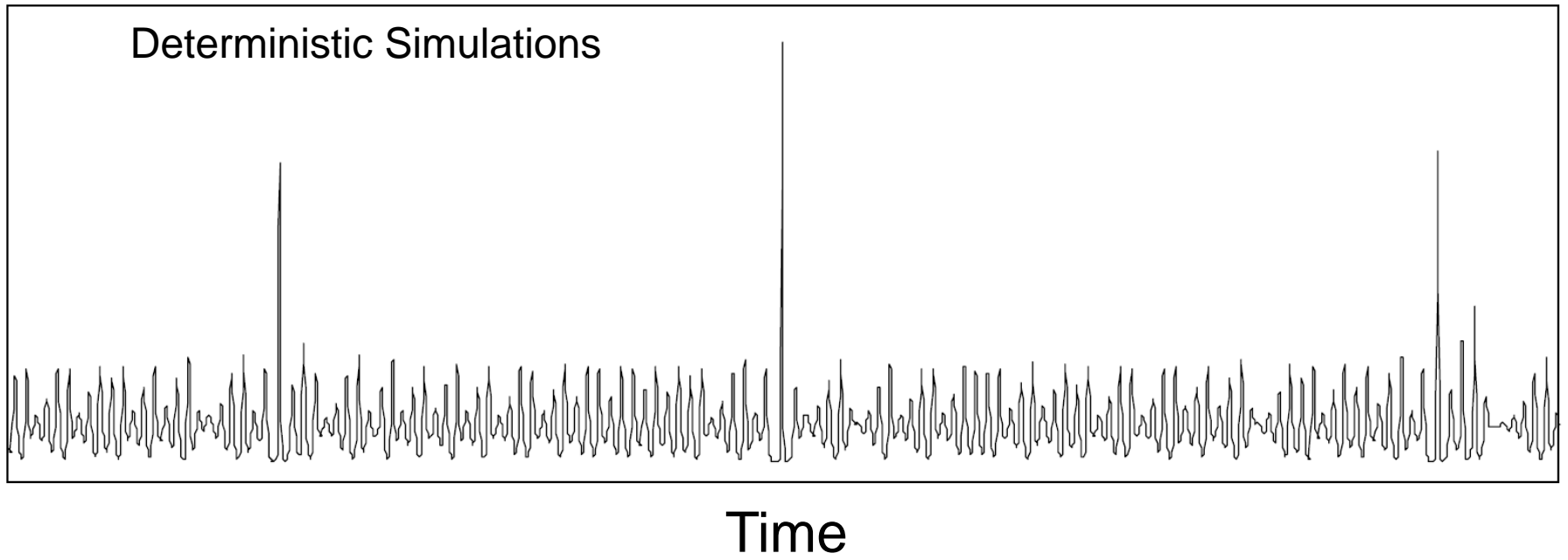
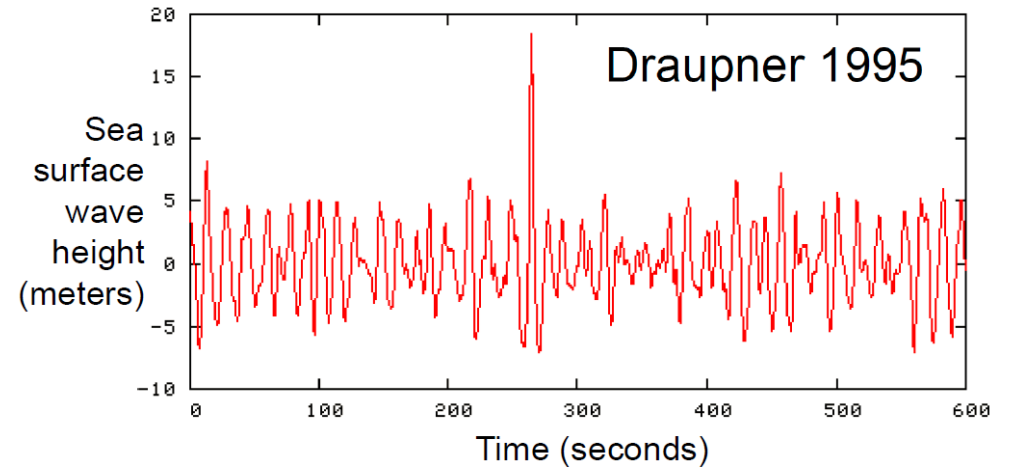
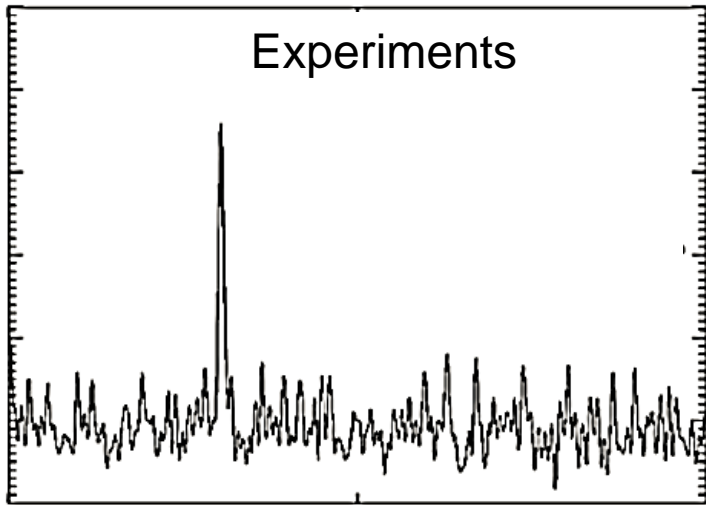
Solitary laser parameters: α τ_p τ_N μ

μ : normalized pump current parameter

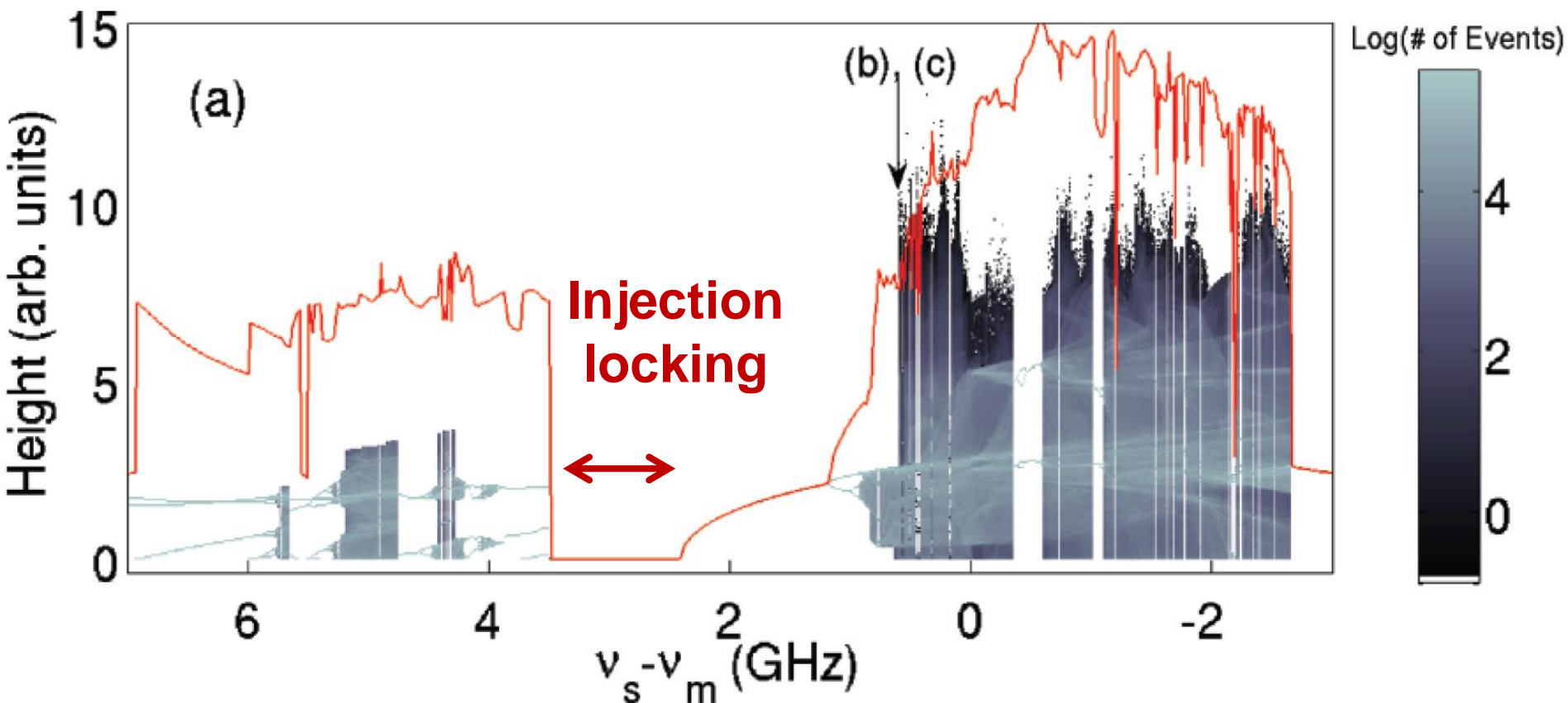
Typical parameter values:

$\alpha = 3$, $\tau_p = 1$ ps, $\tau_N = 1$ ns

These **OD** rate-equations provide good qualitative agreement with the observed intensity dynamics.

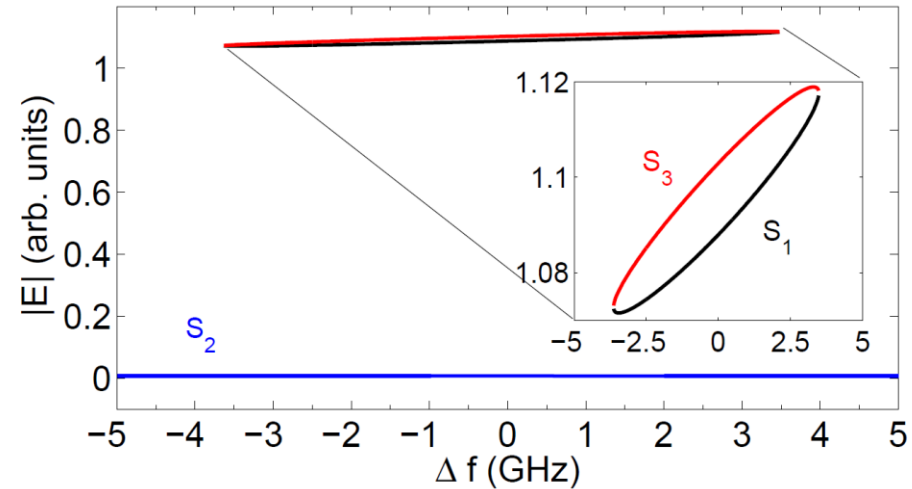
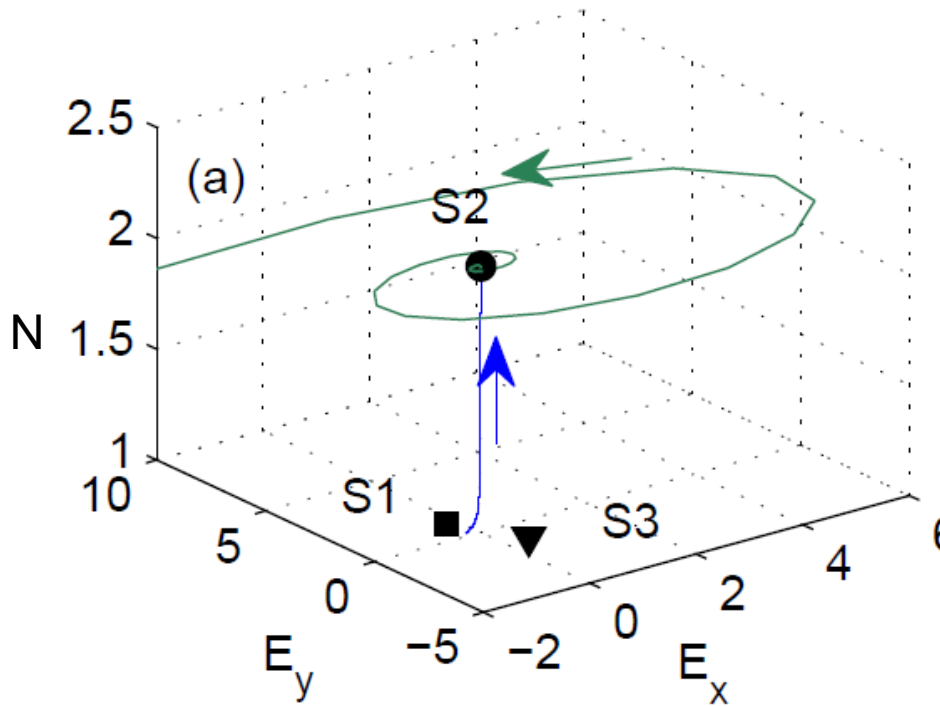


Bifurcation diagram: in color code: $\log(\# \text{ of pulses})$

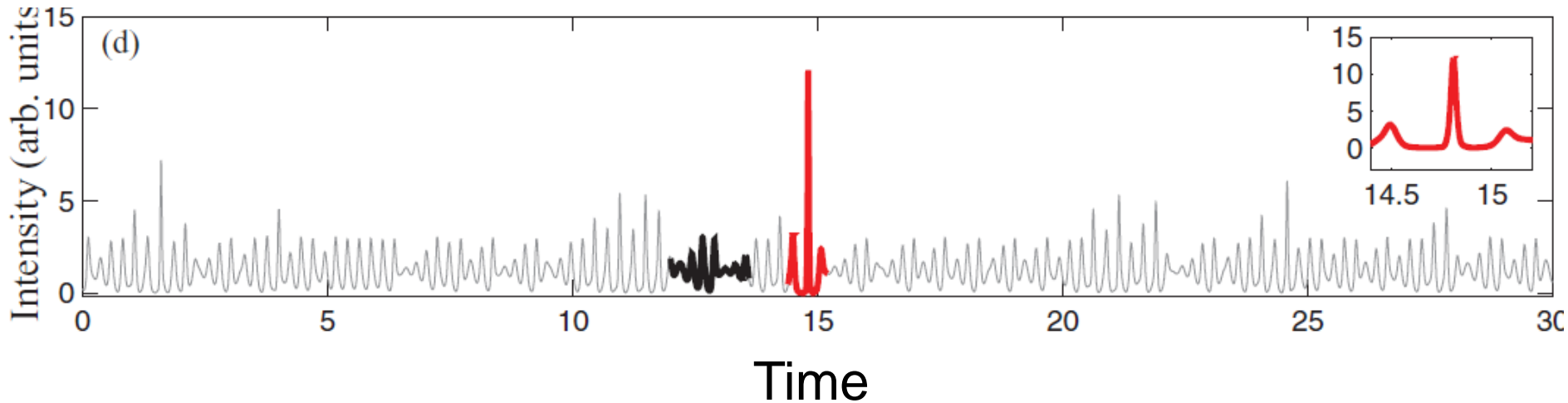
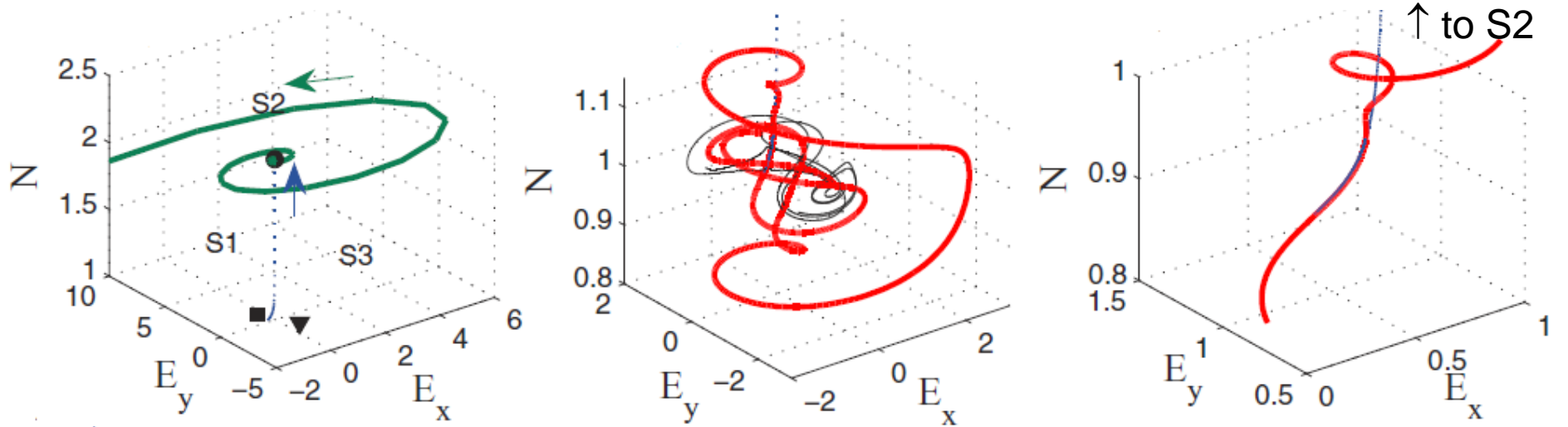


J. Zamora-Munt et al. PRA (2013).

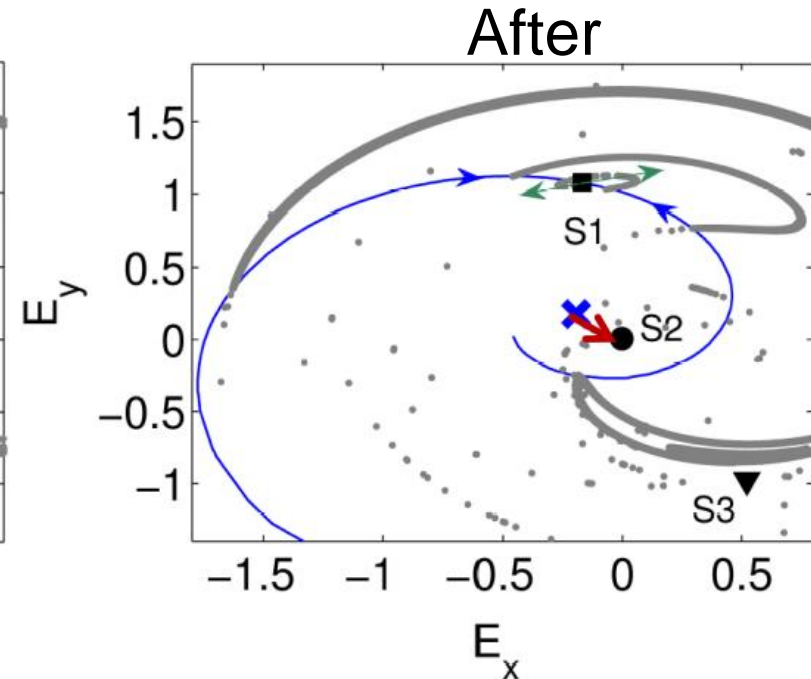
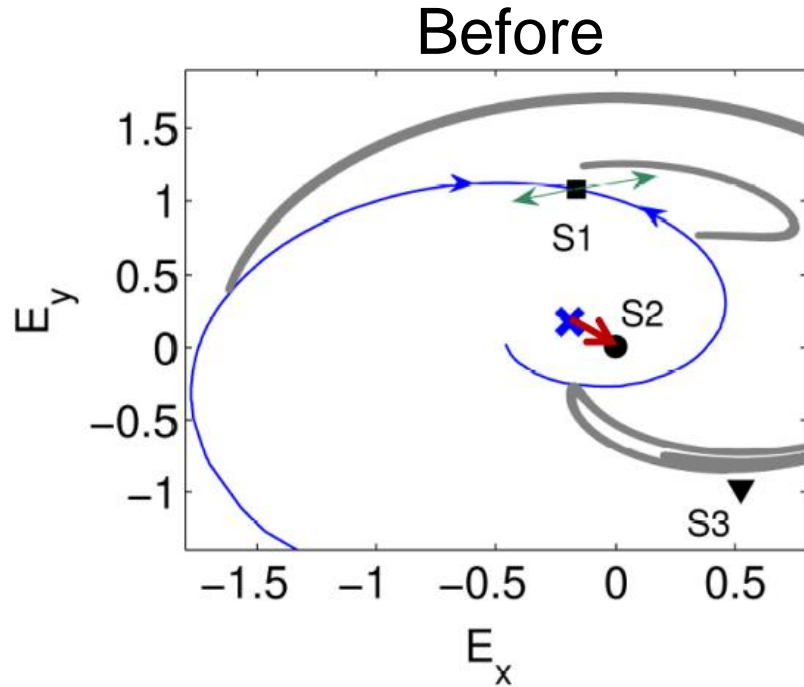
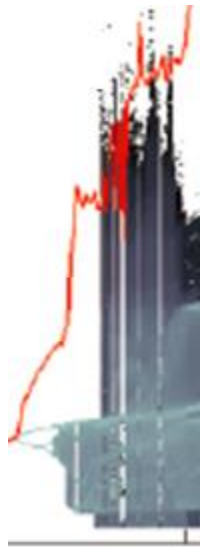
To understand the mechanism underlying the extreme pulses we need to examine the location the **three fixed points**.



An extreme pulse may be triggered when the trajectory closely approaches the stable manifold of S2 (“the door”)



With a Poincare map (N=1) we see the expansion of the attractor when extreme pulses appear



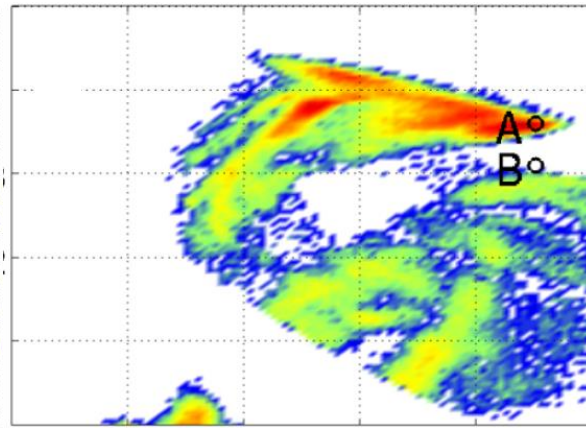
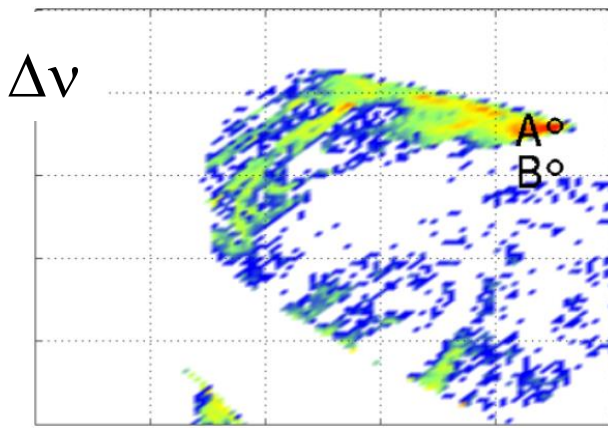
J. Zamora-Munt et al. PRA (2013).

Spontaneous emission noise can induce extreme pulses

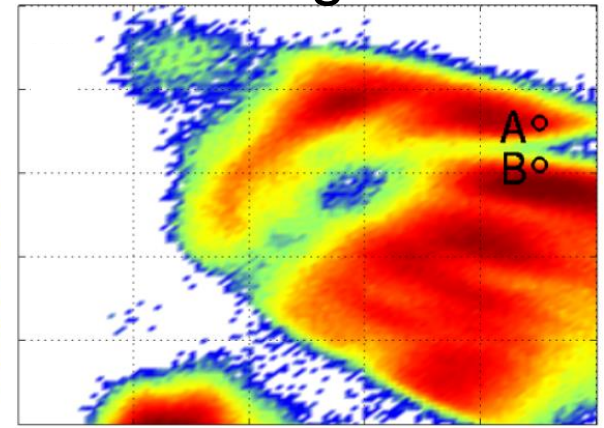
$$\frac{dE}{dt} = \kappa(1 + i\alpha)(N - 1)E + i\Delta\omega E + \sqrt{P_{\text{inj}}} + \sqrt{D}\xi(t),$$

In color code the number of pulses

No noise



Strong noise



Pump current

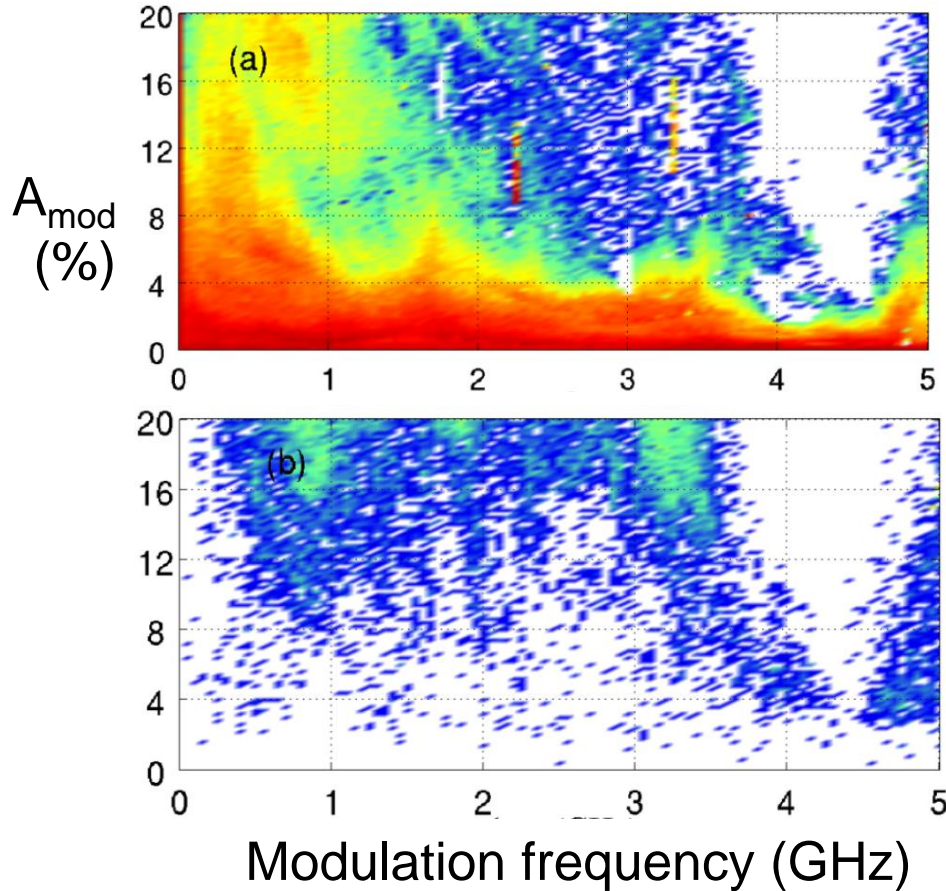
$$T_h = \langle l \rangle + 6\sigma$$

S. Perrone et al., PRA (2014).

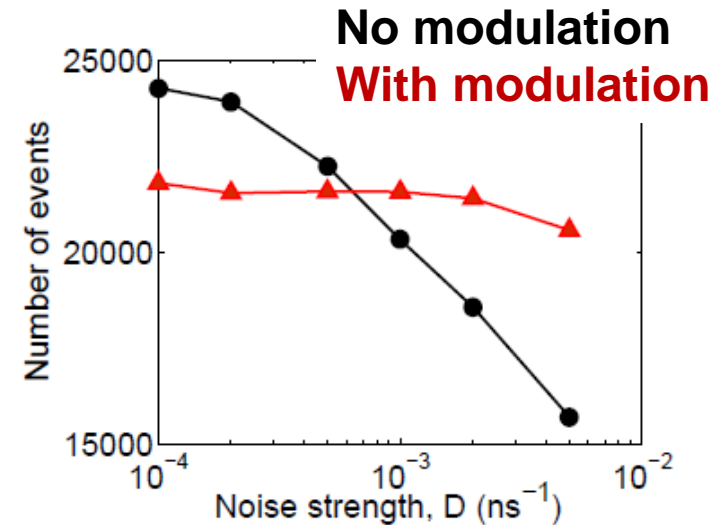
Current modulation suppresses or induces pulses

$$\frac{dN}{dt} = \gamma_N[\mu(t) - N - |E|^2]$$

$$\mu(t) = \mu_0 + \mu_{\text{mod}} \sin(\omega_{\text{mod}} t)$$



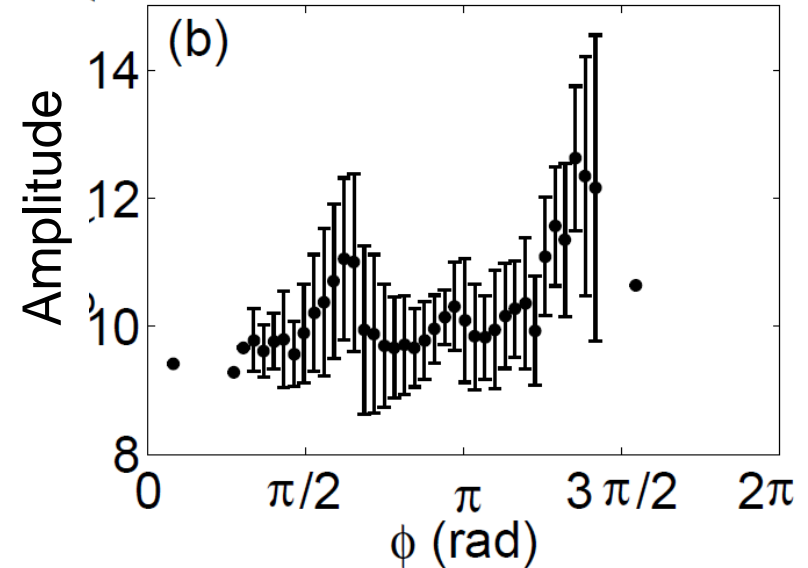
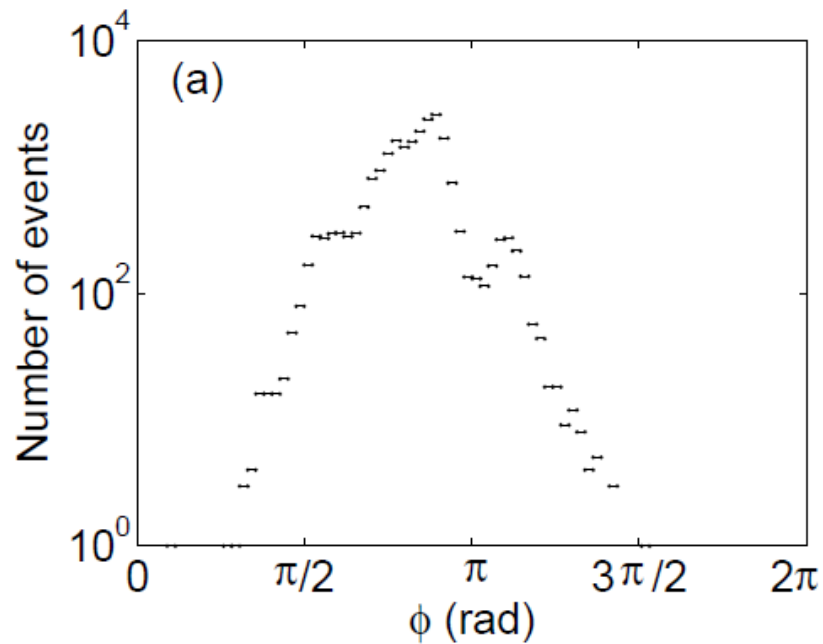
Color code:
log (# of pulses)



$T_h = \langle I \rangle + 6\sigma$

S. Perrone et al., PRA (2014).

When extreme pulses are not suppressed by current modulation, their probability and amplitude depend on the phase of the modulation

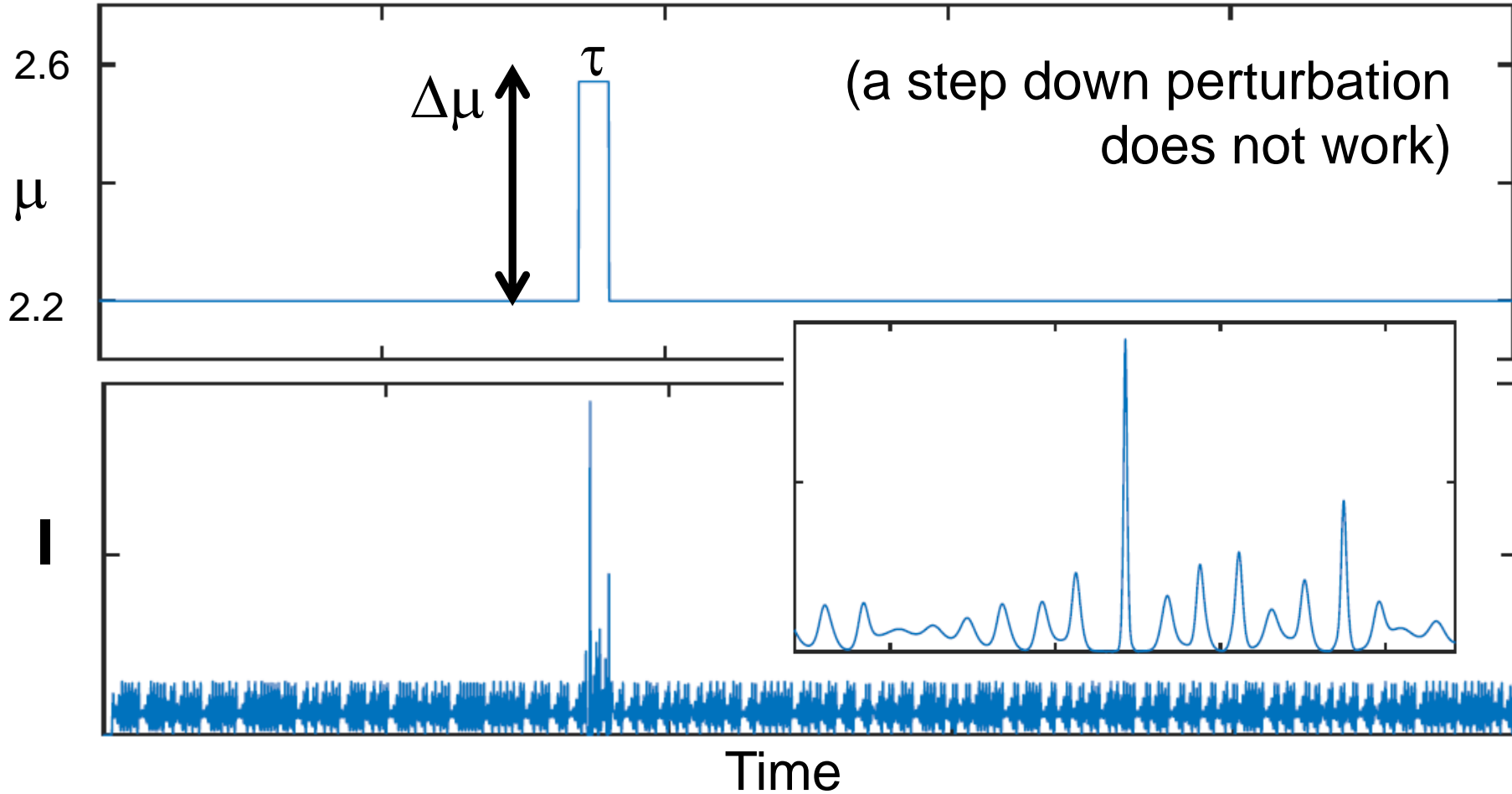


J. Ahuja et al., Opt. Express 2014

Acknowledgments

The work was supported by grants FA9550-14-1-0359, FIS2012-37655-C02-01, FIS2011-29734-C02-01 and ICREA Academia; C. M. and J. Z. M. also acknowledge the Max Planck Institute for the Physics of Complex Systems, Advanced Study Group on Optical Rare Events.

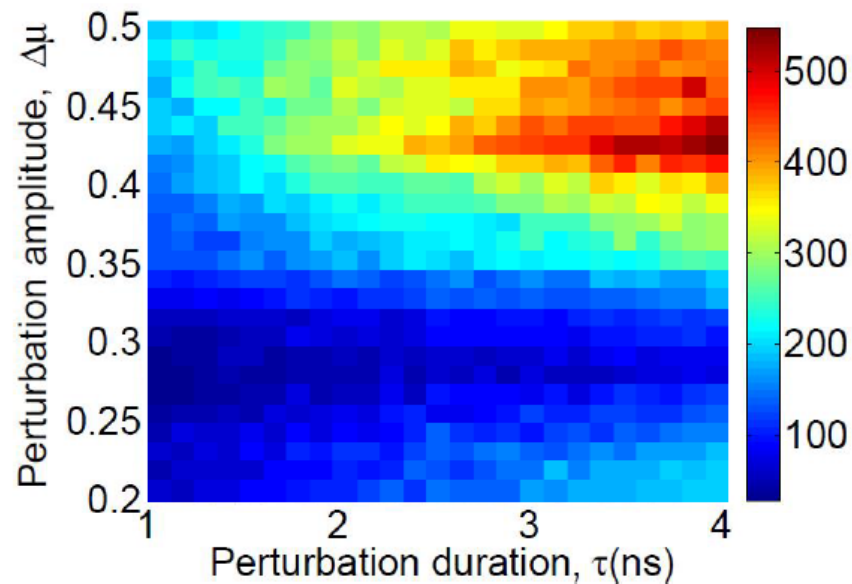
Can extreme pulses be generated “on demand”?



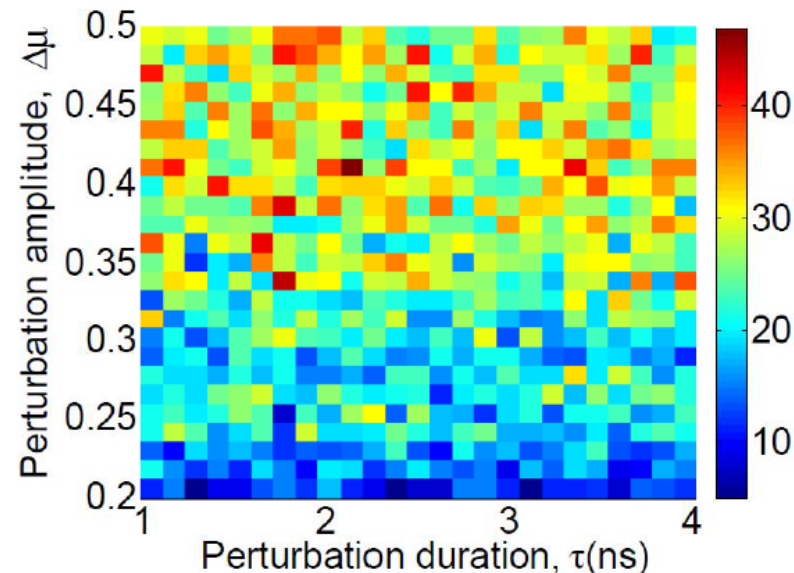
T. Jin et al, Opt. Express (2017).

Number of extreme pulses generated by 1000 perturbations as a function of the perturbation parameters: as large as 50% or as small as 5%

C ($\mu = 1.8$ and $\Delta\nu = 0.22$ GHz)



B ($\mu = 2.2$ and $\Delta\nu = 0.6$ GHz)

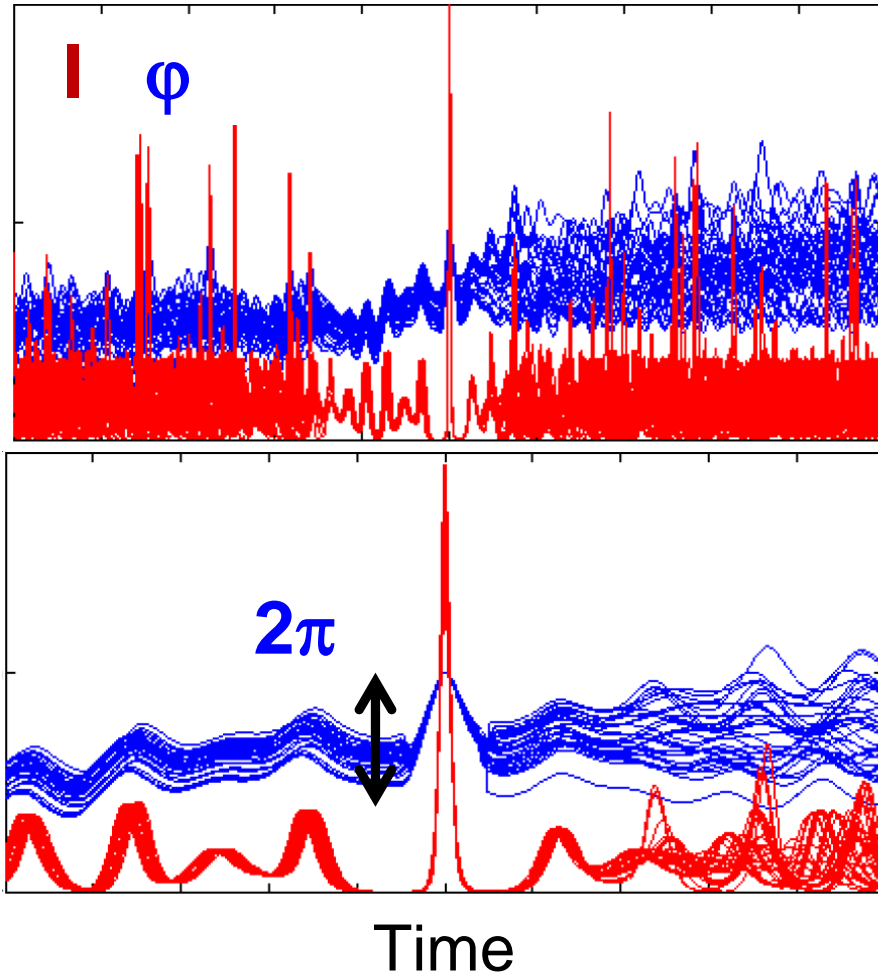


The “success rate” depends on the laser’s parameters and on the perturbation parameters.

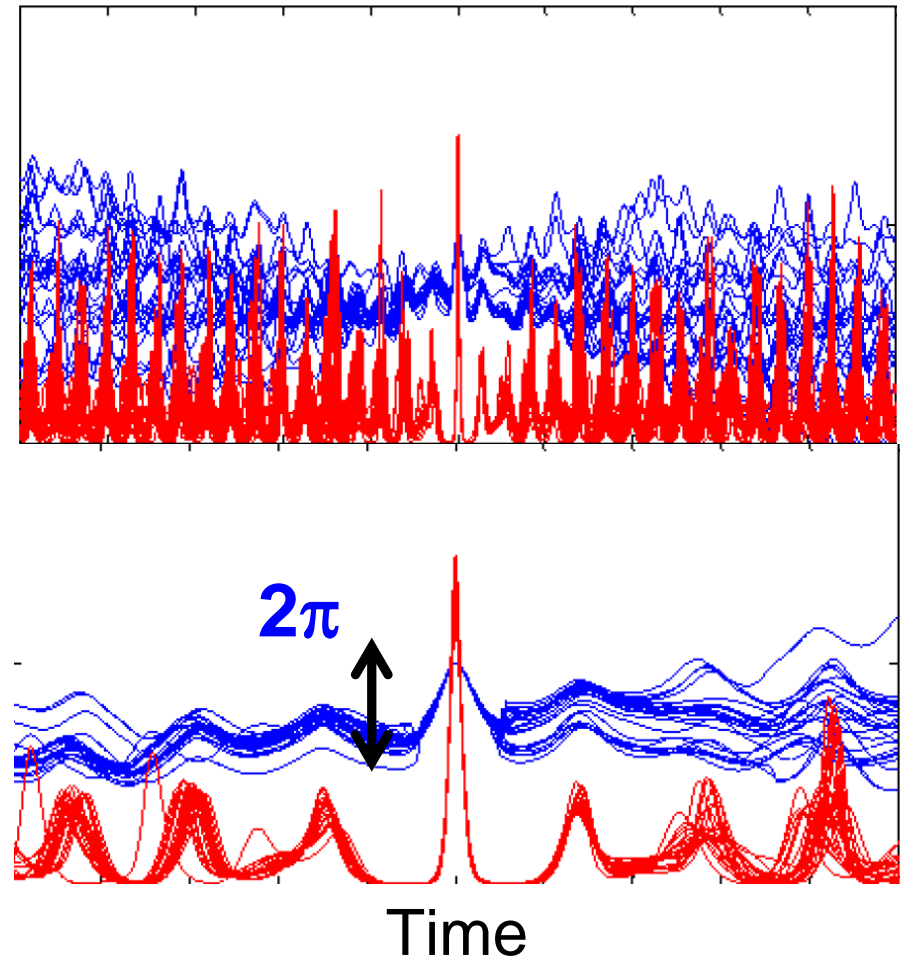
$$Th = \langle I \rangle + 8\sigma$$

Are the generated pulses similar to “natural” ones?

“Natural”



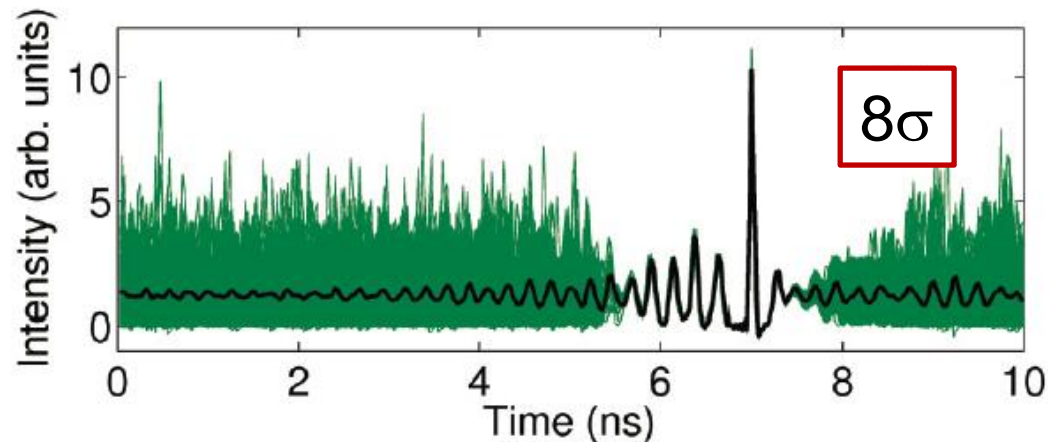
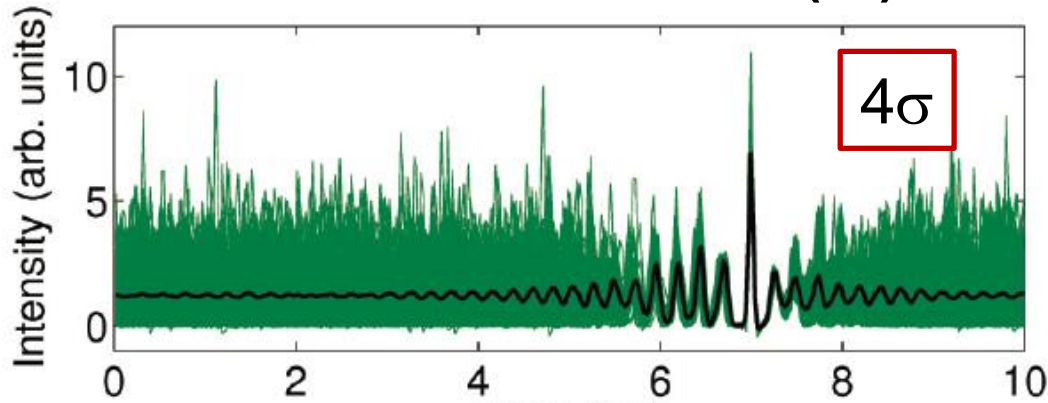
“On demand”



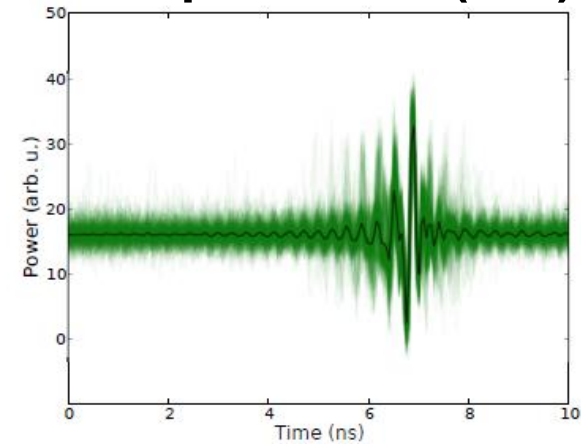
Predictability?

Superposition of 50/500 time series at the peak of the pulse

Deterministic simulations (50)



Experiments (500)

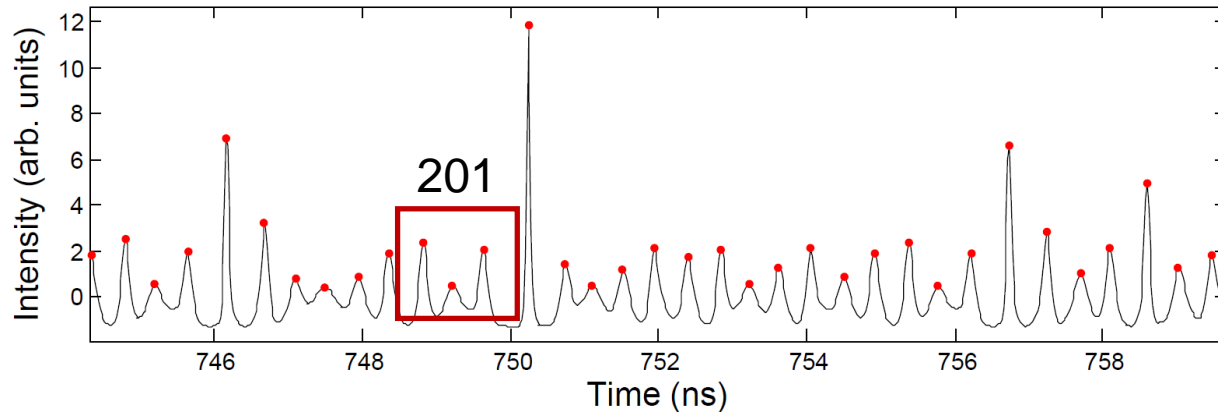


⇒ Well-defined oscillation pattern anticipates extreme pulses.

How can this effect be quantified?

J. Zamora-Munt et al. PRA (2013).

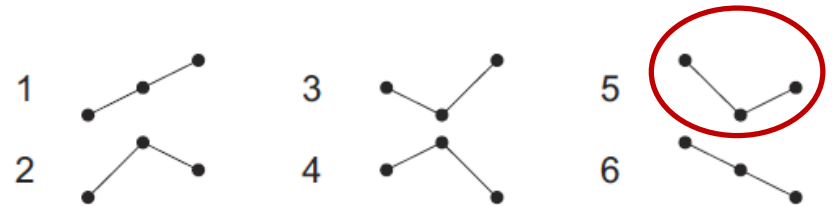
We try to identify a “pattern” that occurs before the pulse



- Consider the sequence of intensity peak heights (red dots):

$$\{\dots, l_i, l_{i+1}, l_{i+2}, \dots\}$$

- Possible order relations of three consecutive values:

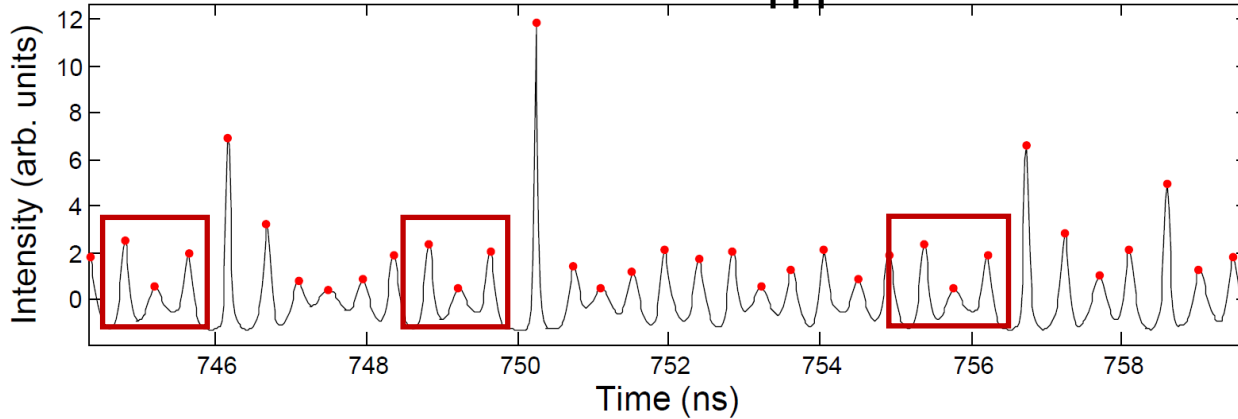
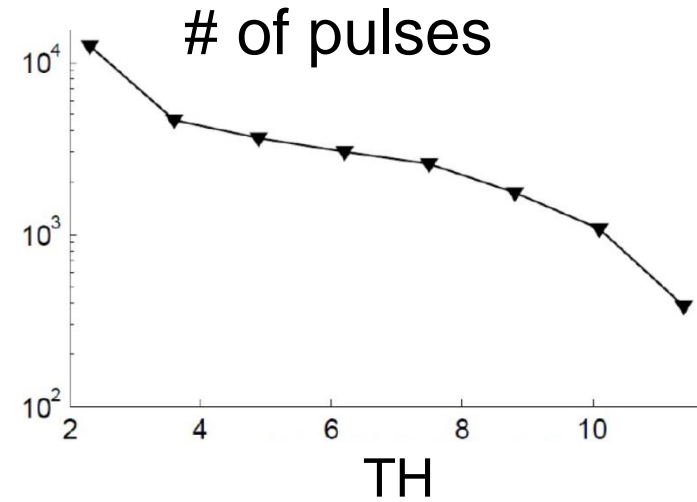
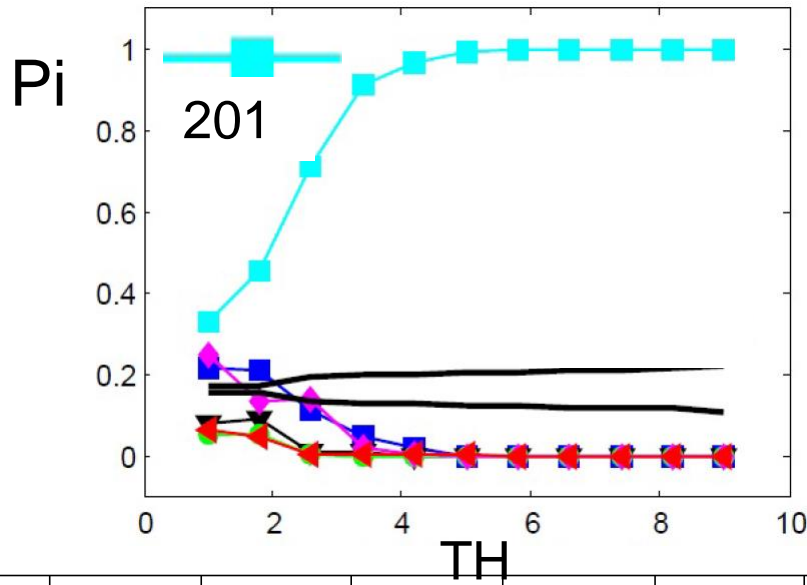


We calculate the probability of these pattern before each extreme pulse:

If $l_i > \mathbf{TH}$, we analyze the pattern defined by $(l_{i-3}, l_{i-2}, l_{i-1})$

“Good” results in deterministic simulations: $P(201)=1$ if $TH > 6$

Black lines:
 3σ
 confidence
 probabilities
 consistent
 with $p_i=1/6$
 $\forall i$

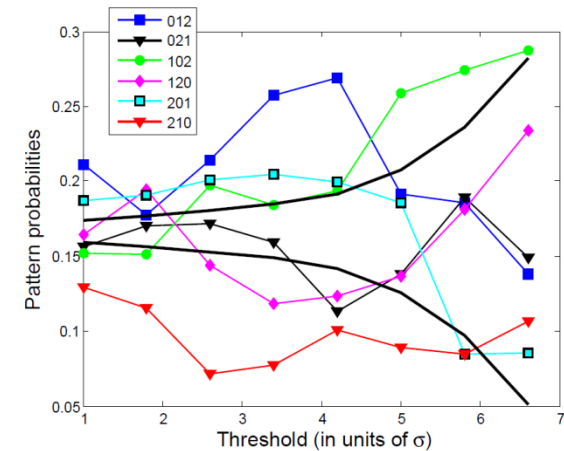
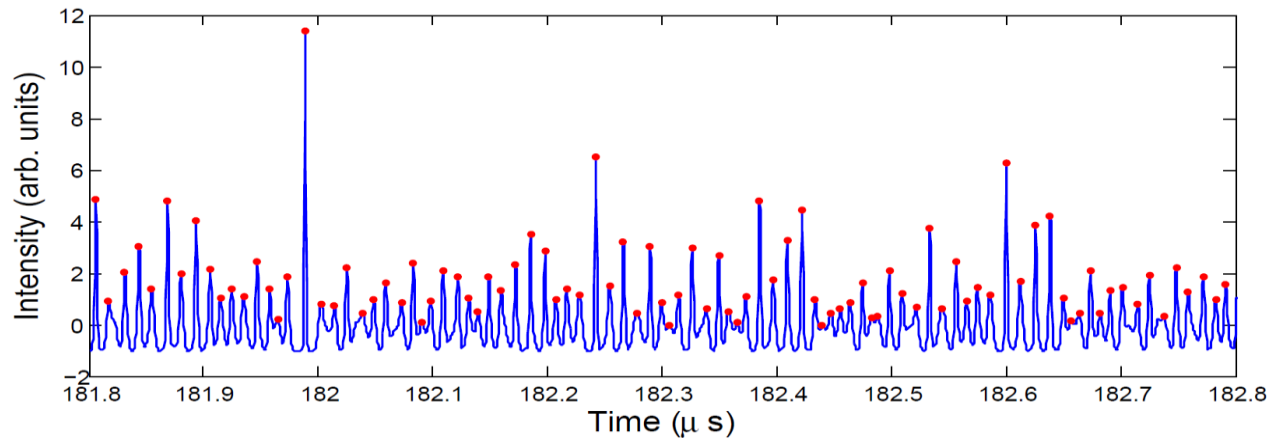
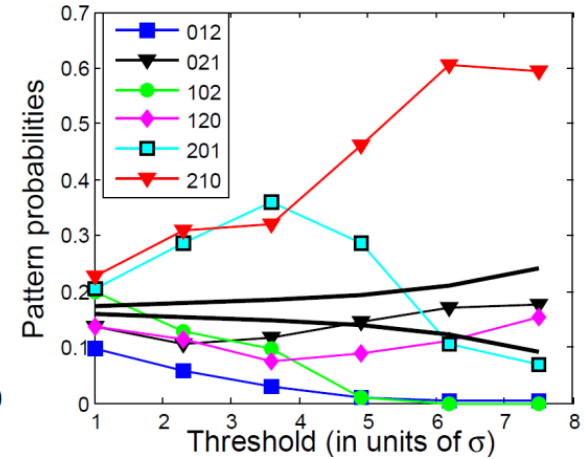
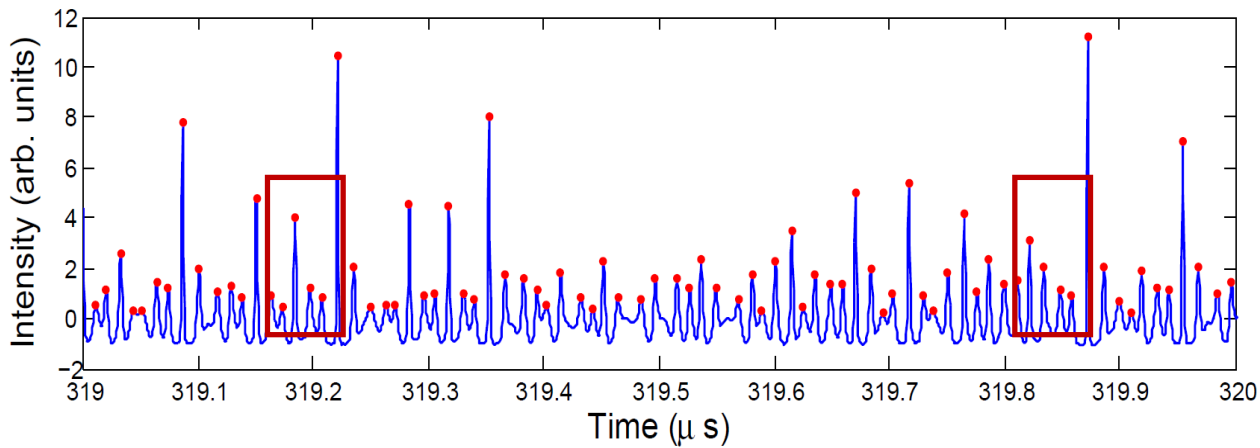


Problem: $P(201) \neq 0$ if $TH < 6$ (pattern 201 also anticipates some small pulses) \Rightarrow false alarms (false positives)

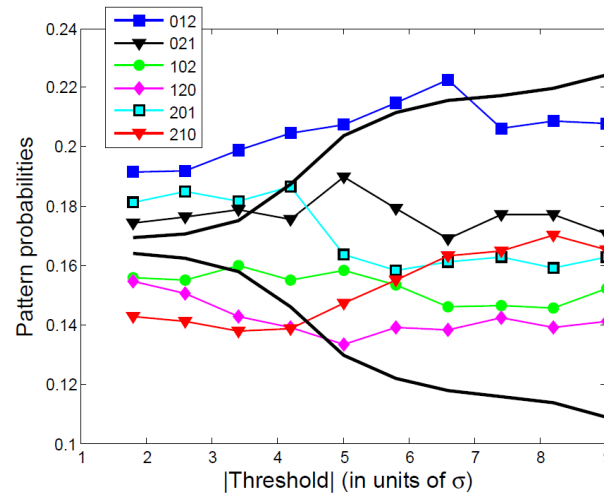
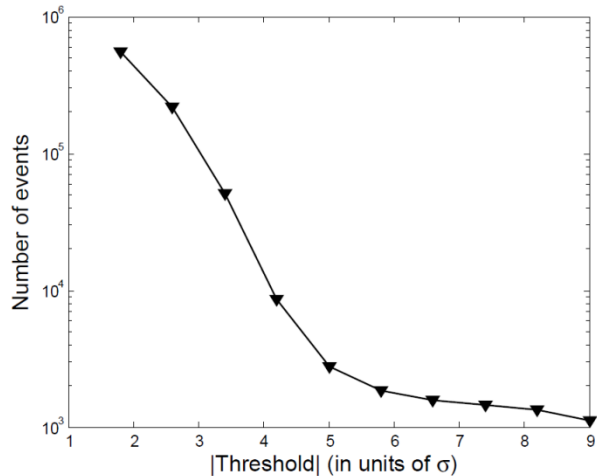
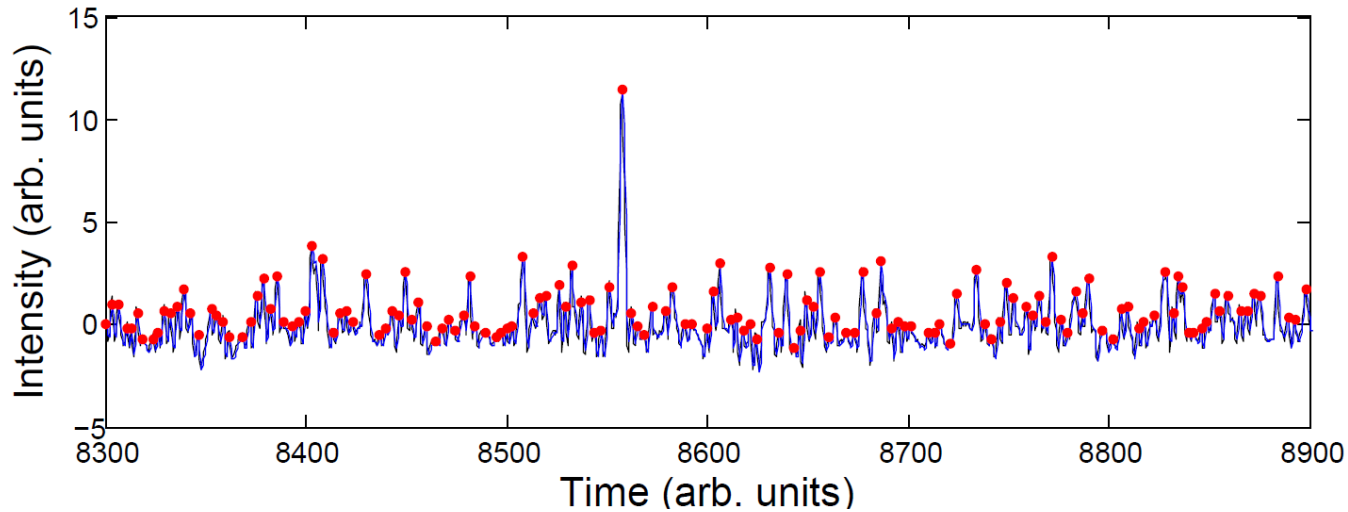
N. Martinez Alvarez et al., EPJST (2017).

The “early warning pattern” varies with the parameters and might not exist

Including noise, two modulation frequencies



Analysis of experimental data



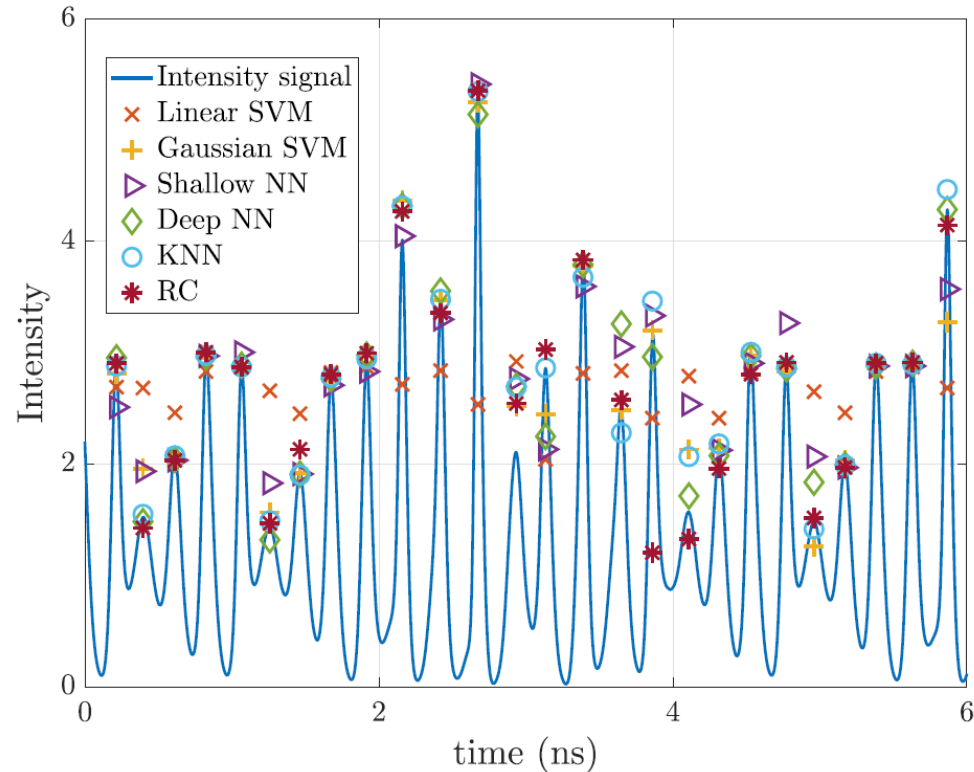
Way to improve the “early warning”:

- Filter noise
- Longer pattern $\{I_{i-4}, I_{i-3}, I_{i-2}, I_{i-1}\}$

Can the amplitude of the next pulse be predicted?

$$I_i = f(I_{i-n} \dots I_{i-3}, I_{i-2}, I_{i-1})$$

- Support Vector Machine (SVM), Linear and Gaussian
- Neural Networks (NN), Shallow and Deep
- k-Nearest Neighbors (KNN)
- Reservoir Computing (RC)



$n = 3$ yields minimum prediction error
(increasing n does not increase the accuracy).

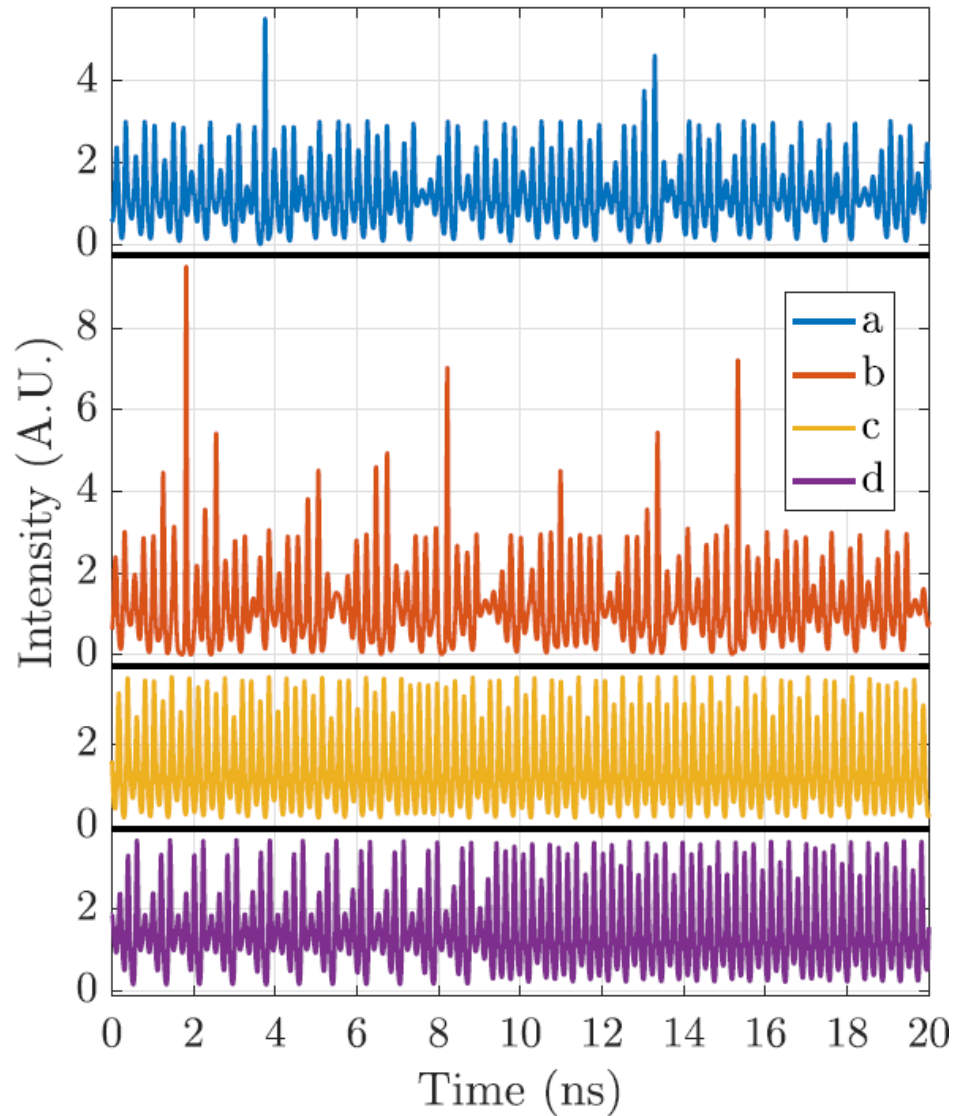
P. Amil et al., Chaos (2019)

Influence of the pump current and noise?

$$\mu = 2.2$$

$$D = 0$$

$$D = 10^{-4} \text{ ns}^{-1}$$



$$\mu = 2.45$$

$$D = 0$$

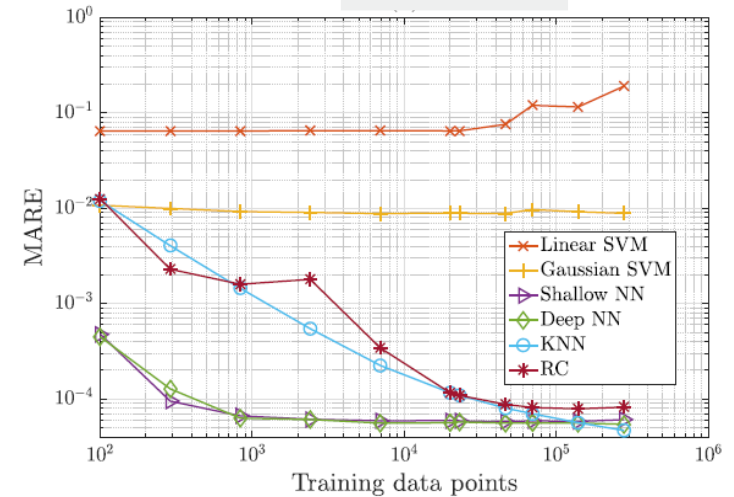
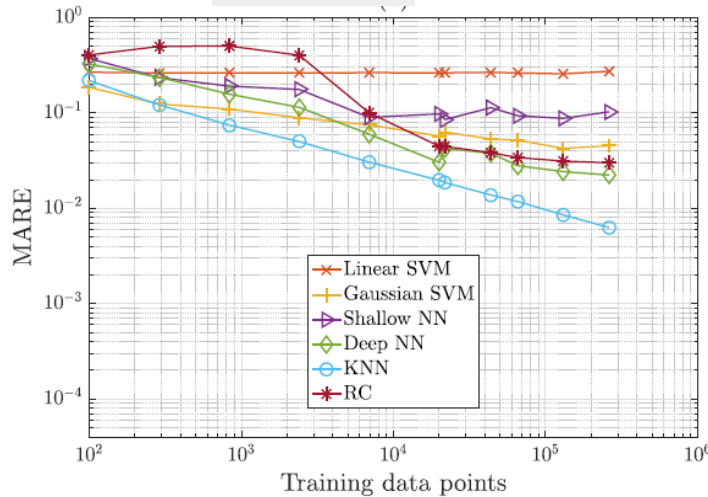
$$D = 10^{-4} \text{ ns}^{-1}$$

Performance quantification: the mean absolute relative error

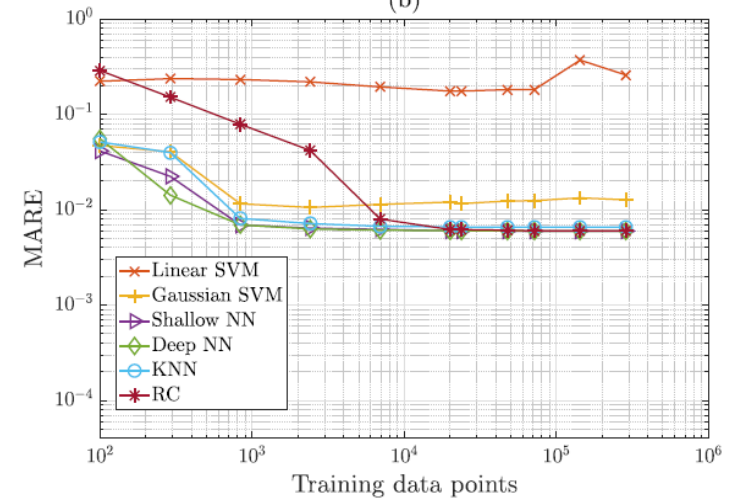
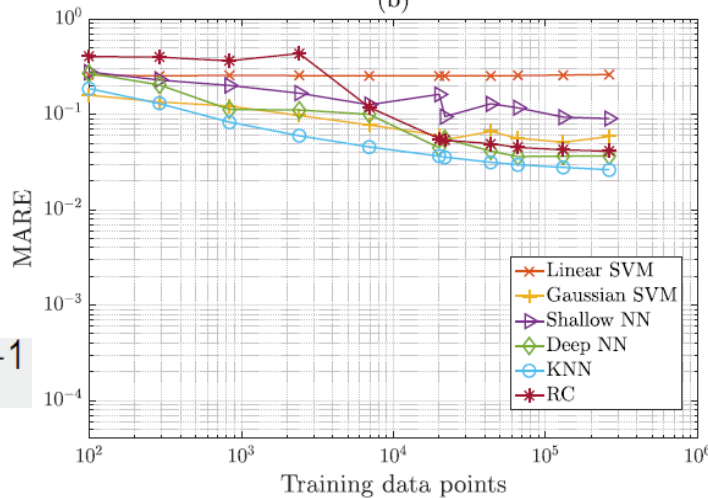
$$\text{MARE} = \frac{1}{N} \sum_{i=1}^{i=N} \frac{|\tilde{y}_i - y_i|}{y_i}$$

$\mu = 2.2$

$\mu = 2.45$



$D = 0$



$D = 10^{-4} \text{ ns}^{-1}$

Summary

- In optically injected semiconductor lasers, extreme pulses can be deterministic, or triggered by noise, and can be suppressed or induced by current modulation.
- Extreme pulses can be generated “on demand” by a small perturbation of the pump current.
- The pulse amplitude can be predicted with good accuracy, even for extreme pulses that have very low probability.
- Future work: potential for sensing applications?
- Future work: the symbolic approach to predict extreme events needs to be further explored.

Co-authors: S. Barland, M. Giudici, J. R. Tredicce, J. R. Rios Leite, J. Zamora, S. Perrone, R. Vilaseca, P. Amil, M. C. Soriano, N. Martinez Alvarez, S. Borkar, J. Ahuja, D. Bhiku Nalawade, T. Jin and C. Siyu

Thank you for your attention!