

Nonlinear laser and neural dynamics: fun and challenges of complex physical phenomena

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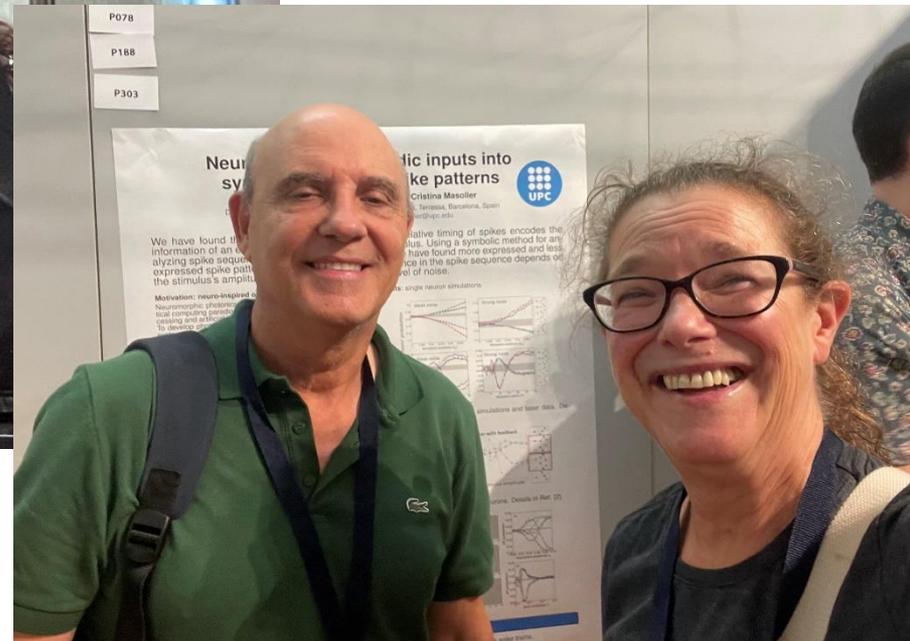
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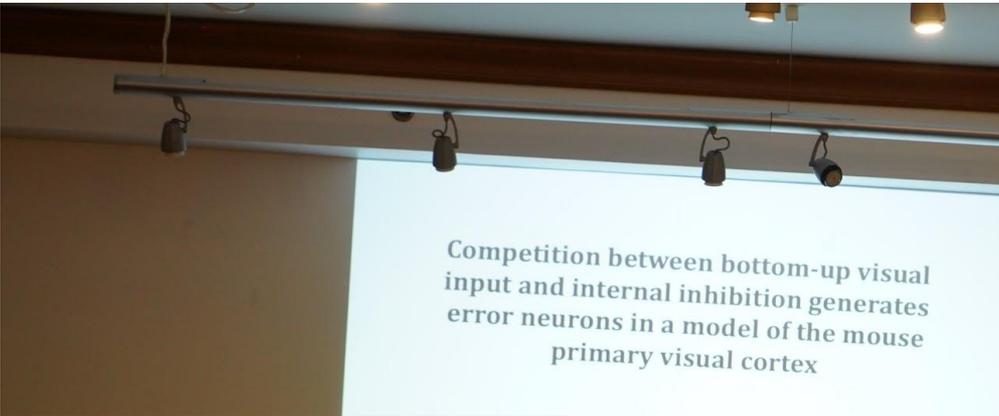
Mallorca 2001



Florence 2025



MEDYFINOL 2022



Our first collaboration: synchronization of chaotic lasers for secure communications

456

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Chaos Shift-Keying Encryption in Chaotic External-Cavity Semiconductor Lasers Using a Single-Receiver Scheme

Claudio R. Mirasso, Josep Mulet, and Cristina Masoller

Abstract—In this letter, we numerically show that chaos shift-keying (CSK) encryption can be achieved by using a single receiver, thus providing a better performance when comparing with the traditional CSK scheme based on two receivers. We analyze the rate equation model for two unidirectionally coupled single-mode external-cavity semiconductor lasers operating in a chaotic regime. The message is encoded in the emitter by slightly varying its injection current. We find that under appropriate conditions, the receiver laser synchronizes to the chaotic emitter, filtering the encoded message and allowing message extraction.

Index Terms—Injection-locked oscillator, nonlinear optics, semiconductor lasers (dynamics).

scheme [2], [8], the carrier is modulated by the message. On the other hand, the CSK scheme [9] is based on the definition of two clearly separated states for bits “1” and “0,” while in the OOSK the system synchronized either to a bit “1” or “0” being unsynchronized for the other bit [10], [11]. In all these methods, the intensity of the message has to be small enough to avoid detection in the time or frequency domains. To the best of our knowledge, in the proposed CSK schemes the decoder consist of two replicas of the transmitting systems, each one configured for detecting either a bit “1” or “0” [9]. Although more secure, this scheme becomes more complicated to implement than the

Modeling bidirectionally coupled single-mode semiconductor lasers

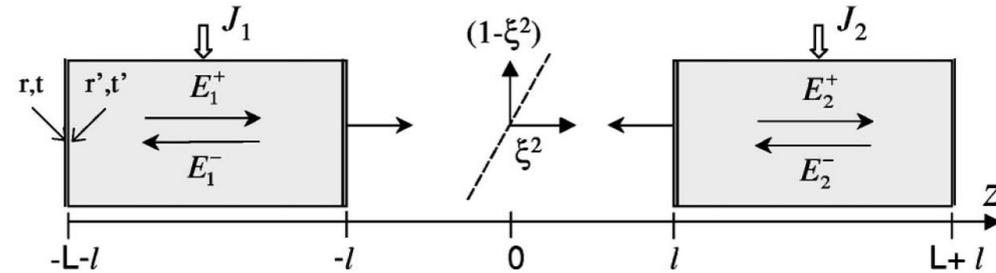
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$$d_t A_{1,2}(t) = \mp i \Delta A_{1,2}(t) + \frac{1}{2} (1 - i\alpha) \gamma [\mathcal{G}_{1,2} - 1] A_{1,2}(t) + \hat{\kappa}_c A_{2,1}(t - \tau), \quad (31a)$$

$$d_t D_{1,2}(t) = \gamma_e [\mu_{1,2} - D_{1,2} - \mathcal{G}_{1,2} |A_{1,2}|^2], \quad (31b)$$

$$\mathcal{G}_{1,2} = \frac{a D_{1,2}}{1 + \varepsilon |A_{1,2}|^2}, \quad (31c)$$

Anticipating the Response of Excitable Systems Driven by Random Forcing

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We study the regime of anticipated synchronization in unidirectionally coupled model neurons subject to a common external aperiodic forcing that makes their behavior unpredictable. We show numerically and by analog hardware electronic circuits that, under appropriate coupling conditions, the pulses fired by the slave neuron anticipate (i.e., predict) the pulses fired by the master neuron. This anticipated synchronization occurs even when the common external forcing is white noise.

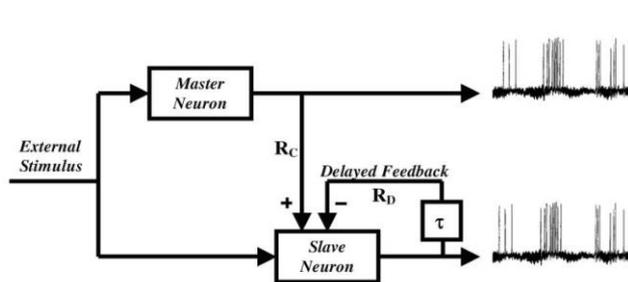
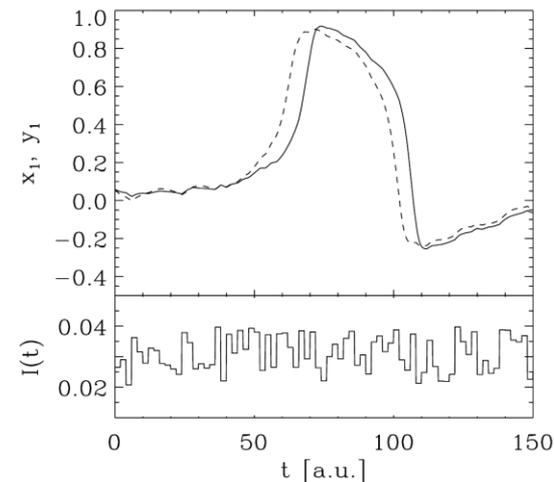


FIG. 1. Schematic diagram of two model neurons coupled in a unidirectional configuration, subjected to the same external forcing and with a feedback loop (with a delay time τ) in the slave neuron.



Anticipating the dynamics of chaotic maps

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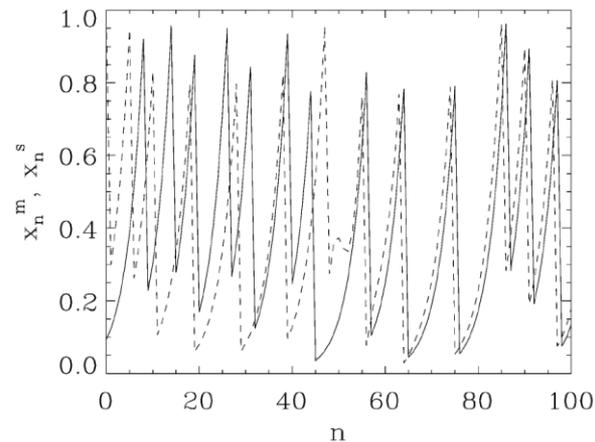


Fig. 4. Time series of the 2D Baker map (x_n^m solid line, x_n^s dashed line) for $a = 1.333$, $b = 0.777$, $\gamma = 0.7$, and $M = 1$. The coupling is set on at $n = 30$, leading to anticipated synchronization.

Synchronization regimes of optical-feedback-induced chaos in unidirectionally coupled semiconductor lasers

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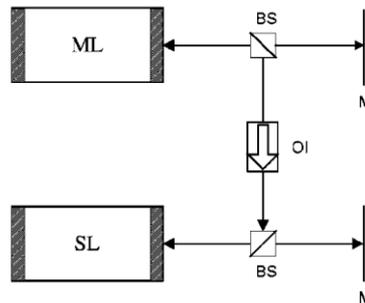
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(Received 5 January 2002; published 29 April 2002)

We numerically study the synchronization of two unidirectionally coupled single-mode semiconductor lasers in a master-slave configuration. The master laser is an external-cavity laser that operates in a chaotic regime while for the slave laser we consider two configurations. In the first one, the slave laser is also an external-cavity laser, subjected to, its own optical feedback and the optical injection from the master laser. In the second one, the slave laser is subject only to the optical injection from the master laser. Depending on the operating conditions the synchronization between the two lasers, whenever it exists, can be either isochronous or anticipated. We perform a detailed study of the parameter regions in which these synchronization regimes occur and how small variations of parameter yield one or the other type of synchronization or an unsynchronized regime.

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PACS number(s): 05.45.Xt, 42.55.Px, 42.65.Sf



95 WOS citations

Almost 20 years later.

Characterizing signal encoding and transmission in class I and class II neurons via ordinal time-series analysis

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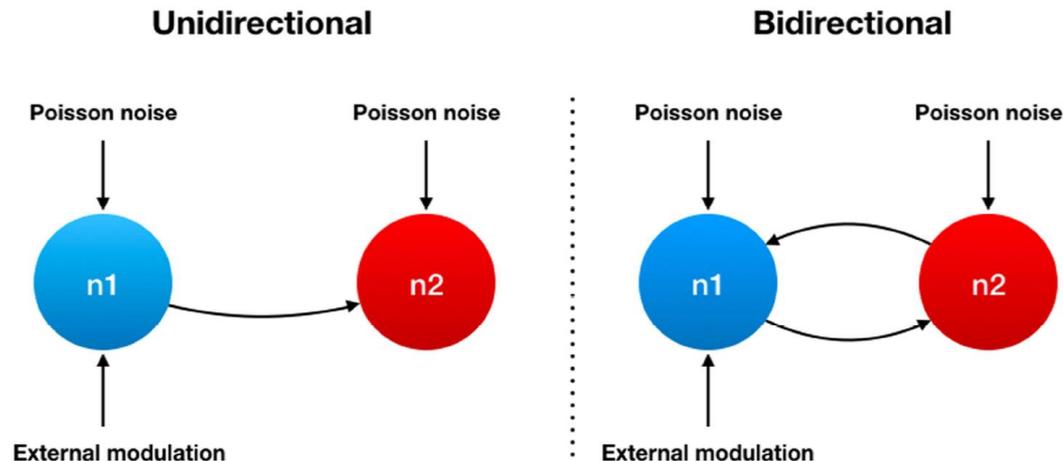
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C. Estarellas,^{1,a)}  M. Masoliver,^{2,a)}  C. Masoller,²  and Claudio R. Mirasso¹ 

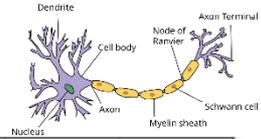
AFFILIATIONS

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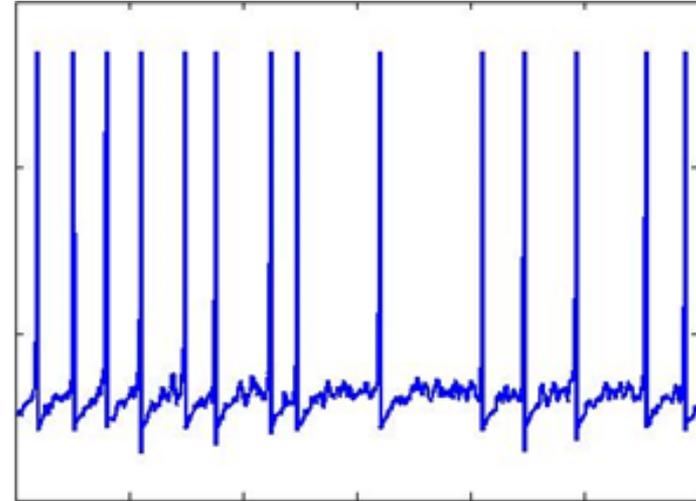
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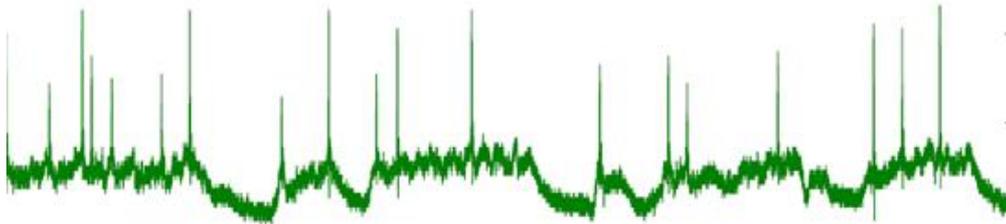
Intensity emitted by a diode laser and simulated spikes



Time 10^{-9} s

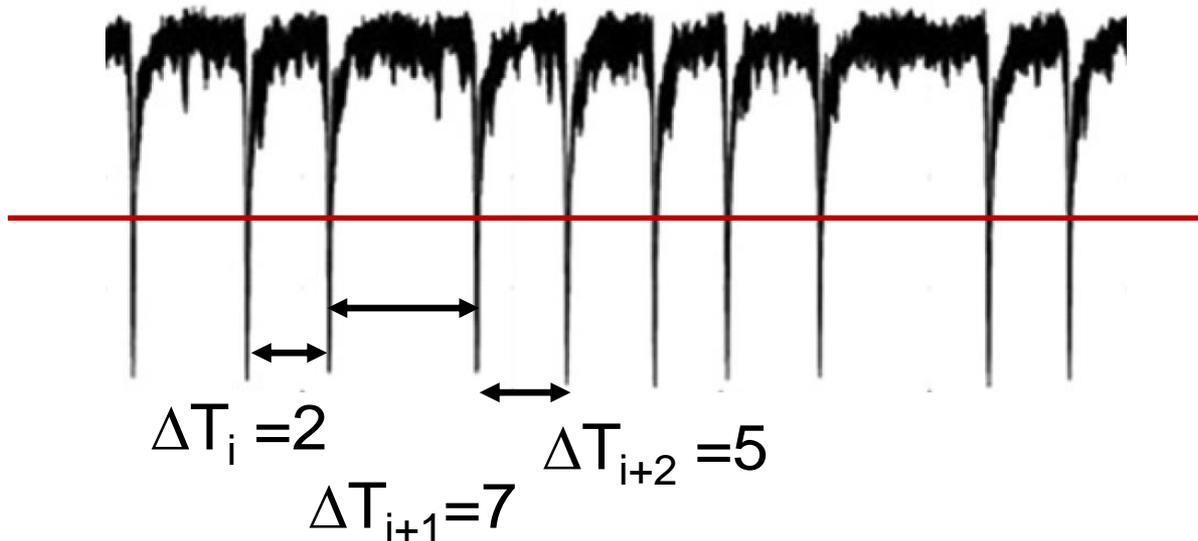
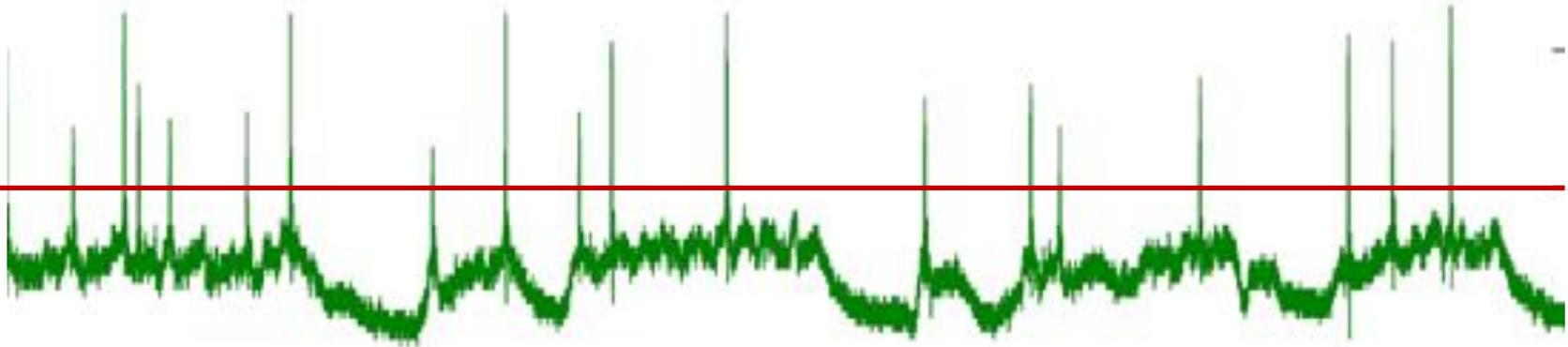


Time 10^{-3} s



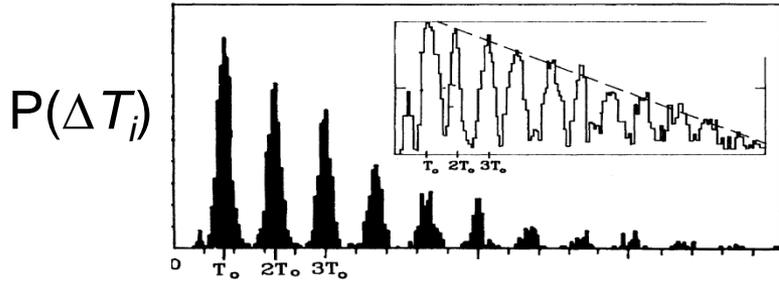
How neurons encode information?

Thresholding detects the spike times \Rightarrow Point Process
 \Rightarrow analyze sequence of inter-spike-intervals (ISIs)



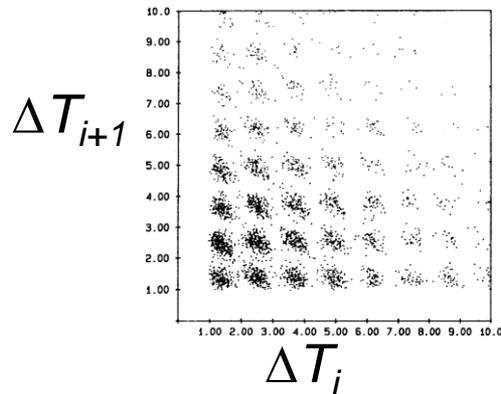
Neuron inter-spike-intervals (ISI) distribution

(spikes in the auditory nerve when a monkey hears a pure tone sound)

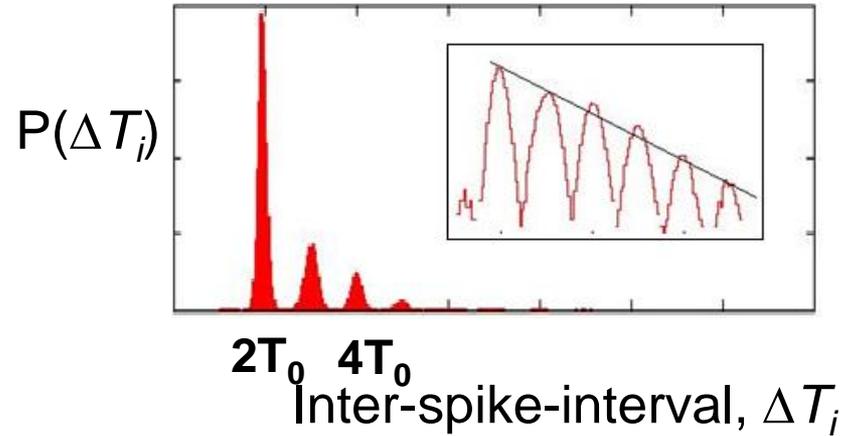


Inter-spike-interval, ΔT_i

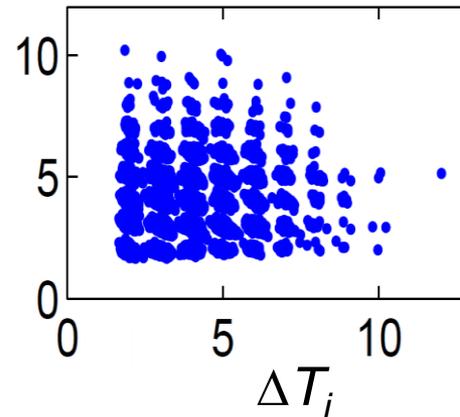
Auditory nerve when a cat hears a pure tone sound:



Diode laser ISI distribution (when the laser is sinusoidally modulated)



How to detect similar temporal order in the ISI sequences?



A. Longtin et al. PRL (1991), IJBC (1993).

A. Aragoneses et al. Opt Express (2014).

Analysis method: ordinal analysis

$$\{\dots X_i, X_{i+1}, X_{i+2}, \dots\}$$

Possible order relations among three numbers (e.g., 2, 5, 7)

$\{\dots 2, 5, 7 \dots\}$

$\{\dots 5, 2, 7 \dots\}$

$\{\dots 7, 2, 5 \dots\}$



$\{\dots 2, 7, 5 \dots\}$

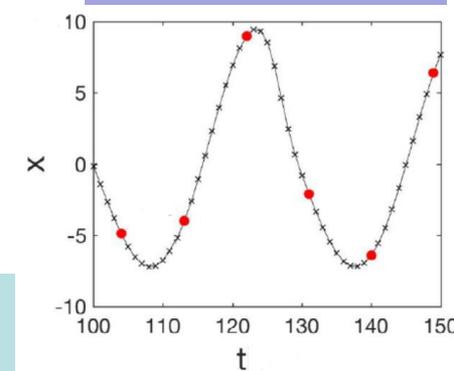


$\{\dots 5, 7, 2 \dots\}$



$\{\dots 7, 5, 2 \dots\}$

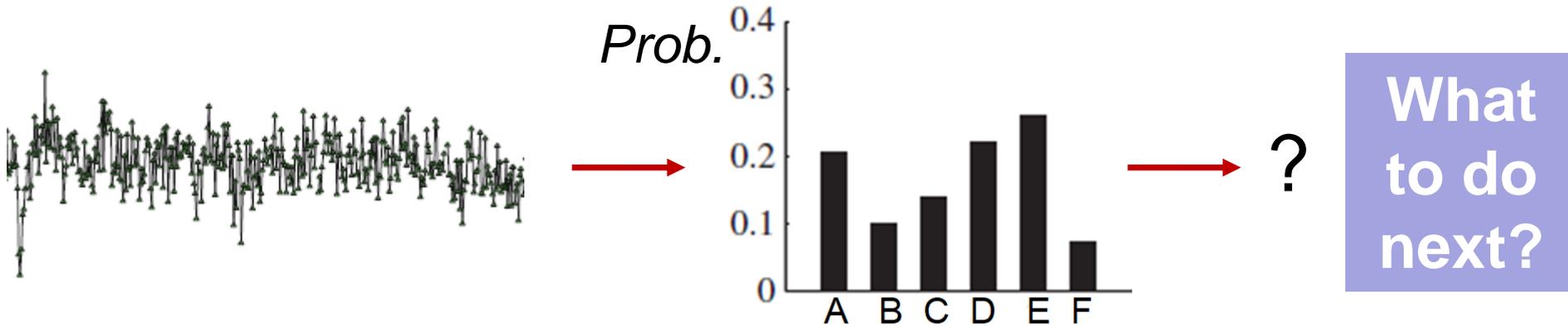
Which is the code?



A B F C

C. Bandt and B. Pompe, Phys. Rev. Lett. 88, 174102 (2002).

From a sequence of data points, by counting the different patterns we can estimate the “ordinal probabilities”



1. Permutation Entropy:

$$p_i = p_j \text{ for all } i, j \Rightarrow H=1$$

$$p_i = 1, p_j = 0 \text{ for all } j \neq i \Rightarrow H=0$$

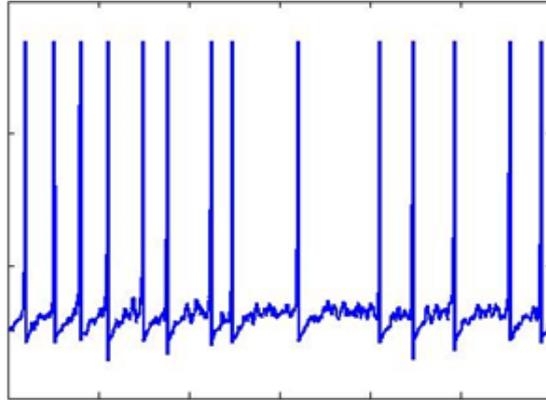
$$H = -\frac{1}{\ln N!} \sum_{i=1}^N p_i \ln p_i$$

(Nonlinear dimensionality reduction)

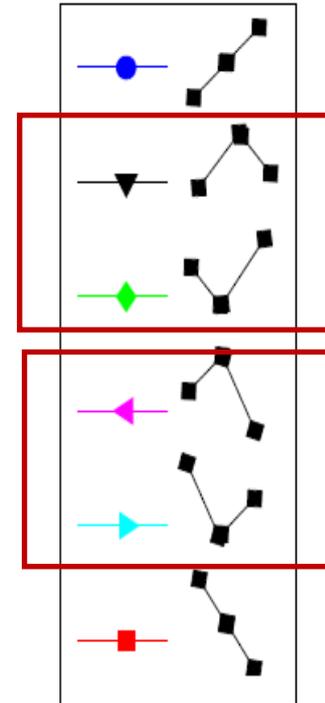
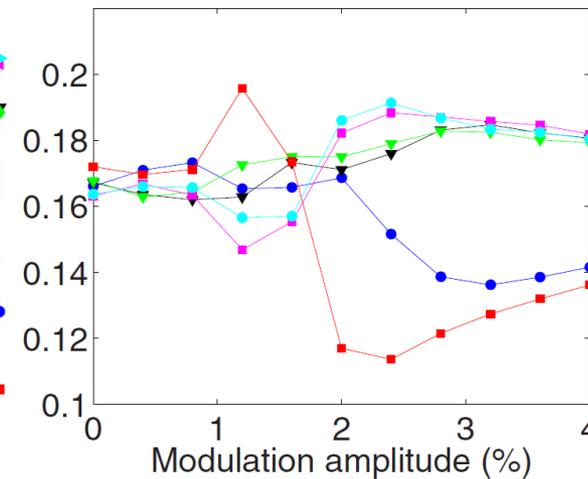
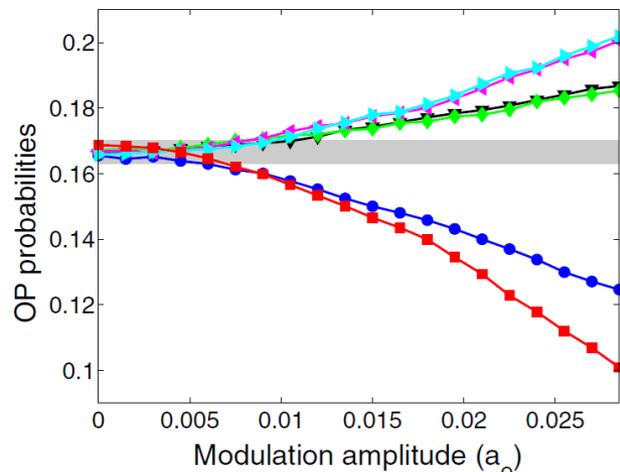
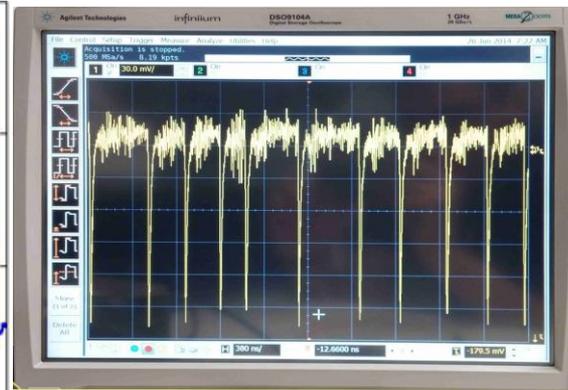
2. Analyze all the probabilities (features for machine learning algorithms)

Comparing the spike timing of a neuron and a laser when they are modulated with a *subthreshold* sinusoidal signal

Neuron model



Diode laser with feedback



J. M. Aparicio-Reinoso et al PRE 94, 032218 (2016) A. Aragoneses et al, Sci. Rep. 4, 4696 (2014)

Characterizing signal encoding and transmission in class I and class II neurons via ordinal time-series analysis

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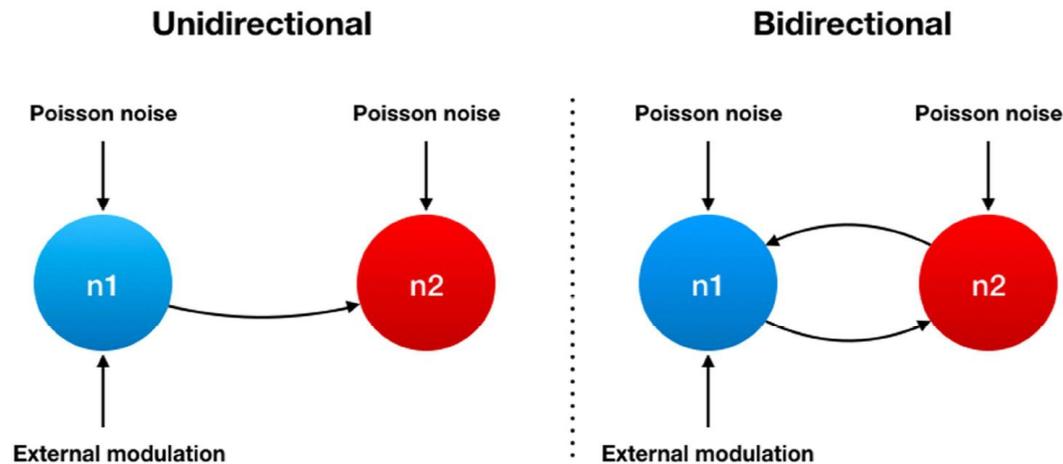
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C. Estarellas,^{1,a)}  M. Masoliver,^{2,a)}  C. Masoller,²  and Claudio R. Mirasso¹ 

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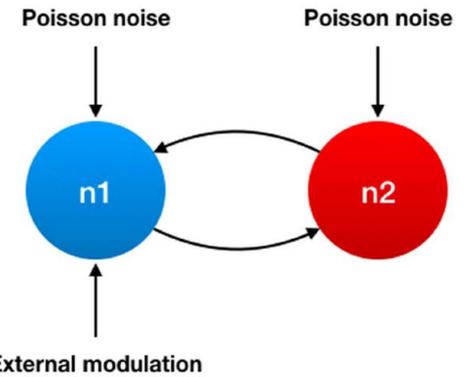
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Neuron 1

Neuron 2

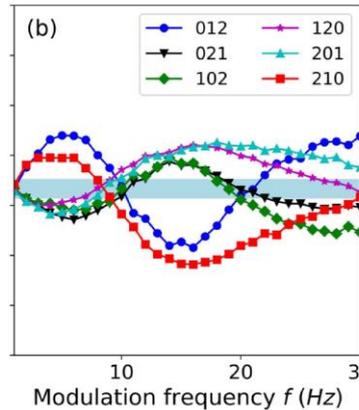
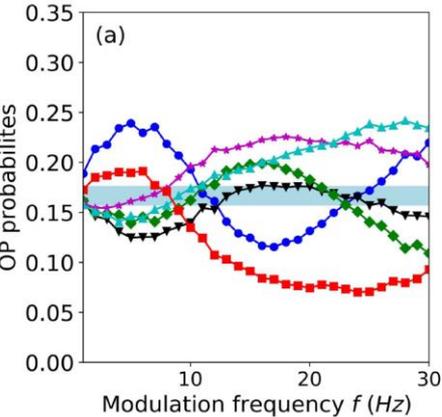
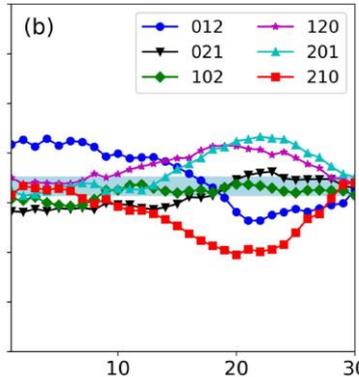
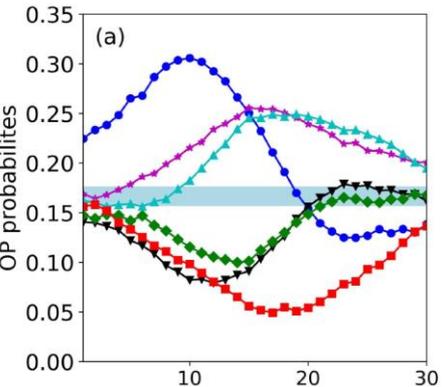
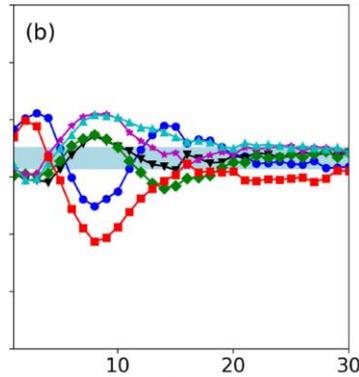
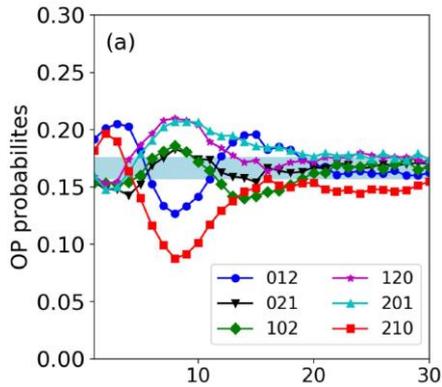


Class I
electrical
coupling.

Class I
chemical
coupling.

Class II
chemical
coupling.

⇒ Depending on the signal frequency, specific combinations of neuron-class and coupling-type allow better encoding or better transmission of the signal.



1. “Self pulsating lasers” (studied by Claudio and co-workers more than 25 years ago).

Self-Pulsating Semiconductor Lasers: Theory and Experiment

Claudio R. Mirasso, Guido H. M. van Tartwijk, E. Hernández-García, Daan Lenstra, S. Lynch, P. Landais, P. Phelan, J. O’Gorman, M. San Miguel, and W. Elsässer

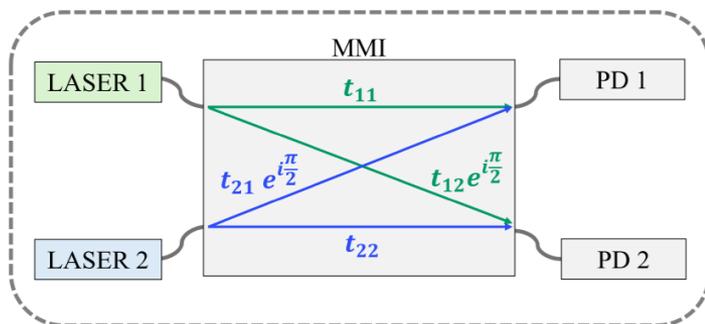
Abstract—We report detailed measurements of the pump-current dependency of the self-pulsating frequency of semiconductor CD lasers. A distinct kink in this dependence is found and explained using a rate-equation model. The kink denotes a transition between a region where the self-pulsations are weakly sustained relaxation oscillations and a region where Q -switching takes place. Simulations show that spontaneous emission noise plays a crucial role in the cross-over.

Index Terms—CD lasers, laser noise, self-pulsations, semiconductor lasers.

In the case of self-pulsations caused by saturable-absorbing effects, the self-pulsation frequency (SPF) dependence on the pump current was investigated in [2]. In later works, the self-pulsations were attributed to undamped relaxation oscillations (RO’s) [3], [4]. The precise values of the RO frequency (ROF), as calculated from a small-signal analysis, and the actual SPF, highly nonlinear, are, however, different, as the SPF is always smaller than the ROF [5].

Saturable absorption effects, causing self-pulsations in

2. Synchronization of mutually coupled lasers



InP PIC with two DFB lasers and a multimode interferometer (MMI) for quantum random number generation.

$$I_{PD1}(t) = |t_{11}E_1(t) + t_{21}e^{i\frac{\pi}{2}}E_2(t)|^2$$

Happy birthday Claudio!!!!